

## xAct‘xPert‘

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### Authors

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### Intro

xPert‘ is a package for efficient construction and manipulation of high–order perturbations of the curvature tensors of a Riemannian manifold. Its main application is Perturbation Theory in General Relativity.

xPert‘ is part of the xAct‘ framework for Tensor Computer Algebra in *Mathematica*, and in particular it needs the underlying package xTensor‘ for abstract tensor manipulations.

### Load the package

---

This loads the package from the default directory, for example `$Home/.Mathematica/Applications/xAct/` for a single–user installation under Linux. xTensor‘ and xPerm‘ are automatically loaded.

```
In[1]:= MemoryInUse[]
```

```
Out[1]= 2125968
```

```
In[2]:= <<xAct`xPert`
```

```
-----
--
Package xAct`xCore` version 0.6.0, {2008, 6, 30}
CopyRight (C) 2007-2008, Jose M.
Martin-Garcia, under the General Public License.
-----
--
Package ExpressionManipulation`
CopyRight (C) 1999-2008, David J. M. Park and Ted Ersek
-----
--
Package xAct`xPerm` version 1.0.1, {2008, 5, 16}
CopyRight (C) 2003-2008, Jose M.
Martin-Garcia, under the General Public License.
Connecting to external linux executable...
Connection established.
-----
--
Package xAct`xTensor` version 0.9.6, {2008, 6, 30}
CopyRight (C) 2002-2008, Jose M.
Martin-Garcia, under the General Public License.
-----
--
Package xAct`xPert` version 1.0.0, {2008, 6, 30}
CopyRight (C) 2005-2008, David Brizuela, Jose M. Martin-Garcia
and Guillermo A. Mena Marugan, under the General Public License.
-----
--
These packages come with ABSOLUTELY NO WARRANTY; for details type
Disclaimer[]. This is free software, and you are welcome to redistribute
it under certain conditions. See the General Public License for details.
-----
--
```

---

Comparing, we see that the packages take about 10Mb in *Mathematica* 5.2:

```
In[3]:= MemoryInUse[]
```

```
Out[3]= 11734032
```

```
In[4]:= Out[3] - Out[1]
```

```
Out[4]= 9608064
```

---

There are seven contexts: `xAct`xPert``, `xAct`xTensor``, `xAct`xPerm``, `xAct`xCore`` and `xAct`ExpressionManipulation`` contain the respective reserved words. `System`` contains *Mathematica*'s reserved words. The current context `Global`` will contain your definitions and right now it is empty.

```
In[5]:= $ContextPath
```

```
Out[5]= {xAct`xPert`, xAct`xTensor`, xAct`xPerm`,
         xAct`xCore`, xAct`ExpressionManipulation`, Global`, System`}
```

```
In[6]:= Context[]
```

```
Out[6]= Global`
```

```
In[7]:= ?Global`*
```

```
Information::nomatch : No symbol matching Global`* found. More...
```

---

The default options of `MakeRule` and `ContractMetric` have been changed:

```
In[8]:= Options[MakeRule]
```

```
Out[8]= {PatternIndices → All, TestIndices → True, MetricOn → All,
         UseSymmetries → True, Verbose → False, ContractMetrics → True}
```

```
In[9]:= Options[ContractMetric]
```

```
Out[9]= {AllowUpperDerivatives → True, OverDerivatives → False}
```

---

We also load the timing package and lower its default threshold of 1 second to 0.1 seconds

```
In[10]:= << xAct`ShowTime1`
```

```
In[11]:= $ShowTimeThreshold = 0.1;
```

---

We define a function to collect and canonicalize equal-order terms:

```
In[12]:= org[expr_] := Collect[ContractMetric[expr], $PerturbationParameter, ToCanonical]
```

## ■ 0. Set up

`xPert`` has been built on `xTensor`` and shares its notation. Then, a session with `xPert`` begins just like one with `xTensor``, defining one or more manifolds and vector bundles and the objects living on them:

---

Define a 4d manifold `M`. Note that most of the results will not depend on the chosen dimension.

```
In[13]:= DefManifold[M, 4, {a, b, c, d, e, f}]
```

```
** DefManifold: Defining manifold M.
```

```
** DefVBundle: Defining vbundle TangentM.
```

---

We also define a metric on this manifold, with negative determinant and associated covariant derivative CD. This derivative is extended to act on tensor densities defined using the determinant of the metric  $g$  in the basis  $AIndex$ :

```
In[14]:= DefMetric[-1, g[-a, -b], CD, {";", "D"}, WeightedWithBasis → AIndex]

** DefTensor: Defining symmetric metric tensor g[-a, -b].
** DefTensor: Defining antisymmetric tensor epsilon[a, b, c, d].
** DefCovD: Defining covariant derivative CD[-a].
** DefTensor: Defining vanishing torsion tensor TorsionCD[a, -b, -c].
** DefTensor: Defining symmetric Christoffel tensor ChristoffelCD[a, -b, -c].
** DefTensor: Defining Riemann tensor RiemannCD[-a, -b, -c, -d].
** DefTensor: Defining symmetric Ricci tensor RicciCD[-a, -b].
** DefCovD: Contractions of Riemann automatically replaced by Ricci.
** DefTensor: Defining Ricci scalar RicciScalarCD[].
** DefCovD: Contractions of Ricci automatically replaced by RicciScalar.
** DefTensor: Defining symmetric Einstein tensor EinsteinCD[-a, -b].
** DefTensor: Defining Weyl tensor WeylCD[-a, -b, -c, -d].
Rules {1, 2, 3, 4, 5, 6, 7, 8} have been declared as DownValues for WeylCD.
** DefTensor: Defining symmetric TFRicci tensor TFRicciCD[-a, -b].
Rules {1, 2} have been declared as DownValues for TFRicciCD.
** DefCovD: Computing RiemannToWeylRules for dim 4
** DefCovD: Computing RicciToTFRicci for dim 4
** DefCovD: Computing RicciToEinsteinRules for dim 4
** DefTensor: Defining weight +2 density Detg[]. Determinant.
** DefTensor: Defining tensor Tetrag[-a, -b, -c, -d].
** DefTensor: Defining tensor Tetrag†[-a, -b, -c, -d].

0.788049 Second
```

---

The absolute value of its determinant in the basis  $AIndex$  is obtained as

```
In[15]:= Determinant[g][[]]
```

```
Out[15]= 3
```

---

and it is actually called

```
In[16]:= InputForm[%]
```

```
Out[16]//InputForm=
Detg[]
```

---

By default, curvature tensors are output showing the name of the derivative they are associated to:

```
In[17]:= RiemannCD[-a, -b, -c, d]
```

```
Out[17]= R[D]abcd
```

---

Because we shall work with a single covariant derivative operator, we can choose a simpler output for the geometric tensors

```
In[18]:= PrintAs[RiemannCD] ^= "R";
PrintAs[RicciCD] ^= "R";
PrintAs[RicciScalarCD] ^= "R";
PrintAs[ChristoffelCD] ^= "Γ";
PrintAs[EinsteinCD] ^= "G";
```

---

so that now

```
In[23]:= RiemannCD[-a, -b, -c, d]
```

```
Out[23]= Rabcd
```

After defining the background manifold in `xTensor``, we now define the perturbative structure with the command `DefMetricPerturbation` of `xPert``.

---

We define the symmetric two-tensor `h` that will be understood as the perturbation of the metric. As can be seen in the output, the first index of this tensor is a label index (LI), which will encode the perturbative order.

```
In[24]:= DefMetricPerturbation[g, h, ε]

** DefParameter: Defining parameter ε.
** DefTensor: Defining tensor h[LI[order], -a, -b].
```

---

The perturbative parameter `ε` is stored in the global variable

```
In[25]:= $PerturbationParameter
```

```
Out[25]= ε
```

---

Finally, we colour in blue the perturbative orders

```
In[26]:= Unprotect[IndexForm];
IndexForm[LI[x_]] := ColorString[ToString[x], RGBColor[0, 0, 1]];
Protect[IndexForm];
```

```
In[29]:= IndexForm[LI[n]]
```

```
Out[29]= n
```

## ■ 1. Perturbations of the metric and fundamental properties

The main objective of `xPert`` is giving the  $n$ -th perturbation of any object in terms of metric perturbations of order  $n$  and lower. Here we introduce the `Perturbation` head and give examples of use of different perturbative orders. For those readers already familiar with `xTensor``, let us point out that the symbol `Perturbation` is a linear inert-head.

---

The perturbation of the metric is represented as follows:

```
In[30]:= Perturbation[g[-a, -b]
```

```
Out[30]= h1ab
```

---

`Perturbation` has actually two arguments in general. The first one is the expression we want to perturb and the second one the perturbative order (with default 1 when it is not specified, as in the previous case). For instance, acting on the metric at fifth perturbative order returns

```
In[31]:= Perturbation[g[-a, -b], 5]
```

```
Out[31]= h5ab
```

---

Or on the perturbed tensor itself,

```
In[32]:= Perturbation[h[LI[3], -a, -b], 5]
```

```
Out[32]= h8ab
```

---

It also knows how to act on a Kronecker delta

```
In[33]:= Perturbation[delta[a, -b], 3]
```

```
Out[33]= 0
```

---

`Perturbation` uses the fundamental properties of a derivative: linearity and Leibnitz rule. When it does not know the answer of the action on a given tensor, for example the inverse of the metric or the Ricci tensor, the head `Perturbation` is kept (this is where its inert-head character plays a role):

```
In[34]:= Perturbation[g[-a, -b] + g[c, d] RicciCD[-a, -c] RicciCD[-d, -b]]
```

```
Out[34]= h1ab + gcd Δ[Rdb] Rac + gcd Δ[Rac] Rdb + Δ[gcd] Rac Rdb
```

---

We can consider the same case but at a higher perturbative order, for instance at seventeenth order and see that the calculation is almost immediate.

```
In[35]:= Perturbation[g[-a, -b] + g[c, d] RicciCD[-a, -c] RicciCD[-d, -b], 17]
```

```
Out[35]= h17ab + 272 Δ[Rac] Δ[Rdb] Δ15[gcd] + 2040 Δ[Rdb] Δ14[gcd] Δ2[Rac] +
9520 Δ[Rdb] Δ13[gcd] Δ3[Rac] + 30940 Δ[Rdb] Δ12[gcd] Δ4[Rac] +
74256 Δ[Rdb] Δ11[gcd] Δ5[Rac] + 136136 Δ[Rdb] Δ10[gcd] Δ6[Rac] +
194480 Δ[Rdb] Δ9[gcd] Δ7[Rac] + 218790 Δ[Rdb] Δ8[gcd] Δ8[Rac] +
194480 Δ[Rdb] Δ7[gcd] Δ9[Rac] + 136136 Δ[Rdb] Δ6[gcd] Δ10[Rac] +
74256 Δ[Rdb] Δ5[gcd] Δ11[Rac] + 30940 Δ[Rdb] Δ4[gcd] Δ12[Rac] +
9520 Δ[Rdb] Δ3[gcd] Δ13[Rac] + 2040 Δ[Rdb] Δ2[gcd] Δ14[Rac] + 272 Δ[gcd] Δ[Rdb] Δ15[Rac] +
17 gcd Δ[Rdb] Δ16[Rac] + 2040 Δ[Rac] Δ14[gcd] Δ2[Rdb] + 14280 Δ13[gcd] Δ2[Rac] Δ2[Rdb] +
61880 Δ12[gcd] Δ3[Rac] Δ2[Rdb] + 185640 Δ11[gcd] Δ4[Rac] Δ2[Rdb] +
408408 Δ10[gcd] Δ5[Rac] Δ2[Rdb] + 680680 Δ9[gcd] Δ6[Rac] Δ2[Rdb] +
875160 Δ8[gcd] Δ7[Rac] Δ2[Rdb] + 875160 Δ7[gcd] Δ8[Rac] Δ2[Rdb] +
680680 Δ6[gcd] Δ9[Rac] Δ2[Rdb] + 408408 Δ5[gcd] Δ10[Rac] Δ2[Rdb] +
185640 Δ4[gcd] Δ11[Rac] Δ2[Rdb] + 61880 Δ3[gcd] Δ12[Rac] Δ2[Rdb] +
14280 Δ2[gcd] Δ13[Rac] Δ2[Rdb] + 2040 Δ[gcd] Δ14[Rac] Δ2[Rdb] +
```

$$\begin{aligned}
 &136 g^{cd} \Delta^{15} [R_{ac}] \Delta^2 [R_{db}] + 9520 \Delta [R_{ac}] \Delta^{13} [g^{cd}] \Delta^3 [R_{db}] + 61880 \Delta^{12} [g^{cd}] \Delta^2 [R_{ac}] \Delta^3 [R_{db}] + \\
 &247520 \Delta^{11} [g^{cd}] \Delta^3 [R_{ac}] \Delta^3 [R_{db}] + 680680 \Delta^{10} [g^{cd}] \Delta^4 [R_{ac}] \Delta^3 [R_{db}] + \\
 &1361360 \Delta^9 [g^{cd}] \Delta^5 [R_{ac}] \Delta^3 [R_{db}] + 2042040 \Delta^8 [g^{cd}] \Delta^6 [R_{ac}] \Delta^3 [R_{db}] + \\
 &2333760 \Delta^7 [g^{cd}] \Delta^7 [R_{ac}] \Delta^3 [R_{db}] + 2042040 \Delta^6 [g^{cd}] \Delta^8 [R_{ac}] \Delta^3 [R_{db}] + \\
 &1361360 \Delta^5 [g^{cd}] \Delta^9 [R_{ac}] \Delta^3 [R_{db}] + 680680 \Delta^4 [g^{cd}] \Delta^{10} [R_{ac}] \Delta^3 [R_{db}] + \\
 &247520 \Delta^3 [g^{cd}] \Delta^{11} [R_{ac}] \Delta^3 [R_{db}] + 61880 \Delta^2 [g^{cd}] \Delta^{12} [R_{ac}] \Delta^3 [R_{db}] + \\
 &9520 \Delta [g^{cd}] \Delta^{13} [R_{ac}] \Delta^3 [R_{db}] + 680 g^{cd} \Delta^{14} [R_{ac}] \Delta^3 [R_{db}] + 30940 \Delta [R_{ac}] \Delta^{12} [g^{cd}] \Delta^4 [R_{db}] + \\
 &185640 \Delta^{11} [g^{cd}] \Delta^2 [R_{ac}] \Delta^4 [R_{db}] + 680680 \Delta^{10} [g^{cd}] \Delta^3 [R_{ac}] \Delta^4 [R_{db}] + \\
 &1701700 \Delta^9 [g^{cd}] \Delta^4 [R_{ac}] \Delta^4 [R_{db}] + 3063060 \Delta^8 [g^{cd}] \Delta^5 [R_{ac}] \Delta^4 [R_{db}] + \\
 &4084080 \Delta^7 [g^{cd}] \Delta^6 [R_{ac}] \Delta^4 [R_{db}] + 4084080 \Delta^6 [g^{cd}] \Delta^7 [R_{ac}] \Delta^4 [R_{db}] + \\
 &3063060 \Delta^5 [g^{cd}] \Delta^8 [R_{ac}] \Delta^4 [R_{db}] + 1701700 \Delta^4 [g^{cd}] \Delta^9 [R_{ac}] \Delta^4 [R_{db}] + \\
 &680680 \Delta^3 [g^{cd}] \Delta^{10} [R_{ac}] \Delta^4 [R_{db}] + 185640 \Delta^2 [g^{cd}] \Delta^{11} [R_{ac}] \Delta^4 [R_{db}] + \\
 &30940 \Delta [g^{cd}] \Delta^{12} [R_{ac}] \Delta^4 [R_{db}] + 2380 g^{cd} \Delta^{13} [R_{ac}] \Delta^4 [R_{db}] + 74256 \Delta [R_{ac}] \Delta^{11} [g^{cd}] \Delta^5 [R_{db}] + \\
 &408408 \Delta^{10} [g^{cd}] \Delta^2 [R_{ac}] \Delta^5 [R_{db}] + 1361360 \Delta^9 [g^{cd}] \Delta^3 [R_{ac}] \Delta^5 [R_{db}] + \\
 &3063060 \Delta^8 [g^{cd}] \Delta^4 [R_{ac}] \Delta^5 [R_{db}] + 4900896 \Delta^7 [g^{cd}] \Delta^5 [R_{ac}] \Delta^5 [R_{db}] + \\
 &5717712 \Delta^6 [g^{cd}] \Delta^6 [R_{ac}] \Delta^5 [R_{db}] + 4900896 \Delta^5 [g^{cd}] \Delta^7 [R_{ac}] \Delta^5 [R_{db}] + \\
 &3063060 \Delta^4 [g^{cd}] \Delta^8 [R_{ac}] \Delta^5 [R_{db}] + 1361360 \Delta^3 [g^{cd}] \Delta^9 [R_{ac}] \Delta^5 [R_{db}] + \\
 &408408 \Delta^2 [g^{cd}] \Delta^{10} [R_{ac}] \Delta^5 [R_{db}] + 74256 \Delta [g^{cd}] \Delta^{11} [R_{ac}] \Delta^5 [R_{db}] + 6188 g^{cd} \Delta^{12} [R_{ac}] \Delta^5 [R_{db}] + \\
 &136136 \Delta [R_{ac}] \Delta^{10} [g^{cd}] \Delta^6 [R_{db}] + 680680 \Delta^9 [g^{cd}] \Delta^2 [R_{ac}] \Delta^6 [R_{db}] + \\
 &2042040 \Delta^8 [g^{cd}] \Delta^3 [R_{ac}] \Delta^6 [R_{db}] + 4084080 \Delta^7 [g^{cd}] \Delta^4 [R_{ac}] \Delta^6 [R_{db}] + \\
 &5717712 \Delta^6 [g^{cd}] \Delta^5 [R_{ac}] \Delta^6 [R_{db}] + 5717712 \Delta^5 [g^{cd}] \Delta^6 [R_{ac}] \Delta^6 [R_{db}] + \\
 &4084080 \Delta^4 [g^{cd}] \Delta^7 [R_{ac}] \Delta^6 [R_{db}] + 2042040 \Delta^3 [g^{cd}] \Delta^8 [R_{ac}] \Delta^6 [R_{db}] + \\
 &680680 \Delta^2 [g^{cd}] \Delta^9 [R_{ac}] \Delta^6 [R_{db}] + 136136 \Delta [g^{cd}] \Delta^{10} [R_{ac}] \Delta^6 [R_{db}] + \\
 &12376 g^{cd} \Delta^{11} [R_{ac}] \Delta^6 [R_{db}] + 194480 \Delta [R_{ac}] \Delta^9 [g^{cd}] \Delta^7 [R_{db}] + 875160 \Delta^8 [g^{cd}] \Delta^2 [R_{ac}] \Delta^7 [R_{db}] + \\
 &2333760 \Delta^7 [g^{cd}] \Delta^3 [R_{ac}] \Delta^7 [R_{db}] + 4084080 \Delta^6 [g^{cd}] \Delta^4 [R_{ac}] \Delta^7 [R_{db}] + \\
 &4900896 \Delta^5 [g^{cd}] \Delta^5 [R_{ac}] \Delta^7 [R_{db}] + 4084080 \Delta^4 [g^{cd}] \Delta^6 [R_{ac}] \Delta^7 [R_{db}] + \\
 &2333760 \Delta^3 [g^{cd}] \Delta^7 [R_{ac}] \Delta^7 [R_{db}] + 875160 \Delta^2 [g^{cd}] \Delta^8 [R_{ac}] \Delta^7 [R_{db}] + \\
 &194480 \Delta [g^{cd}] \Delta^9 [R_{ac}] \Delta^7 [R_{db}] + 19448 g^{cd} \Delta^{10} [R_{ac}] \Delta^7 [R_{db}] + 218790 \Delta [R_{ac}] \Delta^8 [g^{cd}] \Delta^8 [R_{db}] + \\
 &875160 \Delta^7 [g^{cd}] \Delta^2 [R_{ac}] \Delta^8 [R_{db}] + 2042040 \Delta^6 [g^{cd}] \Delta^3 [R_{ac}] \Delta^8 [R_{db}] + \\
 &3063060 \Delta^5 [g^{cd}] \Delta^4 [R_{ac}] \Delta^8 [R_{db}] + 3063060 \Delta^4 [g^{cd}] \Delta^5 [R_{ac}] \Delta^8 [R_{db}] + \\
 &2042040 \Delta^3 [g^{cd}] \Delta^6 [R_{ac}] \Delta^8 [R_{db}] + 875160 \Delta^2 [g^{cd}] \Delta^7 [R_{ac}] \Delta^8 [R_{db}] + \\
 &218790 \Delta [g^{cd}] \Delta^8 [R_{ac}] \Delta^8 [R_{db}] + 24310 g^{cd} \Delta^9 [R_{ac}] \Delta^8 [R_{db}] + 194480 \Delta [R_{ac}] \Delta^7 [g^{cd}] \Delta^9 [R_{db}] + \\
 &680680 \Delta^6 [g^{cd}] \Delta^2 [R_{ac}] \Delta^9 [R_{db}] + 1361360 \Delta^5 [g^{cd}] \Delta^3 [R_{ac}] \Delta^9 [R_{db}] + \\
 &1701700 \Delta^4 [g^{cd}] \Delta^4 [R_{ac}] \Delta^9 [R_{db}] + 1361360 \Delta^3 [g^{cd}] \Delta^5 [R_{ac}] \Delta^9 [R_{db}] + \\
 &680680 \Delta^2 [g^{cd}] \Delta^6 [R_{ac}] \Delta^9 [R_{db}] + 194480 \Delta [g^{cd}] \Delta^7 [R_{ac}] \Delta^9 [R_{db}] + 24310 g^{cd} \Delta^8 [R_{ac}] \Delta^9 [R_{db}] + \\
 &136136 \Delta [R_{ac}] \Delta^6 [g^{cd}] \Delta^{10} [R_{db}] + 408408 \Delta^5 [g^{cd}] \Delta^2 [R_{ac}] \Delta^{10} [R_{db}] + \\
 &680680 \Delta^4 [g^{cd}] \Delta^3 [R_{ac}] \Delta^{10} [R_{db}] + 680680 \Delta^3 [g^{cd}] \Delta^4 [R_{ac}] \Delta^{10} [R_{db}] + \\
 &408408 \Delta^2 [g^{cd}] \Delta^5 [R_{ac}] \Delta^{10} [R_{db}] + 136136 \Delta [g^{cd}] \Delta^6 [R_{ac}] \Delta^{10} [R_{db}] + \\
 &19448 g^{cd} \Delta^7 [R_{ac}] \Delta^{10} [R_{db}] + 74256 \Delta [R_{ac}] \Delta^5 [g^{cd}] \Delta^{11} [R_{db}] + \\
 &185640 \Delta^4 [g^{cd}] \Delta^2 [R_{ac}] \Delta^{11} [R_{db}] + 247520 \Delta^3 [g^{cd}] \Delta^3 [R_{ac}] \Delta^{11} [R_{db}] + \\
 &185640 \Delta^2 [g^{cd}] \Delta^4 [R_{ac}] \Delta^{11} [R_{db}] + 74256 \Delta [g^{cd}] \Delta^5 [R_{ac}] \Delta^{11} [R_{db}] + \\
 &12376 g^{cd} \Delta^6 [R_{ac}] \Delta^{11} [R_{db}] + 30940 \Delta [R_{ac}] \Delta^4 [g^{cd}] \Delta^{12} [R_{db}] + 61880 \Delta^3 [g^{cd}] \Delta^2 [R_{ac}] \Delta^{12} [R_{db}] + \\
 &61880 \Delta^2 [g^{cd}] \Delta^3 [R_{ac}] \Delta^{12} [R_{db}] + 30940 \Delta [g^{cd}] \Delta^4 [R_{ac}] \Delta^{12} [R_{db}] + \\
 &6188 g^{cd} \Delta^5 [R_{ac}] \Delta^{12} [R_{db}] + 9520 \Delta [R_{ac}] \Delta^3 [g^{cd}] \Delta^{13} [R_{db}] + 14280 \Delta^2 [g^{cd}] \Delta^2 [R_{ac}] \Delta^{13} [R_{db}] + \\
 &9520 \Delta [g^{cd}] \Delta^3 [R_{ac}] \Delta^{13} [R_{db}] + 2380 g^{cd} \Delta^4 [R_{ac}] \Delta^{13} [R_{db}] + 2040 \Delta [R_{ac}] \Delta^2 [g^{cd}] \Delta^{14} [R_{db}] + \\
 &2040 \Delta [g^{cd}] \Delta^2 [R_{ac}] \Delta^{14} [R_{db}] + 680 g^{cd} \Delta^3 [R_{ac}] \Delta^{14} [R_{db}] + 272 \Delta [g^{cd}] \Delta [R_{ac}] \Delta^{15} [R_{db}] + \\
 &136 g^{cd} \Delta^2 [R_{ac}] \Delta^{15} [R_{db}] + 17 g^{cd} \Delta [R_{ac}] \Delta^{16} [R_{db}] + 17 \Delta [R_{db}] \Delta^{16} [g^{cd}] R_{ac} + \\
 &136 \Delta^{15} [g^{cd}] \Delta^2 [R_{db}] R_{ac} + 680 \Delta^{14} [g^{cd}] \Delta^3 [R_{db}] R_{ac} + 2380 \Delta^{13} [g^{cd}] \Delta^4 [R_{db}] R_{ac} + \\
 &6188 \Delta^{12} [g^{cd}] \Delta^5 [R_{db}] R_{ac} + 12376 \Delta^{11} [g^{cd}] \Delta^6 [R_{db}] R_{ac} + 19448 \Delta^{10} [g^{cd}] \Delta^7 [R_{db}] R_{ac} + \\
 &24310 \Delta^9 [g^{cd}] \Delta^8 [R_{db}] R_{ac} + 24310 \Delta^8 [g^{cd}] \Delta^9 [R_{db}] R_{ac} + 19448 \Delta^7 [g^{cd}] \Delta^{10} [R_{db}] R_{ac} + \\
 &12376 \Delta^6 [g^{cd}] \Delta^{11} [R_{db}] R_{ac} + 6188 \Delta^5 [g^{cd}] \Delta^{12} [R_{db}] R_{ac} + 2380 \Delta^4 [g^{cd}] \Delta^{13} [R_{db}] R_{ac} + \\
 &680 \Delta^3 [g^{cd}] \Delta^{14} [R_{db}] R_{ac} + 136 \Delta^2 [g^{cd}] \Delta^{15} [R_{db}] R_{ac} + 17 \Delta [g^{cd}] \Delta^{16} [R_{db}] R_{ac} + \\
 &g^{cd} \Delta^{17} [R_{db}] R_{ac} + 17 \Delta [R_{ac}] \Delta^{16} [g^{cd}] R_{db} + 136 \Delta^{15} [g^{cd}] \Delta^2 [R_{ac}] R_{db} +
 \end{aligned}$$

$$\begin{aligned}
& 680 \Delta^{14} [g^{cd}] \Delta^3 [R_{ac}] R_{db} + 2380 \Delta^{13} [g^{cd}] \Delta^4 [R_{ac}] R_{db} + 6188 \Delta^{12} [g^{cd}] \Delta^5 [R_{ac}] R_{db} + \\
& 12376 \Delta^{11} [g^{cd}] \Delta^6 [R_{ac}] R_{db} + 19448 \Delta^{10} [g^{cd}] \Delta^7 [R_{ac}] R_{db} + 24310 \Delta^9 [g^{cd}] \Delta^8 [R_{ac}] R_{db} + \\
& 24310 \Delta^8 [g^{cd}] \Delta^9 [R_{ac}] R_{db} + 19448 \Delta^7 [g^{cd}] \Delta^{10} [R_{ac}] R_{db} + 12376 \Delta^6 [g^{cd}] \Delta^{11} [R_{ac}] R_{db} + \\
& 6188 \Delta^5 [g^{cd}] \Delta^{12} [R_{ac}] R_{db} + 2380 \Delta^4 [g^{cd}] \Delta^{13} [R_{ac}] R_{db} + 680 \Delta^3 [g^{cd}] \Delta^{14} [R_{ac}] R_{db} + \\
& 136 \Delta^2 [g^{cd}] \Delta^{15} [R_{ac}] R_{db} + 17 \Delta [g^{cd}] \Delta^{16} [R_{ac}] R_{db} + g^{cd} \Delta^{17} [R_{ac}] R_{db} + \Delta^{17} [g^{cd}] R_{ac} R_{db}
\end{aligned}$$

---

The perturbative order of an expression can be computed with

```
In[36]:= PerturbationOrder[%]
```

```
Out[36]= 17
```

---

By definition, Perturbation commutes with partial derivatives,

```
In[37]:= Perturbation[PD[-a][RicciCD[-c, -d]], 5]
```

```
Out[37]= ∂a Δ5 [Rcd]
```

---

and with covariant derivatives acting on scalars,

```
In[38]:= Perturbation[CD[-a][RicciScalarCD[]], 6]
```

```
Out[38]= Da Δ6 [R]
```

---

but not with covariant derivatives acting on tensors,

```
In[39]:= Perturbation[CD[-a][RicciCD[a, -b]], 3]
```

```
Out[39]= Δ3 [Rab;a]
```

It is essential to have in mind that the ToCanonical or Simplification commands cannot be used as long as there is any Perturbation head in the expression because internal indices of the expression inside Perturbation could be moved improperly. A warning reminds us of the danger:

---

For example the expression

```
In[40]:= h[LI[3], -a, -e] CD[-d][Perturbation[RicciCD[a, e]]]
```

```
Out[40]= h3ae (Dd Δ[Rae])
```

---

is not properly canonicalized here:

```
In[41]:= % // ToCanonical
```

```
ToCanonical::canonpert :
ToCanonical on Perturbation may incorrectly change index characters.
```

```
Out[41]= h3ae (Dd Δ[Rae])
```



This is not a restriction at all. It just means that the perturbations must be replaced by other tensor expression before canonicalization. How this is done is shown next. In case canonicalization at this stage is really needed the right sequence of commands is

```
In[42]:= ToCanonical[% // SeparateMetric[] // ContractMetric, UseMetricOnVBundle → None]
```

```
Out[42]= 2 h3ae Rab (Dd Δ[geb]) + h3ae (Dd Δ[Rae]) + 2 h3ae Δ[gab] Reb;d
```

which uses explicit metric factors instead of moving indices up and down.

## ■ 2. ExpandPerturbation

The command `ExpandPerturbation` is the most powerful one of this package. As we have seen, the `Perturbation` operator automatically uses linearity and the Leibnitz rule to expand expressions. But up to now it only knows how to act on the metric with covariant indices. All other cases are handled by `ExpandPerturbation`, which makes use of exact formulas to expand the perturbations of any geometric object.

The perturbations of the inverse metric are

```
In[43]:= Perturbation[g[a, b]]
```

```
Out[43]= Δ[gab]
```

```
In[44]:= % // ExpandPerturbation
```

```
Out[44]= -h1ab
```

```
In[45]:= Perturbation[g[a, b], 4]
```

```
Out[45]= Δ4[gab]
```

```
In[46]:= % // ExpandPerturbation
```

```
Out[46]= 24 h1ac h1cd h1de h1eb - 12 h1cd h1db h2ac + 6 h2ac h2cb -
12 h1ac h1db h2cd - 12 h1ac h1cd h2db + 4 h1cb h3ac + 4 h1ac h3cb - h4ab
```

or the perturbations of the Ricci tensor are

```
In[47]:= Perturbation[RicciCD[-a, -b]] // ExpandPerturbation
```

```
Out[47]= 1/2 (-h1cc;b;a - h1cb;c;a + h1bcic;a) + 1/2 (h1cb;a;c + h1ca;b;c - h1baic;c)
```

```
In[48]:= Perturbation[RicciCD[-a, -b], 2] // ExpandPerturbation
```

```
Out[48]= 1/2 (-h2cc;b;a - h2cb;c;a + h2bcic;a) + 1/2 (h2cb;a;c + h2ca;b;c - h2baic;c) +
2 (1/2 h1cd (h1dcb;a + h1dbc;a - h1bcd;a) + 1/4 (h1ecb + h1ebc - h1cbe) (h1eca + h1eaic - h1caie)) -
2 (1/2 h1cd (h1dba;c + h1dab;c - h1bad;c) + 1/4 (h1eba + h1eab - h1abe) (h1ecc + h1ecic - h1ccie))
```

The `ExpandPerturbation` command has been only predefined on certain tensors, and only when those have indices with their expected characters (for the Ricci tensor, for example, both indices are expected to be covariant; all characters are accepted for the metric). All other cases and derivatives of the tensors are handled through optional internal algorithms.

---

`ExpandPerturbation` has two options, which are by default switched on:

```
In[49]:= Options[ExpandPerturbation]
```

```
Out[49]= {SeparateMetric → True, OverDerivatives → True}
```

---

`SeparateMetric` moves indices into their character (covariant or contravariant) specified at definition time, introducing metric factors if required (this is performed using the `SeparateMetric` command in `xTensor``, hence the name of the option). For instance:

```
In[50]:= Perturbation[RicciCD[a, b]]
```

```
Out[50]= Δ[Rab]
```

```
In[51]:= % // ExpandPerturbation
```

```
Out[51]= -gbd h1ac Rcd - gac h1bd Rcd +
gac gbd ( 1/2 (-h1ee;d;c - h1ed;e;c + h1ede;c) + 1/2 (h1ed;c;e + h1ec;d;e - h1edc;e) )
```

---

The additional metric factors can be absorbed using `ContractMetric`:

```
In[52]:= % // ContractMetric
```

```
Out[52]= -h1bc Rac - h1ac Rbc - h1cc;b;a - h1cbic;a + h1bc;c;a + h1cb;aic + h1ca;bic - h1ba;cic
```

```
In[53]:= ExpandPerturbation[%%, SeparateMetric → False]
```

```
Out[53]= Δ[Rab]
```

---

`OverDerivatives` allows to jump over Lie or covariant derivatives. The latter gives perturbations of the Christoffel symbols, that are also expanded into derivatives of metric perturbations. For example,

```
In[54]:= Perturbation[CD[-a][RicciCD[-b, -c]]]
```

```
Out[54]= Δ[Rbc;a]
```

```
In[55]:= % // ExpandPerturbation
```

```
Out[55]= 1/2 (-h1dd;c;b;a - h1dc;d;b;a + h1cdid;b;a) + 1/2 (h1dc;b;d;a + h1db;c;d;a - h1cbid;a) -
1/2 Rdc (h1db;a + h1da;b - h1dab) - 1/2 Rbd (h1dc;a + h1da;c - h1dac)
```

```
In[56]:= ExpandPerturbation[%%, OverDerivatives → False]
```

```
Out[56]= Δ[Rbc;a]
```

---

In order to show an example with Lie derivatives, we need a vector-field

```
In[57]:= DefTensor[v[a], M]
      ** DefTensor: Defining tensor v[a].
```

---

We perturb now at fourth-order the Lie derivative of the vector field over itself

```
In[58]:= Perturbation[LieD[v[a]][v[b]], 4]
Out[58]=  $\Delta^4[\mathcal{L}_v v^b]$ 
```

---

Note that the perturbation of a Lie derivative can be written in terms of Lie derivatives only

```
In[59]:= % // ExpandPerturbation
Out[59]= 4 ( $\mathcal{L}_{\Delta[v^a]} \Delta^3[v^b]$ ) + 6 ( $\mathcal{L}_{\Delta^2[v^a]} \Delta^2[v^b]$ ) + 4 ( $\mathcal{L}_{\Delta^3[v^a]} \Delta[v^b]$ ) +  $\mathcal{L}_{\Delta^4[v^a]} v^b$  +  $\mathcal{L}_v \Delta^4[v^b]$ 
```

---

Again, nothing happens when switching-off the option OverDerivatives

```
In[60]:= ExpandPerturbation[%%, OverDerivatives -> False]
Out[60]=  $\Delta^4[\mathcal{L}_v v^b]$ 
```

## ■ 3. Perturbations of curvature tensors

In this section we show how `ExpandPerturbation` acts on the most important geometric tensors of a Riemannian manifold, at any perturbative order. We shall compare its direct action at a particular order with a recursive computation of that order from lower orders. This is both a consistency check and an example of the efficiency of the direct action.

### 3.1. Perturbations of the inverse metric

Let us start comparing the seventh-order recursive and direct perturbations of the inverse of our metric tensor field.

---

Recursively :

```
In[61]:= Perturbation[g[a, b]] // ExpandPerturbation
Out[61]=  $-h^{1ab}$ 

In[62]:= % // Perturbation // ExpandPerturbation
Out[62]=  $g^{bd} h^{1ac} h^1_{cd} + g^{ac} h^{1bd} h^1_{cd} - g^{ac} g^{bd} h^2_{cd}$ 

In[63]:= % // Perturbation // ExpandPerturbation // org
Out[63]=  $-6 h^{1ac} h^{1bd} h^1_{cd} + 3 h^{1b}_c h^{2ac} + 3 h^{1ac} h^{2b}_c - h^{3ab}$ 
```

```
In[64]:= % // Perturbation // ExpandPerturbation // org
```

0.204013 Second

```
Out[64]= 24 h1ac h1bd h1c e h1de - 12 h1bd h1cd h2ac + 6 h2ac h2bc -
12 h1ac h1cd h2bd - 12 h1ac h1bd h2cd + 4 h1bc h3ac + 4 h1ac h3bc - h4ab
```

```
In[65]:= % // Perturbation // ExpandPerturbation // org
```

0.488031 Second

```
Out[65]= -120 h1ac h1bd h1c e h1df h1ef + 60 h1bd h1c e h1de h2ac + 60 h1ac h1c e h1de h2bd -
30 h1cd h2ac h2bd - 30 h1bd h2ac h2cd - 30 h1ac h2bd h2cd + 60 h1ac h1bd h1d e h2ce +
60 h1ac h1bd h1c e h2de - 20 h1bd h1cd h3ac + 10 h2bc h3ac + 10 h2ac h3bc -
20 h1ac h1cd h3bd - 20 h1ac h1bd h3cd + 5 h1bc h4ac + 5 h1ac h4bc - h5ab
```

```
In[66]:= % // Perturbation // ExpandPerturbation // org
```

1.28408 Second

```
Out[66]= 720 h1ac h1bd h1c e h1df h1ef h1ff1 - 360 h1bd h1c e h1df h1ef h2ac - 360 h1ac h1c e h1df h1ef h2bd +
180 h1c e h1de h2ac h2bd - 90 h2ac h2bd h2cd + 180 h1bd h1d e h2ac h2ce + 180 h1ac h1d e h2bd h2ce -
360 h1ac h1bd h1d e h1ef h2cf + 180 h1bd h1c e h2ac h2de + 180 h1ac h1c e h2bd h2de +
180 h1ac h1bd h2c e h2de - 360 h1ac h1bd h1c e h1ef h2df - 360 h1ac h1bd h1c e h1df h2ef +
120 h1bd h1c e h1de h3ac - 60 h1cd h2bd h3ac - 60 h1bd h2cd h3ac + 20 h3ac h3bc +
120 h1ac h1c e h1de h3bd - 60 h1cd h2ac h3bd - 60 h1ac h2cd h3bd - 60 h1bd h2ac h3cd -
60 h1ac h2bd h3cd + 120 h1ac h1bd h1d e h3ce + 120 h1ac h1bd h1c e h3de - 30 h1bd h1cd h4ac +
15 h2bc h4ac + 15 h2ac h4bc - 30 h1ac h1cd h4bd - 30 h1ac h1bd h4cd + 6 h1bc h5ac + 6 h1ac h5bc - h6ab
```

```
In[67]:= % // Perturbation // ExpandPerturbation // org
```

4.05625 Second

```
Out[67]= -5040 h1ac h1bd h1c e h1df h1ef h1ff1 h1ff2 h1ff1f2 + 2520 h1bd h1c e h1df h1ef h1ff1 h2ac +
2520 h1ac h1c e h1df h1ef h1ff1 h2bd - 1260 h1c e h1df h1ef h2ac h2bd + 630 h1d e h2ac h2bd h2ce -
1260 h1bd h1d e h1ef h2ac h2cf - 1260 h1ac h1d e h1ef h2bd h2cf + 2520 h1ac h1bd h1d e h1ef h1ff1 h2cf1 +
630 h1c e h2ac h2bd h2de + 630 h1bd h2ac h2c e h2de + 630 h1ac h2bd h2c e h2de -
1260 h1bd h1c e h1ef h2df - 1260 h1ac h1c e h1ef h2bd h2df - 1260 h1ac h1bd h1ef h2ce h2df +
2520 h1ac h1bd h1c e h1ef h1ff1 h2df1 - 1260 h1bd h1c e h1df h2ac h2ef - 1260 h1ac h1d e h1ef h2bd h2ef -
1260 h1ac h1bd h1d e h2cf h2ef - 1260 h1ac h1bd h1c e h2df h2ef + 2520 h1ac h1bd h1c e h1df h1ff1 h2ef1 +
2520 h1ac h1bd h1c e h1df h1ff1 h2ff1 - 840 h1bd h1c e h1df h1ef h3ac + 420 h1c e h1de h2bd h3ac -
210 h2bd h2cd h3ac + 420 h1bd h1d e h2ce h3ac + 420 h1bd h1c e h2de h3ac -
840 h1ac h1c e h1df h1ef h3bd + 420 h1c e h1de h2ac h3bd - 210 h2ac h2cd h3bd +
420 h1ac h1d e h2ce h3bd + 420 h1ac h1c e h2de h3bd - 140 h1cd h3ac h3bd - 210 h2ac h2bd h3cd -
140 h1bd h3ac h3cd - 140 h1ac h3bd h3cd + 420 h1bd h1d e h2ac h3ce + 420 h1ac h1d e h2bd h3ce +
420 h1ac h1bd h2d e h3ce - 840 h1ac h1bd h1d e h1ef h3cf + 420 h1bd h1c e h2ac h3de +
420 h1ac h1c e h2bd h3de + 420 h1ac h1bd h2c e h3de - 840 h1ac h1bd h1c e h1ef h3df -
840 h1ac h1bd h1c e h1df h3ef + 210 h1bd h1c e h1de h4ac - 105 h1cd h2bd h4ac -
105 h1bd h2cd h4ac + 35 h3bc h4ac + 35 h3ac h4bc + 210 h1ac h1c e h1de h4bd -
105 h1cd h2ac h4bd - 105 h1ac h2cd h4bd - 105 h1bd h2ac h4cd - 105 h1ac h2bd h4cd +
210 h1ac h1bd h1d e h4ce + 210 h1ac h1bd h1c e h4de - 42 h1bd h1cd h5ac + 21 h2bc h5ac +
21 h2ac h5bc - 42 h1ac h1cd h5bd - 42 h1ac h1bd h5cd + 7 h1bc h6ac + 7 h1ac h6bc - h7ab
```

Applying ExpandPerturbation directly,

```
In[68]:= Perturbation[g[a, b], 7] // ExpandPerturbation
```

```
Out[68]= -5040 h1ac h1cd h1de h1ef h1f1 h1f1 h1f2 h1fb + 2520 h1cd h1de h1ef h1f1 h1f1 h1fb h2ac +
2520 h1ac h1de h1ef h1f1 h1fb h2cd - 1260 h1de h1ef h1fb h2ac h2cd +
2520 h1ac h1cd h1de h1ef h1f1 h1fb h2de - 1260 h1cd h1de h1ef h1fb h2ac h2de - 1260 h1ac h1ef h1fb h2cd h2de +
630 h1eb h2ac h2cd h2de + 630 h1de h2ac h2cd h2eb + 630 h1cd h2ac h2de h2eb + 630 h1ac h2cd h2de h2eb +
2520 h1ac h1cd h1de h1ef h1f1 h1fb h2f - 1260 h1cd h1de h1ef h1fb h2ac h2f - 1260 h1ac h1de h1fb h2cd h2f -
1260 h1ac h1cd h1fb h2de h2f - 1260 h1cd h1de h1ef h2ac h2fb - 1260 h1ac h1de h1ef h2cd h2fb -
1260 h1ac h1cd h1ef h2de h2fb - 1260 h1ac h1cd h1de h2ef h2fb + 2520 h1ac h1cd h1de h1ef h1fb h2f1 +
2520 h1ac h1cd h1de h1ef h1f1 h2fb - 840 h1cd h1de h1ef h1fb h3ac + 420 h1de h1fb h2cd h3ac -
210 h2cd h2fb h3ac + 420 h1cd h1fb h2de h3ac + 420 h1cd h1de h2eb h3ac - 840 h1ac h1de h1fb h3cd +
420 h1de h1fb h2ac h3cd - 210 h2ac h2fb h3cd + 420 h1ac h1fb h2de h3cd + 420 h1ac h1de h2eb h3cd -
140 h1db h3ac h3cd - 210 h2ac h2fb h3cd - 140 h1cd h3ac h3db - 140 h1ac h3cd h3db -
840 h1ac h1cd h1fb h1de h3e + 420 h1cd h1fb h2ac h3e + 420 h1ac h1fb h2cd h3e +
420 h1ac h1cd h2eb h3e + 420 h1cd h1de h2ac h3eb + 420 h1ac h1de h2cd h3eb + 420 h1ac h1cd h2de h3eb -
840 h1ac h1cd h1de h1fb h3f - 840 h1ac h1cd h1de h1fb h3fb + 210 h1cd h1de h1fb h4ac -
105 h1db h2cd h4ac - 105 h1cd h2fb h4ac + 35 h3cb h4ac + 35 h3ac h4cb + 210 h1ac h1de h1fb h4cd -
105 h1db h2ac h4cd - 105 h1ac h2fb h4cd - 105 h1cd h2ac h4db - 105 h1ac h2cd h4db +
210 h1ac h1cd h1de h4e + 210 h1ac h1cd h1de h4eb - 42 h1cd h1fb h5ac + 21 h2cb h5ac +
21 h2ac h5cb - 42 h1ac h1db h5cd - 42 h1ac h1cd h5db + 7 h1cb h6ac + 7 h1ac h6cb - h7ab
```

Comparison and check:

```
In[69]:= % - %% // ToCanonical
```

```
1.5481 Second
```

```
Out[69]= 0
```

Another interesting check is proving that the perturbed inverse of the metric is in fact the inverse of the perturbed metric up to the correct order. For that we will use another xPert ` command named Perturbed that expands an expression into a power series up to the order provided.

The power series of the perturbed metric up to fifth order (note that we do not use the Series construct of Mathematica)

```
In[70]:= Perturbed[g[-a, -b], 5]
```

```
Out[70]= gab + ε h1ab +  $\frac{1}{2}$  ε2 h2ab +  $\frac{1}{6}$  ε3 h3ab +  $\frac{1}{24}$  ε4 h4ab +  $\frac{1}{120}$  ε5 h5ab
```

The corresponding expansion for its inverse

```
In[71]:= Perturbed[g[b, c], 5] // ExpandPerturbation
```

$$\begin{aligned} \text{Out}[71]= & g^{bc} - \epsilon h^{1bc} + \frac{1}{2} \epsilon^2 (2 h_a^1 c h^{1ba} - h^{2bc}) + \\ & \frac{1}{6} \epsilon^3 (-6 h_a^1 d h^{1ba} h_d^1 c + 3 h_f^1 c h^{2bf} + 3 h^{1be} h_e^2 c - h^{3bc}) + \\ & \frac{1}{24} \epsilon^4 (24 h_a^1 d h^{1ba} h_d^1 e h_e^1 c - 12 h_{f4}^1 h_{f5}^1 c h^{2bf4} - 12 h^{1bf} h_f^{f1} h^2_{f1} c - \\ & 12 h^{1bf2} h_{f3}^1 c h^2_{f2} h_{f3}^2 + 6 h^{2bf7} h_{f7}^2 c + 4 h_{f8}^1 c h^{3bf8} + 4 h^{1bf6} h_{f6}^3 c - h^{4bc}) + \\ & \frac{1}{120} \epsilon^5 (-120 h_a^1 d h^{1ba} h_d^1 e h_e^1 f h_f^1 c + 60 h_{f10}^1 h_{f11}^1 h_{f12}^1 c h^{2bf10} - \\ & 30 h^{1bf15} h_{f15}^2 h_{f16}^2 h_{f16}^2 c - 30 h_{f19}^1 h_{f20}^2 h^{2bf19} h_{f20}^2 c - 30 h_{f22}^1 c h^{2bf21} h_{f21}^2 h_{f22}^2 + \\ & 60 h^{1bf1} h_{f1}^1 h_{f2}^1 h_{f3}^2 h_{f3}^2 c + 60 h^{1bf4} h_{f4}^1 h_{f5}^1 c h_{f5}^2 h_{f5}^2 c + 60 h^{1bf7} h_{f8}^1 h_{f9}^1 c h_{f7}^2 h_{f7}^2 c - \\ & 20 h_{f23}^1 h_{f24}^1 c h^{3bf23} + 10 h_{f27}^2 c h^{3bf27} - 20 h^{1bf13} h_{f13}^1 h_{f14}^1 h_{f14}^3 c - \\ & 20 h^{1bf17} h_{f18}^1 c h_{f17}^3 h_{f18}^3 c + 10 h^{2bf26} h_{f26}^3 c + 5 h_{f28}^1 c h^{4bf28} + 5 h^{1bf25} h_{f25}^4 c - h^{5bc}) \end{aligned}$$

And the check

```
In[72]:= %% // org
```

2.33215 Second

$$\begin{aligned} \text{Out}[72]= & \delta_a^c + \\ & \epsilon^6 \left( -h_a^1 b h_b^1 e h^{1cd} h_d^1 f h_e^1 f_1 h_{ff1}^1 + \frac{1}{2} h_b^1 e h^{1cd} h_d^1 f h_{ef}^1 h_a^2 b - \frac{1}{4} h^{1cd} h_d^1 e h_a^2 b h_{be}^2 + \frac{1}{2} h_a^1 b \right. \\ & h^{1cd} h_d^1 e h_e^1 f h_{bf}^2 + \frac{1}{2} h_a^1 b h_b^1 e h_d^1 f h_{ef}^1 h^{2cd} - \frac{1}{4} h_b^1 e h_{de}^1 h_a^2 b h^{2cd} + \frac{1}{8} h_a^2 b h_{bd}^2 h^{2cd} - \\ & \frac{1}{4} h_a^1 b h_d^1 e h_{be}^2 h^{2cd} - \frac{1}{4} h_b^1 e h^{1cd} h_a^2 b h_{de}^2 - \frac{1}{4} h_a^1 b h^{1cd} h_b^2 e h_{de}^2 - \frac{1}{4} h_a^1 b h_b^1 e h^{2cd} h_{de}^2 + \\ & \frac{1}{2} h_a^1 b h_b^1 e h^{1cd} h_e^1 f h_{df}^2 + \frac{1}{2} h_a^1 b h_b^1 e h^{1cd} h_d^1 f h_{ef}^2 - \frac{1}{6} h_b^1 e h^{1cd} h_{de}^1 h_a^3 b + \\ & \frac{1}{12} h^{1cd} h_{bd}^2 h_a^3 b + \frac{1}{12} h_{bd}^1 h^{2cd} h_a^3 b + \frac{1}{12} h^{1cd} h_a^2 b h_{bd}^3 + \frac{1}{12} h_a^1 b h^{2cd} h_{bd}^3 - \\ & \frac{1}{6} h_a^1 b h^{1cd} h_d^1 e h_{be}^3 - \frac{1}{36} h_a^3 b h_b^3 c - \frac{1}{6} h_a^1 b h_b^1 e h_{de}^1 h^{3cd} + \frac{1}{12} h_{bd}^1 h_a^2 b h^{3cd} + \\ & \frac{1}{12} h_a^1 b h_{bd}^2 h^{3cd} - \frac{1}{6} h_a^1 b h_b^1 e h^{1cd} h_{de}^3 + \frac{1}{24} h_{bd}^1 h^{1cd} h_a^4 b - \frac{1}{48} h_b^2 c h_a^4 b + \\ & \left. \frac{1}{24} h_a^1 b h^{1cd} h_{bd}^4 - \frac{1}{48} h_a^2 b h_b^4 c + \frac{1}{24} h_a^1 b h_{bd}^1 h^{4cd} - \frac{1}{120} h_b^1 c h_a^5 b - \frac{1}{120} h_a^1 b h_b^5 c \right) + \\ & \epsilon^7 \left( -\frac{1}{2} h_b^1 e h^{1cd} h_d^1 f h_e^1 f_1 h_{ff1}^1 h_a^2 b + \frac{1}{4} h^{1cd} h_d^1 e h_e^1 f h_a^2 b h_{bf}^2 + \frac{1}{4} h_b^1 e h_d^1 f h_{ef}^1 h_a^2 b h^{2cd} - \right. \\ & \frac{1}{8} h_d^1 e h_a^2 b h_{be}^2 h^{2cd} - \frac{1}{8} h^{1cd} h_a^2 b h_b^2 e h_{de}^2 - \frac{1}{8} h_b^1 e h_a^2 b h^{2cd} h_{de}^2 + \\ & \frac{1}{4} h_b^1 e h^{1cd} h_e^1 f h_a^2 b h_{df}^2 + \frac{1}{4} h_b^1 e h^{1cd} h_d^1 f h_a^2 b h_{ef}^2 + \frac{1}{6} h_b^1 e h^{1cd} h_d^1 f h_{ef}^1 h_a^3 b - \\ & \frac{1}{12} h^{1cd} h_d^1 e h_{be}^2 h_a^3 b - \frac{1}{12} h_b^1 e h_{de}^1 h^{2cd} h_a^3 b + \frac{1}{24} h_{bd}^2 h^{2cd} h_a^3 b - \frac{1}{12} h_b^1 e h^{1cd} h_{de}^2 h_a^3 b + \\ & \frac{1}{24} h_a^2 b h^{2cd} h_{bd}^3 + \frac{1}{36} h^{1cd} h_a^3 b h_{bd}^3 - \frac{1}{12} h^{1cd} h_d^1 e h_a^2 b h_{be}^3 - \frac{1}{12} h_b^1 e h_{de}^1 h_a^2 b h^{3cd} + \\ & \frac{1}{24} h_a^2 b h_{bd}^2 h^{3cd} + \frac{1}{36} h_{bd}^1 h_a^3 b h_{bd}^3 - \frac{1}{12} h_b^1 e h^{1cd} h_a^2 b h_{de}^3 - \frac{1}{24} h_b^1 e h^{1cd} h_{de}^1 h_a^4 b + \\ & \frac{1}{48} h^{1cd} h_{bd}^2 h_a^4 b + \frac{1}{48} h_{bd}^1 h^{2cd} h_a^4 b - \frac{1}{144} h_b^3 c h_a^4 b + \frac{1}{48} h^{1cd} h_a^2 b h_{bd}^4 - \\ & \left. \frac{1}{144} h_a^3 b h_b^4 c + \frac{1}{48} h_{bd}^1 h_a^2 b h^{4cd} + \frac{1}{120} h_{bd}^1 h^{1cd} h_a^5 b - \frac{1}{240} h_b^2 c h_a^5 b - \frac{1}{240} h_a^2 b h_b^5 c \right) + \end{aligned}$$

$$\begin{aligned}
& \epsilon^8 \left( -\frac{1}{6} h_b^1 e h^{1cd} h_d^1 f h_e^1 f_1 h_{ff1}^1 h_a^3 b + \frac{1}{12} h^{1cd} h_d^1 e h_e^1 f h_{bf}^2 h_a^3 b + \frac{1}{12} h_b^1 e h_d^1 f h_{ef}^1 h^{2cd} h_a^3 b - \right. \\
& \quad \frac{1}{24} h_d^1 e h_{be}^2 h^{2cd} h_a^3 b - \frac{1}{24} h^{1cd} h_b^2 e h_{de}^2 h_a^3 b - \frac{1}{24} h_b^1 e h^{2cd} h_{de}^2 h_a^3 b + \\
& \quad \frac{1}{12} h_b^1 e h^{1cd} h_e^1 f h_{df}^2 h_a^3 b + \frac{1}{12} h_b^1 e h^{1cd} h_d^1 f h_{ef}^2 h_a^3 b + \frac{1}{72} h^{2cd} h_a^3 b h_{bd}^3 - \\
& \quad \frac{1}{36} h^{1cd} h_d^1 e h_a^3 b h_{be}^3 - \frac{1}{36} h_b^1 e h_{de}^1 h_a^3 b h^{3cd} + \frac{1}{72} h_{bd}^2 h_a^3 b h^{3cd} - \frac{1}{36} h_b^1 e h^{1cd} h_a^3 b h_{de}^3 + \\
& \quad \frac{1}{24} h_b^1 e h^{1cd} h_d^1 f h_{ef}^1 h_a^4 b - \frac{1}{48} h^{1cd} h_d^1 e h_{be}^2 h_a^4 b - \frac{1}{48} h_b^1 e h_{de}^1 h^{2cd} h_a^4 b + \\
& \quad \frac{1}{96} h_{bd}^2 h^{2cd} h_a^4 b - \frac{1}{48} h_b^1 e h^{1cd} h_{de}^2 h_a^4 b + \frac{1}{144} h^{1cd} h_{bd}^3 h_a^4 b + \frac{1}{144} h_{bd}^1 h^{3cd} h_a^4 b + \\
& \quad \frac{1}{144} h^{1cd} h_a^3 b h_{bd}^4 - \frac{1}{576} h_a^4 b h_{bc}^4 + \frac{1}{144} h_{bd}^1 h_a^3 b h^{4cd} - \frac{1}{120} h_b^1 e h^{1cd} h_{de}^1 h_a^5 b + \\
& \quad \left. \frac{1}{240} h^{1cd} h_{bd}^2 h_a^5 b + \frac{1}{240} h_{bd}^1 h^{2cd} h_a^5 b - \frac{1}{720} h_b^3 c h_a^5 b - \frac{1}{720} h_a^3 b h_{bc}^5 \right) + \\
& \epsilon^9 \left( -\frac{1}{24} h_b^1 e h^{1cd} h_d^1 f h_e^1 f_1 h_{ff1}^1 h_a^4 b + \frac{1}{48} h^{1cd} h_d^1 e h_e^1 f h_{bf}^2 h_a^4 b + \right. \\
& \quad \frac{1}{48} h_b^1 e h_d^1 f h_{ef}^1 h^{2cd} h_a^4 b - \frac{1}{96} h_d^1 e h_{be}^2 h^{2cd} h_a^4 b - \frac{1}{96} h^{1cd} h_b^2 e h_{de}^2 h_a^4 b - \\
& \quad \frac{1}{96} h_b^1 e h^{2cd} h_{de}^2 h_a^4 b + \frac{1}{48} h_b^1 e h^{1cd} h_e^1 f h_{df}^2 h_a^4 b + \frac{1}{48} h_b^1 e h^{1cd} h_d^1 f h_{ef}^2 h_a^4 b + \\
& \quad \frac{1}{288} h^{2cd} h_{bd}^3 h_a^4 b - \frac{1}{144} h^{1cd} h_d^1 e h_{be}^3 h_a^4 b - \frac{1}{144} h_b^1 e h_{de}^1 h^{3cd} h_a^4 b + \\
& \quad \frac{1}{288} h_{bd}^2 h^{3cd} h_a^4 b - \frac{1}{144} h_b^1 e h^{1cd} h_{de}^3 h_a^4 b + \frac{1}{576} h^{1cd} h_a^4 b h_{bd}^4 + \\
& \quad \frac{1}{576} h_{bd}^1 h_a^4 b h^{4cd} + \frac{1}{120} h_b^1 e h^{1cd} h_d^1 f h_{ef}^1 h_a^5 b - \frac{1}{240} h^{1cd} h_d^1 e h_{be}^2 h_a^5 b - \\
& \quad \frac{1}{240} h_b^1 e h_{de}^1 h^{2cd} h_a^5 b + \frac{1}{480} h_{bd}^2 h^{2cd} h_a^5 b - \frac{1}{240} h_b^1 e h^{1cd} h_{de}^2 h_a^5 b + \\
& \quad \left. \frac{1}{720} h^{1cd} h_{bd}^3 h_a^5 b + \frac{1}{720} h_{bd}^1 h^{3cd} h_a^5 b - \frac{h_b^4 c h_a^5 b}{2880} - \frac{h_a^4 b h_{bc}^5}{2880} \right) + \\
& \epsilon^{10} \left( -\frac{1}{120} h_b^1 e h^{1cd} h_d^1 f h_e^1 f_1 h_{ff1}^1 h_a^5 b + \frac{1}{240} h^{1cd} h_d^1 e h_e^1 f h_{bf}^2 h_a^5 b + \right. \\
& \quad \frac{1}{240} h_b^1 e h_d^1 f h_{ef}^1 h^{2cd} h_a^5 b - \frac{1}{480} h_d^1 e h_{be}^2 h^{2cd} h_a^5 b - \frac{1}{480} h^{1cd} h_b^2 e h_{de}^2 h_a^5 b - \\
& \quad \frac{1}{480} h_b^1 e h^{2cd} h_{de}^2 h_a^5 b + \frac{1}{240} h_b^1 e h^{1cd} h_e^1 f h_{df}^2 h_a^5 b + \frac{1}{240} h_b^1 e h^{1cd} h_d^1 f h_{ef}^2 h_a^5 b + \\
& \quad \frac{h_{bd}^{2cd} h_{bd}^3 h_a^5 b}{1440} - \frac{1}{720} h^{1cd} h_d^1 e h_{be}^3 h_a^5 b - \frac{1}{720} h_b^1 e h_{de}^1 h^{3cd} h_a^5 b + \frac{h_{bd}^2 h^{3cd} h_a^5 b}{1440} - \\
& \quad \left. \frac{1}{720} h_b^1 e h^{1cd} h_{de}^3 h_a^5 b + \frac{h^{1cd} h_{bd}^4 h_a^5 b}{2880} + \frac{h_{bd}^1 h^{4cd} h_a^5 b}{2880} - \frac{h_a^5 b h_{bc}^5}{14400} \right)
\end{aligned}$$

### 3.2. Perturbations of the determinant of the metric and the antisymmetric tensor

Now we compute the perturbations of the determinant of the metric up to sixth order. Two comments are in order:

First, note that the result is always proportional to the determinant itself times a combination of contracted metric perturbations. This proportionality scalar depends on the dimension we are considering. In fact, for a given dimension  $N$ , there cannot be products of more than  $N$  metric perturbations (recall that  $N=4$  in our examples). Timings are highly dependent on the dimensions because the internal algorithms use determinants of  $N \times N$  matrices.

Second, the concept of determinant is basis-dependent (actually a density) and hence can only be treated properly using `xCoba`, the `xAct` package for component computations in a given basis. `xTensor` implements a fake basis called `AIndex`, which we will assume throughout this subsection.

Low order perturbations of the determinant:

```
In[73]:= Perturbation[Detg[]] // ExpandPerturbation // Simplification
```

0.180011 Second

```
Out[73]=  $\tilde{g} h^1_a$ 
```

```
In[74]:= Perturbation[Detg[], 2] // ExpandPerturbation // Simplification
```

0.408026 Second

```
Out[74]=  $\tilde{g} (-h^1_{ab} h^1_{ab} + h^1_a h^1_b + h^2_a)$ 
```

```
In[75]:= Perturbation[Detg[], 3] // ExpandPerturbation // Simplification
```

0.904057 Second

```
Out[75]=  $\tilde{g} (2 h^1_a h^1_{ab} h^1_{bc} - 3 h^1_{ab} h^2_{ab} + h^1_a (-3 h^1_{bc} h^1_{bc} + h^1_b h^1_c + 3 h^2_b) + h^3_a)$ 
```

```
In[76]:= Perturbation[Detg[], 4] // ExpandPerturbation // Simplification
```

1.95212 Second

```
Out[76]=  $\tilde{g} (3 h^1_{ab} h^1_{ab} h^1_{cd} h^1_{cd} - 3 h^2_{ab} h^2_{ab} - 6 h^1_a h^1_{ab} (h^1_b h^1_{cd} - 2 h^2_{bc}) + 3 h^2_a h^2_b - 6 h^1_{ab} h^1_{ab} h^2_c - 4 h^1_{ab} h^3_{ab} + h^1_a (8 h^1_b h^1_{bc} h^1_{cd} + h^1_b (-6 h^1_{cd} h^1_{cd} + h^1_c h^1_d + 6 h^2_c) + 4 (-3 h^1_{bc} h^2_{bc} + h^3_b)) + h^4_a)$ 
```

Now the perturbative order will exceed the dimension. Still we get just products of up to four metric perturbations. In other dimensions these formulas would be wrong:

```
In[77]:= Perturbation[Detg[], 5] // ExpandPerturbation // Simplification
```

3.74823 Second

```
Out[77]=  $\tilde{g} (-10 h^2_{ab} h^3_{ab} + 10 h^2_a h^3_b + 5 h^1_{ab} (6 h^2_a h^2_{bc} + 6 h^1_{ab} h^1_{cd} h^2_{cd} - 6 h^2_{ab} h^2_c + 4 h^1_a (-3 h^1_b h^2_{cd} + h^1_{bc} h^2_d + h^3_{bc}) - 2 h^1_{ab} h^3_c - h^4_{ab}) - 5 h^1_a (3 h^2_{bc} h^2_{bc} - 12 h^1_b h^1_{bc} h^2_{cd} + 6 h^1_b h^1_{cd} h^2_{cd} - 3 h^2_b h^2_c + 6 h^1_{bc} h^1_{bc} h^2_d - 2 h^1_b h^1_c h^2_d + 4 h^1_{bc} h^3_{bc} - 2 h^1_b h^3_c - h^4_b) + h^5_a)$ 
```

```
In[78]:= Perturbation[Detg[], 6] // ExpandPerturbation // Simplification
```

7.05644 Second

```
Out[78]=  $\tilde{g} (-45 h^2_a h^2_{bc} h^2_{bc} + 180 h^1_a h^1_{bc} h^2_d h^2_{cd} + 15 h^2_a h^2_b h^2_c - 45 h^1_a h^1_b h^2_{cd} h^2_{cd} - 180 h^1_a h^1_{bc} h^2_{bc} h^2_d + 45 h^1_a h^1_b h^2_c h^2_d - 10 h^3_{ab} h^3_{ab} - 60 h^1_a h^2_{bc} h^3_{bc} + 30 h^2_a (h^2_{ab} h^2_{bc} + 4 h^1_{ab} h^3_{bc}) + 10 h^3_a h^3_b + 120 h^1_a h^1_d h^1_{bc} h^3_{cd} - 60 h^1_a h^1_b h^1_{cd} h^3_{cd} + 60 h^1_a h^2_b h^3_c - 60 h^1_a h^1_{bc} h^1_{bc} h^3_d + 20 h^1_a h^1_b h^1_c h^3_d - 15 h^2_{ab} h^4_{ab} - 30 h^1_a h^1_{bc} h^4_{bc} + 15 h^2_a h^4_b + 15 h^1_a h^1_b h^4_c - h^1_{ab} (-45 h^1_{ab} h^2_{cd} h^2_{cd} + 45 h^1_{ab} h^2_c h^2_d + 60 h^2_c h^3_{ab} + 30 h^1_{ac} h^2_{bd} - 3 h^2_{ab} h^2_{cd} - 2 h^1_{ab} h^3_{cd}) + 60 h^2_{ab} h^3_c + 10 h^1_a (18 h^2_b h^2_{cd} - 18 h^2_{bc} h^2_d + 12 h^1_b h^3_{cd} - 4 h^1_{bc} h^3_d - 3 h^4_{bc}) + 15 h^1_{ab} h^4_c + 6 h^5_{ab}) + 6 h^1_a h^5_b + h^6_a)$ 
```



Perturbations of the antisymmetric volume–form tensor are given in terms of the perturbations of the determinant of the metric and have a similar structure.

---

At first–order

```
In[79]:= Perturbation[epsilon[-a, -b, -c, -d]] // ExpandPerturbation // Simplification
0.16801 Second
```

```
Out[79]=  $\frac{1}{2} \epsilon_{abcd} h^1{}^e{}_e$ 
```

---

At third–order (note the head `Scalar`; see the documentation of `xTensor`` for a description)

```
In[80]:= Perturbation[epsilon[-a, -b, -c, -d], 3] // ExpandPerturbation // Simplification
1.85212 Second
```

```
Out[80]=  $\frac{1}{8} \epsilon_{abcd} \left( 8 h^1{}^e{}_{ff1} h^1{}^{ef} h^1{}_{ff1} - 12 h^1{}^{ef} h^2{}_{ef} - \right.$   

 $\left. 2 h^1{}^e{}_e \left( 3 h^1{}_{ff1} h^1{}^{ff1} + h^1{}^f{}_f h^1{}^{f1}{}_{f1} - 3 h^2{}^f{}_f \right) + 4 h^3{}^e{}_e + 3 \text{Scalar}[h^1{}^a{}_a]^3 \right)$ 
```

---

Note that now there are five `h` factors multiplying due to the perturbation of `Sqrt[Detg[]]`, and not only of `Detg[]`. The formula is still valid only in dimension 4 for the reasons above:

```
In[81]:= Perturbation[epsilon[-a, -b, -c, -d], 5] // ExpandPerturbation // Simplification
20.3413 Second
```

```
Out[81]=  $\frac{1}{32} \epsilon_{abcd} \left( 480 h^1{}^{ef} h^2{}^e{}_{ff1} h^2{}_{ff1} - 960 h^1{}^e{}_{ff1} h^1{}^{ef} h^1{}^f{}_{f10} h^2{}_{f1f10} - 240 h^1{}^{ef} h^2{}_{ef} h^2{}^f{}_{f1} + \right.$   

 $160 h^1{}^e{}_{ff1} h^1{}^{ef} h^1{}_{ff1} h^2{}^f{}_{f10} - 160 h^2{}^{ef} h^3{}_{ef} + 320 h^1{}^e{}_{ff1} h^1{}^{ef} h^3{}_{ff1} + 80 h^2{}^e{}_e h^3{}^f{}_f -$   

 $80 h^1{}^{ef} h^4{}_{ef} + 16 h^5{}^e{}_e + 240 h^1{}^e{}_{ff1} h^1{}^{ef} h^1{}_{ff1} \text{Scalar}[h^1{}^a{}_a]^2 - 360 h^1{}^{ef} h^2{}_{ef} \text{Scalar}[h^1{}^a{}_a]^2 +$   

 $120 h^3{}^e{}_e \text{Scalar}[h^1{}^a{}_a]^2 - 300 h^2{}^e{}_e \text{Scalar}[h^1{}^a{}_a]^3 + 105 \text{Scalar}[h^1{}^a{}_a]^5 +$   

 $20 h^1{}_{ef} h^1{}^{ef} \left( 8 h^1{}^f{}_{f1} h^1{}^{f1}{}_{f10} h^1{}_{f10f11} + 12 h^1{}^{f1}{}_{f10} h^2{}_{f1f10} - 4 h^3{}^f{}_{f1} + 15 \text{Scalar}[h^1{}^a{}_a]^3 \right) +$   

 $20 h^1{}^e{}_e \left( -18 h^1{}_{ff1} h^1{}^{ff1} h^1{}_{f10f11} h^1{}^{f10f11} - 6 h^2{}_{ff1} h^2{}^{ff1} + \right.$   

 $12 h^1{}^f{}_{f10} h^1{}^{ff1} \left( h^1{}_{f1} h^1{}^{f11} h^1{}_{f10f11} + 2 h^2{}_{f1f10} \right) - 6 h^2{}^f{}_f h^2{}^f{}_{f1} + 12 h^1{}_{ff1} h^1{}^{ff1} h^2{}^f{}_{f10} -$   

 $8 h^1{}^{ff1} h^3{}_{ff1} + 2 h^4{}^f{}_f - 18 h^1{}_{ff1} h^1{}^{ff1} \text{Scalar}[h^1{}^a{}_a]^2 + 18 h^2{}^f{}_f \text{Scalar}[h^1{}^a{}_a]^2 -$   

 $h^1{}^f{}_f \left( 24 h^1{}_{f1} h^1{}^{f11} h^1{}^{f1}{}_{f10} h^1{}_{f10f11} - 12 h^1{}^{f1}{}_{f10} h^2{}_{f1f10} + 4 h^3{}^f{}_{f1} + 15 \text{Scalar}[h^1{}^a{}_a]^3 + \right.$   

 $\left. h^1{}^f{}_{f1} \left( -28 h^1{}_{f10f11} h^1{}^{f10f11} + 6 h^1{}^f{}_{f10} h^1{}^{f11}{}_{f11} + 20 h^2{}^f{}_{f10} - 6 \text{Scalar}[h^1{}^a{}_a]^2 \right) \right) +$   

 $9 \text{Scalar}[h^1{}_{ab} h^1{}^{ab}]^2 - 18 \text{Scalar}[h^1{}_{ab} h^1{}^{ab}] \text{Scalar}[h^1{}^a{}_a h^1{}^b{}_b] +$   

 $9 \text{Scalar}[h^1{}^a{}_a h^1{}^b{}_b]^2 - 18 \text{Scalar}[h^1{}_{ab} h^1{}^{ab}] \text{Scalar}[h^2{}^a{}_a] +$   

 $18 \text{Scalar}[h^1{}^a{}_a h^1{}^b{}_b] \text{Scalar}[h^2{}^a{}_a] + 9 \text{Scalar}[h^2{}^a{}_a]^2 \left. \right)$ 
```

### 3.3. Perturbations of the Christoffel symbols

In order to obtain the perturbation of the Christoffel symbols, we can write them in terms of partial derivatives of the metric and then apply the `ExpandPerturbation` command that will only act on metrics and inverse metrics. It is faster to do it directly applying `ExpandPerturbation`, that will calculate the perturbation by making use of non–recursive formulas for the  $n$ –th order perturbation of the Christoffel symbols. We shall check and compare the relative efficiency of both methods.

---

First-order perturbation of the Christoffel in terms of the metric and its inverse

`In[82]:= Perturbation[ChristoffelCD[a, -b, -c] // ChristoffelToMetric]`

$$\text{Out}[82]= \frac{1}{2} (\Delta[g^{\text{ad}}] (g_{\text{cd},b} + g_{\text{bd},c} - g_{\text{bc},d}) + g^{\text{ad}} (h^1_{\text{cd},b} + h^1_{\text{bd},c} - h^1_{\text{bc},d}))$$

---

We will frequently convert partial derivatives into covariant derivatives to help canonicalizing:

`In[83]:= CovDToChristoffel[% // ExpandPerturbation, PD, CD] // org`

$$\text{Out}[83]= -\frac{h^1_{bc}{}^{ia}}{2} + \frac{h^1{}^a{}_{c;b}}{2} + \frac{h^1{}^a{}_{b;c}}{2}$$

---

Comparison with the direct action of ExpandPerturbation

`In[84]:= % - (Perturbation[ChristoffelCD[a, -b, -c] // ExpandPerturbation) // Simplification`

`Out[84]= 0`

---

We do the same for the second-order perturbation

`In[85]:= Perturbation[ChristoffelCD[a, -b, -c] // ChristoffelToMetric, 2]`

$$\text{Out}[85]= \frac{1}{2} (\Delta^2[g^{\text{ad}}] (g_{\text{cd},b} + g_{\text{bd},c} - g_{\text{bc},d}) + 2\Delta[g^{\text{ad}}] (h^1_{\text{cd},b} + h^1_{\text{bd},c} - h^1_{\text{bc},d}) + g^{\text{ad}} (h^2_{\text{cd},b} + h^2_{\text{bd},c} - h^2_{\text{bc},d}))$$

`In[86]:= CovDToChristoffel[% // ExpandPerturbation, PD, CD] // org`

0.244015 Second

$$\text{Out}[86]= -\frac{h^2_{bc}{}^{ia}}{2} - h^1{}^{\text{ad}} h^1{}_{\text{cd};b} + \frac{h^2{}^a{}_{c;b}}{2} - h^1{}^{\text{ad}} h^1{}_{\text{bd};c} + \frac{h^2{}^a{}_{b;c}}{2} + h^1{}^{\text{ad}} h^1{}_{\text{bc};d}$$

`In[87]:= % - ExpandPerturbation[Perturbation[ChristoffelCD[a, -b, -c], 2]] // ToCanonical`

`Out[87]= 0`

---

At third order

`In[88]:= Perturbation[ChristoffelCD[a, -b, -c] // ChristoffelToMetric, 3]`

$$\text{Out}[88]= \frac{1}{2} (\Delta^3[g^{\text{ad}}] (g_{\text{cd},b} + g_{\text{bd},c} - g_{\text{bc},d}) + 3\Delta^2[g^{\text{ad}}] (h^1_{\text{cd},b} + h^1_{\text{bd},c} - h^1_{\text{bc},d}) + 3\Delta[g^{\text{ad}}] (h^2_{\text{cd},b} + h^2_{\text{bd},c} - h^2_{\text{bc},d}) + g^{\text{ad}} (h^3_{\text{cd},b} + h^3_{\text{bd},c} - h^3_{\text{bc},d}))$$

`In[89]:= CovDToChristoffel[% // ExpandPerturbation, PD, CD] // org`

0.516032 Second

$$\text{Out}[89]= -\frac{h^3_{bc}{}^{ia}}{2} - \frac{3}{2} h^2{}^{\text{ad}} h^1{}_{\text{cd};b} + 3 h^1{}^{\text{ad}} h^1{}_{\text{d}}{}^e h^1{}_{\text{ce};b} - \frac{3}{2} h^1{}^{\text{ad}} h^2{}_{\text{cd};b} + \frac{h^3{}^a{}_{c;b}}{2} - \frac{3}{2} h^2{}^{\text{ad}} h^1{}_{\text{bd};c} + 3 h^1{}^{\text{ad}} h^1{}_{\text{d}}{}^e h^1{}_{\text{be};c} - \frac{3}{2} h^1{}^{\text{ad}} h^2{}_{\text{bd};c} + \frac{h^3{}^a{}_{b;c}}{2} + \frac{3}{2} h^2{}^{\text{ad}} h^1{}_{\text{bc};d} + \frac{3}{2} h^1{}^{\text{ad}} h^2{}_{\text{bc};d} - 3 h^1{}^{\text{ad}} h^1{}_{\text{d}}{}^e h^1{}_{\text{bc};e}$$

```
In[90]:= % - ExpandPerturbation[Perturbation[ChristoffelCD[a, -b, -c], 3]] // ToCanonical
```

```
Out[90]= 0
```

---

And finally at four order

```
In[91]:= Perturbation[ChristoffelCD[a, -b, -c] // ChristoffelToMetric, 4]
```

```
Out[91]=  $\frac{1}{2} (\Delta^4 [g^{ad}] (g_{cd,b} + g_{bd,c} - g_{bc,d}) +$   

 $4 \Delta^3 [g^{ad}] (h^1_{cd,b} + h^1_{bd,c} - h^1_{bc,d}) + 6 \Delta^2 [g^{ad}] (h^2_{cd,b} + h^2_{bd,c} - h^2_{bc,d}) +$   

 $4 \Delta [g^{ad}] (h^3_{cd,b} + h^3_{bd,c} - h^3_{bc,d}) + g^{ad} (h^4_{cd,b} + h^4_{bd,c} - h^4_{bc,d}))$ 
```

```
In[92]:= CovDToChristoffel[% // ExpandPerturbation, PD, CD] // org
```

```
1.18007 Second
```

```
Out[92]=  $-\frac{h^4_{bc}{}^{ia}}{2} - 2 h^3{}^{ad} h^1_{cd;b} + 6 h^1{}_d{}^e h^2{}^{ad} h^1_{ce;b} + 6 h^1{}^{ad} h^2{}_d{}^e h^1_{ce;b} - 12 h^1{}^{ad} h^1{}_d{}^e h^1{}_e{}^f h^1_{cf;b} -$   

 $3 h^2{}^{ad} h^2_{cd;b} + 6 h^1{}^{ad} h^1{}_d{}^e h^2_{ce;b} - 2 h^1{}^{ad} h^3_{cd;b} + \frac{h^4{}_c{}^a{}^b}{2} - 2 h^3{}^{ad} h^1_{bd;c} +$   

 $6 h^1{}_d{}^e h^2{}^{ad} h^1_{be;c} + 6 h^1{}^{ad} h^2{}_d{}^e h^1_{be;c} - 12 h^1{}^{ad} h^1{}_d{}^e h^1{}_e{}^f h^1_{bf;c} - 3 h^2{}^{ad} h^2_{bd;c} +$   

 $6 h^1{}^{ad} h^1{}_d{}^e h^2_{be;c} - 2 h^1{}^{ad} h^3_{bd;c} + \frac{h^4{}_b{}^a{}^c}{2} + 2 h^3{}^{ad} h^1_{bc;d} + 3 h^2{}^{ad} h^2_{bc;d} + 2 h^1{}^{ad} h^3_{bc;d} -$   

 $6 h^1{}_d{}^e h^2{}^{ad} h^1_{bc;e} - 6 h^1{}^{ad} h^2{}_d{}^e h^1_{bc;e} - 6 h^1{}^{ad} h^1{}_d{}^e h^2_{bc;e} + 12 h^1{}^{ad} h^1{}_d{}^e h^1{}_e{}^f h^1_{bc;f}$ 
```

```
In[93]:= % - ExpandPerturbation[Perturbation[ChristoffelCD[a, -b, -c], 4]] // ToCanonical
```

```
0.256015 Second
```

```
Out[93]= 0
```

---

Note that the first method is much slower: let us compute the seventh-order perturbations:

```
In[94]:= tmp1 =  
CovDToChristoffel[Perturbation[ChristoffelCD[a, -b, -c] // ChristoffelToMetric,  
7] // ExpandPerturbation, PD, CD] // org;
```

```
14.2929 Second
```

```
In[95]:= tmp2 = ExpandPerturbation[Perturbation[ChristoffelCD[a, -b, -c], 7]] // org;
```

```
2.37615 Second
```

```
In[96]:= tmp1 - tmp2 // ToCanonical
```

```
0.996062 Second
```

```
Out[96]= 0
```

### 3.4. Perturbations of the Riemann tensor

As we have done for the Christoffel symbols, we can calculate the perturbations of Riemann in two ways. The first one will start from the expression of Riemann tensor in terms of the metric, its inverse and partial derivatives. The second will apply directly `ExpandPerturbation`.

At first order going through the metric

```
In[97]:= Perturbation[RiemannCD[-a, -b, -c, d], 1] // RiemannToChristoffel //
ChristoffelToMetric
```

$$\begin{aligned}
\text{Out}[97]= & \frac{1}{2} (\Delta[g^{de}] (\mathcal{G}_{ce,a,b} + \mathcal{G}_{ae,c,b} - \mathcal{G}_{ac,e,b}) + g^{de} (h^1_{ce,a,b} + h^1_{ae,c,b} - h^1_{ac,e,b}) + \\
& \partial_b \Delta[g^{de}] (\mathcal{G}_{ce,a} + \mathcal{G}_{ae,c} - \mathcal{G}_{ac,e}) - g^{df} g^{ef1} g_{ff1,b} (h^1_{ce,a} + h^1_{ae,c} - h^1_{ac,e})) + \\
& \frac{1}{2} (-\Delta[g^{de}] (\mathcal{G}_{ce,b,a} + \mathcal{G}_{be,c,a} - \mathcal{G}_{bc,e,a}) - g^{de} (h^1_{ce,b,a} + h^1_{be,c,a} - h^1_{bc,e,a}) - \\
& \partial_a \Delta[g^{de}] (\mathcal{G}_{ce,b} + \mathcal{G}_{be,c} - \mathcal{G}_{bc,e}) + g^{df} g^{ef1} g_{ff1,a} (h^1_{ce,b} + h^1_{be,c} - h^1_{bc,e})) + \\
& \frac{1}{4} g^{df} (\mathcal{G}_{ef,b} + \mathcal{G}_{bf,e} - \mathcal{G}_{be,f}) (\Delta[g^{ef1}] (\mathcal{G}_{cf1,a} + \mathcal{G}_{af1,c} - \mathcal{G}_{ac,f1}) + g^{ef1} (h^1_{cf1,a} + h^1_{af1,c} - h^1_{ac,f1})) \cdot \\
& \frac{1}{4} g^{ef} (\mathcal{G}_{cf,b} + \mathcal{G}_{bf,c} - \mathcal{G}_{bc,f}) (\Delta[g^{df1}] (\mathcal{G}_{ef1,a} + \mathcal{G}_{af1,e} - \mathcal{G}_{ae,f1}) + g^{df1} (h^1_{ef1,a} + h^1_{af1,e} - h^1_{ae,f1})) \cdot \\
& \frac{1}{4} g^{df} (\mathcal{G}_{ef,a} + \mathcal{G}_{af,e} - \mathcal{G}_{ae,f}) (\Delta[g^{ef1}] (\mathcal{G}_{cf1,b} + \mathcal{G}_{bf1,c} - \mathcal{G}_{bc,f1}) + g^{ef1} (h^1_{cf1,b} + h^1_{bf1,c} - h^1_{bc,f1})) \cdot \\
& \frac{1}{4} g^{ef} (\mathcal{G}_{cf,a} + \mathcal{G}_{af,c} - \mathcal{G}_{ac,f}) (\Delta[g^{df1}] (\mathcal{G}_{ef1,b} + \mathcal{G}_{bf1,e} - \mathcal{G}_{be,f1}) + g^{df1} (h^1_{ef1,b} + h^1_{bf1,e} - h^1_{be,f1}))
\end{aligned}$$

```
In[98]:= ChangeCovD[% // ExpandPerturbation, PD, CD] // org
```

5.30433 Second

$$\text{Out}[98]= -\frac{h^1_{c;ib;a}}{2} - \frac{h^1_{b;ic;a}}{2} + \frac{h^1_{bc;ia}}{2} + \frac{h^1_{c;ia;b}}{2} + \frac{h^1_{a;ic;b}}{2} - \frac{h^1_{ac;ib}}{2}$$

Applying directly ExpandPerturbation

```
In[99]:= Perturbation[RiemannCD[-a, -b, -c, d], 1] // ExpandPerturbation
```

$$\text{Out}[99]= \frac{1}{2} (-h^1_{c;ib;a} - h^1_{b;ic;a} + h^1_{cb;ia}) + \frac{1}{2} (h^1_{c;ia;b} + h^1_{a;ic;b} - h^1_{ca;ib})$$

Comparison

```
In[100]:= %% - % // ToCanonical
```

```
Out[100]= 0
```

That result is usually presented with only four covariant derivatives, but with 10 Riemann terms (which vanish in vacuum):

```
In[101]:= %% // SortCovDs // ToCanonical
```

0.124008 Second

$$\begin{aligned}
\text{Out}[101]= & -\frac{1}{2} h^{1de} R_{abce} - \frac{1}{2} h^1_{c^e} R_{ab^d} - \frac{1}{2} h^{1de} R_{acbe} - \frac{1}{2} h^1_{b^e} R_{ac^d} + \\
& \frac{1}{2} h^1_{c^e} R_{a^d_{be}} + \frac{1}{2} h^1_{b^e} R_{a^d_{ce}} + \frac{1}{2} h^{1de} R_{aebc} - \frac{1}{2} h^1_{c^e} R_{aeb^d} + \\
& \frac{1}{2} h^1_{a^e} R_{bc^d} - \frac{1}{2} h^1_{a^e} R_{b^d_{ce}} - \frac{h^1_{b;ia;c}}{2} + \frac{h^1_{a;ib;c}}{2} + \frac{h^1_{bc;ia}}{2} - \frac{h^1_{ac;ib}}{2}
\end{aligned}$$

At second order going through the metric

In[102]:=

**Perturbation[RiemannCD[-a, -b, -c, d], 2] // RiemannToChristoffel // ChristoffelToMetric**

Out[102]=

$$\begin{aligned}
& \frac{1}{2} (\Delta^2 [g^{de}] (g_{ce,a,b} + g_{ae,c,b} - g_{ac,e,b}) + \\
& \quad g^{de} (h_{ce,a,b}^2 + h_{ae,c,b}^2 - h_{ac,e,b}^2) + \partial_b \Delta^2 [g^{de}] (g_{ce,a} + g_{ae,c} - g_{ac,e}) + \\
& \quad 2 (\Delta [g^{de}] (h_{ce,a,b}^1 + h_{ae,c,b}^1 - h_{ac,e,b}^1) + \partial_b \Delta [g^{de}] (h_{ce,a}^1 + h_{ae,c}^1 - h_{ac,e}^1)) - \\
& \quad g^{df} g^{ef1} g_{ff1,b} (h_{ce,a}^2 + h_{ae,c}^2 - h_{ac,e}^2)) + \\
& \frac{1}{2} (-\Delta^2 [g^{de}] (g_{ce,b,a} + g_{be,c,a} - g_{bc,e,a}) - g^{de} (h_{ce,b,a}^2 + h_{be,c,a}^2 - h_{bc,e,a}^2) - \\
& \quad \partial_a \Delta^2 [g^{de}] (g_{ce,b} + g_{be,c} - g_{bc,e}) - \\
& \quad 2 (\Delta [g^{de}] (h_{ce,b,a}^1 + h_{be,c,a}^1 - h_{bc,e,a}^1) + \partial_a \Delta [g^{de}] (h_{ce,b}^1 + h_{be,c}^1 - h_{bc,e}^1)) + \\
& \quad g^{df} g^{ef1} g_{ff1,a} (h_{ce,b}^2 + h_{be,c}^2 - h_{bc,e}^2)) + \\
& \frac{1}{2} (\Delta [g^{df}] (g_{ef,b} + g_{bf,e} - g_{be,f}) + g^{df} (h_{ef,b}^1 + h_{bf,e}^1 - h_{be,f}^1)) \\
& \quad (\Delta [g^{ef1}] (g_{cf1,a} + g_{af1,c} - g_{ac,f1}) + g^{ef1} (h_{cf1,a}^1 + h_{af1,c}^1 - h_{ac,f1}^1)) - \\
& \frac{1}{2} (\Delta [g^{df}] (g_{ef,a} + g_{af,e} - g_{ae,f}) + g^{df} (h_{ef,a}^1 + h_{af,e}^1 - h_{ae,f}^1)) \\
& \quad (\Delta [g^{ef1}] (g_{cf1,b} + g_{bf1,c} - g_{bc,f1}) + g^{ef1} (h_{cf1,b}^1 + h_{bf1,c}^1 - h_{bc,f1}^1)) + \\
& \frac{1}{4} g^{df} (g_{ef,b} + g_{bf,e} - g_{be,f}) (\Delta^2 [g^{ef1}] (g_{cf1,a} + g_{af1,c} - g_{ac,f1}) + \\
& \quad 2 \Delta [g^{ef1}] (h_{cf1,a}^1 + h_{af1,c}^1 - h_{ac,f1}^1) + g^{ef1} (h_{cf1,a}^2 + h_{af1,c}^2 - h_{ac,f1}^2)) - \\
& \frac{1}{4} g^{ef} (g_{cf,b} + g_{bf,c} - g_{bc,f}) (\Delta^2 [g^{df1}] (g_{ef1,a} + g_{af1,e} - g_{ae,f1}) + \\
& \quad 2 \Delta [g^{df1}] (h_{ef1,a}^1 + h_{af1,e}^1 - h_{ae,f1}^1) + g^{df1} (h_{ef1,a}^2 + h_{af1,e}^2 - h_{ae,f1}^2)) - \\
& \frac{1}{4} g^{df} (g_{ef,a} + g_{af,e} - g_{ae,f}) (\Delta^2 [g^{ef1}] (g_{cf1,b} + g_{bf1,c} - g_{bc,f1}) + \\
& \quad 2 \Delta [g^{ef1}] (h_{cf1,b}^1 + h_{bf1,c}^1 - h_{bc,f1}^1) + g^{ef1} (h_{cf1,b}^2 + h_{bf1,c}^2 - h_{bc,f1}^2)) + \\
& \frac{1}{4} g^{ef} (g_{cf,a} + g_{af,c} - g_{ac,f}) (\Delta^2 [g^{df1}] (g_{ef1,b} + g_{bf1,e} - g_{be,f1}) + \\
& \quad 2 \Delta [g^{df1}] (h_{ef1,b}^1 + h_{bf1,e}^1 - h_{be,f1}^1) + g^{df1} (h_{ef1,b}^2 + h_{bf1,e}^2 - h_{be,f1}^2))
\end{aligned}$$

In[103]:=

**ChangeCovD[% // ExpandPerturbation, PD, CD] // org**

19.3452 Second

Out[103]=

$$\begin{aligned}
& h^{1de} h^1_{ce;bia} - \frac{h^2_{c;bia}}{2} + h^{1de} h^1_{be;c;a} - \frac{h^2_{b;cia}}{2} + \frac{h^2_{bc;ia}}{2} - h^{1de} h^1_{bc;eia} + \\
& \frac{1}{2} h^{1d}_{e;a} h^1_{c;ib} - \frac{1}{2} h^1_{c;ia} h^1_{e;b} - h^{1de} h^1_{ce;a;b} + \frac{h^2_{c;ia;b}}{2} - h^{1de} h^1_{ae;c;b} + \\
& \frac{h^2_{a;ic;b}}{2} - \frac{h^2_{ac;ib}}{2} + h^{1de} h^1_{ac;ieb} - \frac{1}{2} h^{1d}_{e;b} h^1_{a;ic} + \frac{1}{2} h^{1d}_{e;a} h^1_{b;ic} + \\
& \frac{1}{2} h^1_{ce;ib} h^1_{a;eid} + \frac{1}{2} h^1_{be;ic} h^1_{a;eid} - \frac{1}{2} h^1_{a;ic} h^1_{be;id} - \frac{1}{2} h^1_{ce;a} h^1_{b;eid} - \frac{1}{2} h^1_{a;eid} h^1_{bc;ie} + \\
& \frac{1}{2} h^1_{a;ic} h^1_{b;ie} + \frac{1}{2} h^1_{e;ib} h^1_{ac;ie} + \frac{1}{2} h^1_{be;id} h^1_{ac;ie} - \frac{1}{2} h^1_{b;ie} h^1_{ac;ie} - \\
& \frac{1}{2} h^1_{ce;ib} h^1_{a;die} - \frac{1}{2} h^1_{be;ic} h^1_{a;die} + \frac{1}{2} h^1_{bc;ie} h^1_{a;die} - \frac{1}{2} h^1_{e;a} h^1_{bc;ie} + \frac{1}{2} h^1_{ce;a} h^1_{b;die}
\end{aligned}$$

---

Applying directly ExpandPerturbation

```
In[104]:=
  Perturbation[RiemannCD[-a, -b, -c, d], 2] // ExpandPerturbation

Out[104]=

$$\frac{1}{2} (-h_{c;b;a}^{2d} - h_{b;c;a}^{2d} + h_{cb;a}^{2d}) + \frac{1}{2} (h_{c;a;b}^{2d} + h_{a;c;b}^{2d} - h_{ca;b}^{2d}) +$$


$$2 \left( \frac{1}{2} h^{1de} (h_{ec;b;a}^1 + h_{eb;c;a}^1 - h_{cb;e;a}^1) + \frac{1}{4} (h_{fc;b}^1 + h_{fb;c}^1 - h_{bc;f}^1) (h_{fd;a}^1 + h_{fa;d}^1 - h_{a;f}^1) \right) -$$


$$2 \left( \frac{1}{2} h^{1de} (h_{ec;a;b}^1 + h_{ea;c;b}^1 - h_{ca;e;b}^1) + \frac{1}{4} (h_{fc;a}^1 + h_{fa;c}^1 - h_{ac;f}^1) (h_{fd;b}^1 + h_{fb;d}^1 - h_{b;f}^1) \right)$$

```

---

Comparison

```
In[105]:=
  %% - % // ToCanonical

0.512031 Second

Out[105]=
0
```

There is a large difference between the timings of both methods.

---

Here we show that the recursive method takes more or less the same time to compute the third-order perturbation as the direct method takes for the seventh-order one:

```
In[106]:=
  ChangeCovD[Perturbation[RiemannCD[-a, -b, -c, d], 3] // RiemannToChristoffel //
  ChristoffelToMetric // ExpandPerturbation, PD, CD] // org;

57.2796 Second

In[107]:=
  ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 7]] // org;

58.1196 Second
```

### 3.5. Perturbations of the Ricci tensor

The Ricci tensor is a simple trace of the Riemann tensor and hence behaves very similarly.

---

First-order perturbation going through the metric

```
In[108]:=
  ChangeCovD[Perturbation[RiemannCD[-a, -c, -b, c] // RiemannToChristoffel //
  ChristoffelToMetric], PD, CD] // ExpandPerturbation // org;

5.69636 Second

Out[108]=

$$-\frac{h_{c;a;b}^{1c}}{2} + \frac{h_{b;a;c}^{1c}}{2} + \frac{h_{a;b;c}^{1c}}{2} - \frac{h_{ab;c}^{1c}}{2}$$

```

---

Applying directly ExpandPerturbation

```
In[109]:=
  Perturbation[RicciCD[-a, -b]] // ExpandPerturbation // org
```

```
Out[109]=
  -  $\frac{h^{1c}_{c;a;b}}{2} + \frac{h^{1c}_{b;a;c}}{2} + \frac{h^{1c}_{a;b;c}}{2} - \frac{h^{1c}_{ab;ic}}{2}$ 
```

---

Comparison

```
In[110]:=
  %% - % // ToCanonical
```

```
Out[110]=
  0
```

---

Second-order perturbation going through the metric

```
In[111]:=
  ChangeCovD[Perturbation[RiemannCD[-a, -c, -b, c] // RiemannToChristoffel //
  ChristoffelToMetric, 2], PD, CD] // ExpandPerturbation // org
```

20.6453 Second

```
Out[111]=
   $h^{1cd} h^{1cd;b;a} + \frac{1}{2} h^{1cd}{}_{;a} h^{1cd;b} - \frac{h^{2c}_{c;a;b}}{2} + \frac{1}{2} h^{1c}{}_{;b;a} h^{1d}{}_{;c} + \frac{1}{2} h^{1c}{}_{;a;b} h^{1d}{}_{;c} +$   

 $\frac{h^{2c}_{b;a;c}}{2} + \frac{h^{2c}_{a;b;c}}{2} - \frac{h^{2ab}{}^{ic}}{2} - \frac{1}{2} h^{1d}{}_{;c} h^{1ab}{}^{ic} - h^{1c}{}_{;b;a} h^{1d}{}_{;c;d} - h^{1c}{}_{;a;b} h^{1d}{}_{;c;d} +$   

 $h^{1ab}{}^{ic} h^{1d}{}_{;c;d} - h^{1cd} h^{1bc;a;d} - h^{1cd} h^{1ac;b;d} + h^{1cd} h^{1ab;c;d} - h^{1bd;c} h^{1c;d} + h^{1bc;d} h^{1c;d}$ 
```

---

Applying directly ExpandPerturbation

```
In[112]:=
  Perturbation[RicciCD[-a, -b], 2] // ExpandPerturbation // org
```

0.240015 Second

```
Out[112]=
   $h^{1cd} h^{1cd;b;a} + \frac{1}{2} h^{1cd}{}_{;a} h^{1cd;b} - \frac{h^{2c}_{c;a;b}}{2} + \frac{1}{2} h^{1c}{}_{;b;a} h^{1d}{}_{;c} + \frac{1}{2} h^{1c}{}_{;a;b} h^{1d}{}_{;c} +$   

 $\frac{h^{2c}_{b;a;c}}{2} + \frac{h^{2c}_{a;b;c}}{2} - \frac{h^{2ab}{}^{ic}}{2} - \frac{1}{2} h^{1d}{}_{;c} h^{1ab}{}^{ic} - h^{1c}{}_{;b;a} h^{1d}{}_{;c;d} - h^{1c}{}_{;a;b} h^{1d}{}_{;c;d} +$   

 $h^{1ab}{}^{ic} h^{1d}{}_{;c;d} - h^{1cd} h^{1bc;a;d} - h^{1cd} h^{1ac;b;d} + h^{1cd} h^{1ab;c;d} - h^{1bd;c} h^{1c;d} + h^{1bc;d} h^{1c;d}$ 
```

---

Comparison

```
In[113]:=
  %% - % // ToCanonical
```

```
Out[113]=
  0
```

Third-order perturbation going through the metric

In[114]:=

```
ChangeCovD[Perturbation[RiemannCD[-a, -c, -b, c] // RiemannToChristoffel //
ChristoffelToMetric, 3], PD, CD] // ExpandPerturbation // org
```

64.388 Second

Out[114]=

$$\begin{aligned}
& \frac{3}{2} h^{2cd} h^1_{cd;ba} - 3 h^1_c{}^e h^1{}^{cd} h^1_{de;ba} + \frac{3}{2} h^1{}^{cd} h^2_{cd;ba} + \frac{3}{4} h^2{}^{cd}{}_{;a} h^1{}_{cd;b} - \\
& 3 h^1{}^{cd} h^1_c{}^e{}_{;a} h^1_{de;b} + \frac{3}{4} h^1{}^{cd}{}_{;a} h^2_{cd;b} - \frac{h^3{}^c{}_{;a;b}}{2} + \frac{3}{4} h^2{}^c{}_{;a} h^1{}^d{}_{;c} + \frac{3}{4} h^2{}^c{}_{;a;b} h^1{}^d{}_{;c} + \\
& \frac{3}{4} h^1{}^c{}_{;a} h^2{}^d{}_{;c} + \frac{3}{4} h^1{}^c{}_{;a;b} h^2{}^d{}_{;c} + \frac{h^3{}^c{}_{;a;b}}{2} + \frac{h^3{}^c{}_{;a;b;c}}{2} - \frac{h^3{}^c{}_{;a;b;c}}{2} - \frac{3}{4} h^2{}^d{}_{;c} h^1{}^{;ic} - \\
& \frac{3}{4} h^1{}^d{}_{;c} h^2{}^{;ic} - 3 h^1{}^{cd} h^1{}^e{}_{;c} h^1{}_{be;d} - \frac{3}{2} h^2{}^c{}_{;a} h^1{}^d{}_{;c} - \frac{3}{2} h^2{}^c{}_{;a;b} h^1{}^d{}_{;c} + \frac{3}{2} h^2{}^{;ic} h^1{}^d{}_{;c} + \\
& 3 h^1{}^{cd} h^1{}^e{}_{;a} h^1{}_{ce;d} + 3 h^1{}^{cd} h^1{}^e{}_{;a;b} h^1{}_{ce;d} - \frac{3}{2} h^1{}^{cd} h^1{}_{bc;a} h^1{}^e{}_{;d} - \frac{3}{2} h^1{}^{cd} h^1{}_{ac;b} h^1{}^e{}_{;d} + \\
& \frac{3}{2} h^1{}^{cd} h^1{}_{ab;c} h^1{}^e{}_{;d} - \frac{3}{2} h^1{}^c{}_{;a} h^2{}^d{}_{;c} - \frac{3}{2} h^1{}^c{}_{;a;b} h^2{}^d{}_{;c} + \frac{3}{2} h^1{}_{ab}{}^{;ic} h^2{}^d{}_{;c} - \frac{3}{2} h^2{}^{cd} h^1{}_{bc;a;d} - \\
& \frac{3}{2} h^1{}^{cd} h^2{}_{bc;a;d} - \frac{3}{2} h^2{}^{cd} h^1{}_{ac;b;d} - \frac{3}{2} h^1{}^{cd} h^2{}_{ac;b;d} + \frac{3}{2} h^2{}^{cd} h^1{}_{ab;c;d} + \frac{3}{2} h^1{}^{cd} h^2{}_{ab;c;d} - \\
& \frac{3}{2} h^2{}_{bd;c} h^1{}^c{}_{;d} + \frac{3}{2} h^2{}_{bc;d} h^1{}^c{}_{;d} - \frac{3}{2} h^1{}_{bd;c} h^2{}^c{}_{;d} + \frac{3}{2} h^1{}_{bc;d} h^2{}^c{}_{;d} + 3 h^1{}^{cd} h^1{}^e{}_{;c} h^1{}_{bd;e} - \\
& \frac{3}{2} h^1{}^{cd} h^1{}^e{}_{;a} h^1{}_{cd;e} - \frac{3}{2} h^1{}^{cd} h^1{}^e{}_{;a;b} h^1{}_{cd;e} + 3 h^1{}^{cd} h^1{}_{bc;a} h^1{}^e{}_{;e} + 3 h^1{}^{cd} h^1{}_{ac;b} h^1{}^e{}_{;e} - \\
& 3 h^1{}^{cd} h^1{}_{ab;c} h^1{}^e{}_{;e} + 3 h^1{}^e{}_{;c} h^1{}^{cd} h^1{}_{bd;a;e} + 3 h^1{}^e{}_{;c} h^1{}^{cd} h^1{}_{ad;b;e} - 3 h^1{}^e{}_{;c} h^1{}^{cd} h^1{}_{ab;d;e} - \\
& 3 h^1{}^{cd} h^1{}_{ce;d} h^1{}_{ab}{}^{;e} + \frac{3}{2} h^1{}^{cd} h^1{}_{cd;e} h^1{}_{ab}{}^{;e} + 3 h^1{}^{cd} h^1{}_{be;d} h^1{}_{ac}{}^{;e} - 3 h^1{}^{cd} h^1{}_{bd;e} h^1{}_{ac}{}^{;e}
\end{aligned}$$



Applying directly ExpandPerturbation

In[115]:=

**Perturbation[RicciCD[-a, -b], 3] // ExpandPerturbation // org**

0.876055 Second

Out[115]=

$$\begin{aligned}
& \frac{3}{2} h^{2cd} h^1_{cd;bia} - 3 h^1_c h^1_{de;bia} h^{1cd} h^1_{de;bia} + \frac{3}{2} h^{1cd} h^2_{cd;bia} + \frac{3}{4} h^{2cd}_{;ia} h^1_{cd;b} - \\
& 3 h^{1cd} h^1_{c;ia} h^1_{de;b} + \frac{3}{4} h^{1cd}_{;ia} h^2_{cd;b} - \frac{h^3_{c;ia;b}}{2} + \frac{3}{4} h^2_{b;ia} h^1_{d;c} + \frac{3}{4} h^2_{a;ib} h^1_{d;c} + \\
& \frac{3}{4} h^1_{b;ia} h^2_{d;c} + \frac{3}{4} h^1_{a;ib} h^2_{d;c} + \frac{h^3_{b;ia;c}}{2} + \frac{h^3_{a;ib;c}}{2} - \frac{h^3_{ab;ic}}{2} - \frac{3}{4} h^2_{d;c} h^1_{ab;ic} - \\
& \frac{3}{4} h^1_{d;c} h^2_{ab;ic} - 3 h^{1cd} h^1_{a;ic} h^1_{be;d} - \frac{3}{2} h^2_{b;ia} h^1_{c;d} - \frac{3}{2} h^2_{a;ib} h^1_{c;d} + \frac{3}{2} h^2_{ab;ic} h^1_{c;d} + \\
& 3 h^{1cd} h^1_{b;ie} h^1_{ce;d} + 3 h^{1cd} h^1_{a;ib} h^1_{ce;d} - \frac{3}{2} h^{1cd} h^1_{bc;ia} h^1_{e;d} - \frac{3}{2} h^{1cd} h^1_{ac;ib} h^1_{e;d} + \\
& \frac{3}{2} h^{1cd} h^1_{ab;ic} h^1_{e;d} - \frac{3}{2} h^1_{b;ia} h^2_{c;d} - \frac{3}{2} h^1_{a;ib} h^2_{c;d} + \frac{3}{2} h^1_{ab;ic} h^2_{c;d} - \frac{3}{2} h^{2cd} h^1_{bc;ia;d} - \\
& \frac{3}{2} h^{1cd} h^2_{bc;ia;d} - \frac{3}{2} h^{2cd} h^1_{ac;ib;d} - \frac{3}{2} h^{1cd} h^2_{ac;ib;d} + \frac{3}{2} h^{2cd} h^1_{ab;ic;d} + \frac{3}{2} h^{1cd} h^2_{ab;ic;d} - \\
& \frac{3}{2} h^2_{bd;ic} h^1_{a;c;d} + \frac{3}{2} h^2_{bc;id} h^1_{a;c;d} - \frac{3}{2} h^1_{bd;ic} h^2_{a;c;d} + \frac{3}{2} h^1_{bc;id} h^2_{a;c;d} + 3 h^{1cd} h^1_{a;ic} h^1_{bd;ie} - \\
& \frac{3}{2} h^{1cd} h^1_{b;ia} h^1_{cd;ie} - \frac{3}{2} h^{1cd} h^1_{a;ib} h^1_{cd;ie} + 3 h^{1cd} h^1_{bc;ia} h^1_{d;ie} + 3 h^{1cd} h^1_{ac;ib} h^1_{d;ie} - \\
& 3 h^{1cd} h^1_{ab;ic} h^1_{d;ie} + 3 h^1_c h^1_{cd} h^1_{bd;ia;ie} + 3 h^1_c h^1_{cd} h^1_{ad;ib;ie} - 3 h^1_c h^1_{cd} h^1_{ab;id;ie} - \\
& 3 h^{1cd} h^1_{ce;id} h^1_{ab;ie} + \frac{3}{2} h^{1cd} h^1_{cd;ie} h^1_{ab;ie} + 3 h^{1cd} h^1_{be;id} h^1_{ac;ie} - 3 h^{1cd} h^1_{bd;ie} h^1_{ac;ie}
\end{aligned}$$

Comparison

In[116]:=

**%% - % // ToCanonical**

0.200013 Second

Out[116]=

0

Let us just compute the fourth-order perturbation making use only of the direct approach. Note that what really takes times is not to construct the perturbation itself but to simplify it.

In[117]:=

**Perturbation[RicciCD[-a, -b], 4] // ExpandPerturbation;**

In[118]:=

% // org

3.1522 Second

Out[118]=

$$\begin{aligned}
& 2 h_{b;ia}^{3cd} h_{cd;bia}^1 - 12 h_{c^e}^{1cd} h_{de;bia}^2 h_{de;bia}^1 + 12 h_c^{1e} h_{cd}^1 h_d^f h_{ef;bia}^1 + 3 h_{cd;bia}^{2cd} h_{cd;bia}^2 - \\
& 6 h_c^{1e} h_{cd}^1 h_{de;bia}^2 + 2 h_{cd}^1 h_{cd;bia}^3 + h_{cd;ia}^3 h_{cd;ib}^1 - 6 h_{cd}^2 h_c^{1e} h_{de;ib}^1 - 6 h_{cd}^1 h_c^{2e} h_{de;ib}^1 h_{de;ib}^1 + \\
& 6 h_{cd}^1 h_{de;ib}^1 h_{ce;ia}^1 h_{df;ib}^1 + 12 h_c^{1e} h_{cd}^1 h_d^f h_{ef;ib}^1 + \frac{3}{2} h_{cd;ia}^{2cd} h_{cd;ib}^2 - 6 h_{cd}^1 h_c^{1e} h_{de;ib}^2 + \\
& h_{cd;ia}^1 h_{cd;ib}^3 - \frac{h_{cd;ia;ib}^4}{2} + h_{b;ia}^3 h_{d;ic}^1 + h_{a;ib}^3 h_{d;ic}^1 + \frac{3}{2} h_{b;ia}^2 h_{d;ic}^2 + \frac{3}{2} h_{a;ib}^2 h_{d;ic}^2 + \\
& h_{b;ia}^1 h_{d;ic}^3 + h_{a;ib}^1 h_{d;ic}^3 + \frac{h_{b;ia;ic}^4}{2} + \frac{h_{a;ib;ic}^4}{2} - \frac{h_{ab;ic}^4}{2} - h_{d;ic}^3 h_{ab;ic}^1 - \frac{3}{2} h_{d;ic}^2 h_{ab;ic}^2 - \\
& h_{d;ic}^1 h_{ab;ic}^3 - 6 h_{cd}^2 h_{a;ic}^1 h_{be;id}^1 - 6 h_{cd}^1 h_{a;ic}^2 h_{be;id}^1 - 2 h_{b;ia}^3 h_{c;id}^1 - 2 h_{a;ib}^3 h_{c;id}^1 + \\
& 2 h_{ab;ic}^3 h_{c;id}^1 + 6 h_{cd}^2 h_{b;ia}^1 h_{ce;id}^1 + 6 h_{cd}^1 h_{b;ia}^2 h_{ce;id}^1 + 6 h_{cd}^2 h_{a;ib}^1 h_{ce;id}^1 + 6 h_{cd}^1 h_{a;ib}^2 h_{ce;id}^1 + \\
& 6 h_{cd}^1 h_{ef}^1 h_{bc;ia}^1 h_{ef;id}^1 + 6 h_{cd}^1 h_{ef}^1 h_{ac;ib}^1 h_{ef;id}^1 - 6 h_{cd}^1 h_{ef}^1 h_{ab;ic}^1 h_{ef;id}^1 - 3 h_{cd}^2 h_{bc;ia}^1 h_{ef;id}^1 - \\
& 3 h_{cd}^1 h_{bc;ia}^2 h_{ef;id}^1 - 3 h_{cd}^2 h_{ac;ib}^1 h_{ef;id}^1 - 3 h_{cd}^1 h_{ac;ib}^2 h_{ef;id}^1 + 3 h_{cd}^2 h_{ab;ic}^1 h_{ef;id}^1 + \\
& 3 h_{cd}^1 h_{ab;ic}^2 h_{ef;id}^1 - 6 h_{cd}^1 h_{a;ic}^1 h_{be;id}^2 - 3 h_{cd}^2 h_{a;ib}^1 h_{c;id}^2 - 3 h_{cd}^1 h_{a;ib}^2 h_{c;id}^2 + 3 h_{ab;ic}^2 h_{c;id}^2 + \\
& 6 h_{cd}^1 h_{a;ib}^1 h_{ce;id}^2 + 6 h_{cd}^1 h_{a;ib}^2 h_{ce;id}^2 - 3 h_{cd}^1 h_{bc;ia}^1 h_{ce;id}^2 - 3 h_{cd}^1 h_{ac;ib}^1 h_{ce;id}^2 + \\
& 3 h_{cd}^1 h_{ab;ic}^1 h_{ce;id}^2 - 2 h_{b;ia}^3 h_{c;id}^2 - 2 h_{a;ib}^3 h_{c;id}^2 + 2 h_{ab;ic}^3 h_{c;id}^2 - 2 h_{bc;ia}^3 h_{c;id}^2 + \\
& 6 h_{cd}^1 h_{c^e}^{2e} h_{be;ia;id}^1 - 3 h_{cd}^2 h_{bc;ia;id}^2 - 2 h_{cd}^1 h_{bc;ia;id}^3 - 2 h_{cd}^3 h_{ac;ib;id}^1 + 6 h_{cd}^1 h_{c^e}^{2e} h_{ae;ib;id}^1 - \\
& 3 h_{cd}^2 h_{ac;ib;id}^2 - 2 h_{cd}^1 h_{ac;ib;id}^3 + 2 h_{cd}^3 h_{ab;ic;id}^1 + 3 h_{cd}^2 h_{ab;ic;id}^2 + 2 h_{cd}^1 h_{ab;ic;id}^3 - \\
& 6 h_{cd}^1 h_{c^e}^{2e} h_{bd;ic}^1 h_{a;id}^2 + 2 h_{bd;ic}^3 h_{a;id}^2 - 3 h_{bd;ic}^2 h_{a;id}^3 + 3 h_{bd;ic}^1 h_{a;id}^3 - \\
& 2 h_{bd;ic}^1 h_{a;id}^3 + 2 h_{bd;ic}^1 h_{a;id}^3 - 12 h_{cd}^1 h_{ef}^1 h_{bf;id}^1 h_{ac;ie}^1 + 6 h_{cd}^2 h_{a;ic}^1 h_{bd;ie}^1 + \\
& 6 h_{cd}^1 h_{a;ic}^2 h_{bd;ie}^1 + 12 h_c^{1e} h_{cd}^1 h_{a;id}^1 h_{bf;ie}^1 - 3 h_{cd}^2 h_{b;ia}^1 h_{cd;ie}^1 - 3 h_{cd}^1 h_{b;ia}^2 h_{cd;ie}^1 - \\
& 3 h_{cd}^2 h_{a;ib}^1 h_{cd;ie}^1 - 3 h_{cd}^1 h_{a;ib}^2 h_{cd;ie}^1 + 6 h_{cd}^2 h_{bc;ia}^1 h_{d;ie}^1 + 6 h_{cd}^1 h_{bc;ia}^2 h_{d;ie}^1 + \\
& 6 h_{cd}^2 h_{ac;ib}^1 h_{d;ie}^1 + 6 h_{cd}^1 h_{ac;ib}^2 h_{d;ie}^1 - 6 h_{cd}^2 h_{ab;ic}^1 h_{d;ie}^1 - 6 h_{cd}^1 h_{ab;ic}^2 h_{d;ie}^1 - \\
& 12 h_c^{1e} h_{cd}^1 h_{b;ia}^1 h_{df;ie}^1 - 12 h_c^{1e} h_{cd}^1 h_{a;ib}^1 h_{df;ie}^1 + 6 h_c^{1e} h_{cd}^1 h_{bd;ia}^1 h_{f;ie}^1 + \\
& 6 h_c^{1e} h_{cd}^1 h_{ad;ib}^1 h_{f;ie}^1 - 6 h_c^{1e} h_{cd}^1 h_{ab;id}^1 h_{f;ie}^1 + 6 h_{cd}^1 h_{a;ic}^1 h_{bd;ie}^2 - 3 h_{cd}^1 h_{b;ia}^1 h_{cd;ie}^2 - \\
& 3 h_{cd}^1 h_{a;ib}^1 h_{cd;ie}^2 + 6 h_{cd}^1 h_{bc;ia}^1 h_{d;ie}^2 + 6 h_{cd}^1 h_{ac;ib}^1 h_{d;ie}^2 - 6 h_{cd}^1 h_{ab;ic}^1 h_{d;ie}^2 + \\
& 6 h_{cd}^1 h_{c^e}^{2e} h_{bd;ia;ie}^1 + 6 h_c^{1e} h_{cd}^1 h_{bd;ia;ie}^2 + 6 h_{cd}^1 h_{c^e}^{2e} h_{ad;ib;ie}^1 + 6 h_c^{1e} h_{cd}^1 h_{ad;ib;ie}^2 - \\
& 6 h_{cd}^1 h_{c^e}^{2e} h_{ab;id;ie}^1 - 6 h_c^{1e} h_{cd}^1 h_{ab;id;ie}^2 - 6 h_{cd}^2 h_{ce;id}^1 h_{ab;ie}^1 - 6 h_{cd}^1 h_{ce;id}^2 h_{ab;ie}^1 + \\
& 3 h_{cd}^2 h_{cd;ie}^1 h_{ab;ie}^1 + 3 h_{cd}^1 h_{cd;ie}^2 h_{ab;ie}^1 + 6 h_{cd}^2 h_{be;id}^1 h_{ac;ie}^1 + 6 h_{cd}^1 h_{be;id}^2 h_{ac;ie}^1 - \\
& 6 h_{cd}^2 h_{bd;ie}^1 h_{ac;ie}^1 - 6 h_{cd}^1 h_{bd;ie}^2 h_{ac;ie}^1 - 6 h_{cd}^1 h_{ce;id}^1 h_{ab;ie}^2 + 3 h_{cd}^1 h_{cd;ie}^1 h_{ab;ie}^2 + \\
& 6 h_{cd}^1 h_{be;id}^1 h_{ac;ie}^2 - 6 h_{cd}^1 h_{bd;ie}^1 h_{ac;ie}^2 + 12 h_{cd}^1 h_{ef}^1 h_{ac;ie}^1 h_{bd;if}^1 - 12 h_c^{1e} h_{cd}^1 h_{a;id}^1 h_{be;if}^1 - \\
& 12 h_{cd}^1 h_{ef}^1 h_{bc;ia}^1 h_{de;if}^1 + 6 h_c^{1e} h_{cd}^1 h_{b;ia}^1 h_{de;if}^1 - 12 h_{cd}^1 h_{ef}^1 h_{ac;ib}^1 h_{de;if}^1 + \\
& 6 h_c^{1e} h_{cd}^1 h_{a;ib}^1 h_{de;if}^1 + 12 h_{cd}^1 h_{ef}^1 h_{ab;ic}^1 h_{de;if}^1 - 12 h_c^{1e} h_{cd}^1 h_{bd;ia}^1 h_{e;if}^1 - \\
& 12 h_c^{1e} h_{cd}^1 h_{ad;ib}^1 h_{e;if}^1 + 12 h_c^{1e} h_{cd}^1 h_{ab;id}^1 h_{e;if}^1 - 12 h_c^{1e} h_{cd}^1 h_{d;f}^1 h_{be;ia;if}^1 - \\
& 12 h_c^{1e} h_{cd}^1 h_{d;f}^1 h_{ae;ib;if}^1 + 12 h_c^{1e} h_{cd}^1 h_{d;f}^1 h_{ab;ie;if}^1 + 12 h_c^{1e} h_{cd}^1 h_{d;f}^1 h_{df;ie}^1 h_{ab;if}^1 - \\
& 6 h_c^{1e} h_{cd}^1 h_{d;f}^1 h_{de;if}^1 h_{ab;if}^1 - 12 h_c^{1e} h_{cd}^1 h_{bf;ie}^1 h_{ad;if}^1 + 12 h_c^{1e} h_{cd}^1 h_{be;if}^1 h_{ad;if}^1
\end{aligned}$$

Some more examples:

In[119]:=

**Perturbation[RicciCD[-a, -b], 8] // ExpandPerturbation;**

1.39609 Second

In[120]:=

**Perturbation[RicciCD[-a, -b], 12] // ExpandPerturbation;**

38.4104 Second

### 3.6. Perturbations of the Ricci scalar

The Ricci scalar is computed from a double trace of the Riemann tensor.

```
In[121]:=
ricciscalar =
g[a, b] RiemannCD[-a, -c, -b, c] // RiemannToChristoffel // ChristoffelToMetric //
Simplification
0.624039 Second

Out[121]=
- $\frac{1}{4} g^{ab} g^{cd}$ 
(-4 gac,b,d + 4 gab,c,d + gef (2 gbf,d gac,e - 3 gac,e gbd,f + gab,c (gef,d - 4 gde,f) + 4 gac,b gde,f))
```

---

First-order perturbation from the metric

```
In[122]:=
ChangeCovD[Perturbation[ricciscalar] // ExpandPerturbation, PD, CD] // org
1.89212 Second

Out[122]=
- $\Gamma^{abc} \Gamma_{ba}^d h^1_{cd} + \Gamma_a^b \Gamma_b^{cd} h^1_{cd} + h^{1ab} \Gamma^c_{ac;b} + h^{1ab}{}_{;a;b} - h^{1a}{}_{;b}{}^b - h^{1ab} \Gamma^c_{ab;c}$ 
```

---

Applying directly ExpandPerturbation

```
In[123]:=
Perturbation[RicciscalarCD[]] // ExpandPerturbation // org

Out[123]=
-h1ab Rab + h1ab;a;b - h1a;bb
```

---

Comparison

```
In[124]:=
ChangeCovD[% - %% // RiemannToChristoffel, PD, CD] // Simplification

Out[124]=
0
```

---

Second-order perturbation from the metric

```
In[125]:=
ChangeCovD[Perturbation[ricciscalar, 2] // ExpandPerturbation, PD, CD] // org
7.17245 Second

Out[125]=
2  $\Gamma^{abc} \Gamma_{ba}^d h^1_c{}^e h^1_{de} - 2 \Gamma_a^b \Gamma_b^{cd} h^1_c{}^e h^1_{de} - \Gamma^{abc} \Gamma_{ba}^d h^2_{cd} +$ 
 $\Gamma_a^b \Gamma_b^{cd} h^2_{cd} + h^{2ab} \Gamma^c_{ac;b} + 2 h^{1ab} h^1_{c;a;b} + h^{2ab}{}_{;a;b} - h^{2a}{}_{;b}{}^b - 2 h^{1ab} h^1_{a;c;b} -$ 
 $\frac{1}{2} h^1_{c;b} h^1_a{}^b - h^{2ab} \Gamma^c_{ab;c} - 2 h^1_a{}^c h^1_{ab} \Gamma^d_{bd;c} - 2 h^{1ab}{}_{;a} h^1_b{}^c + 2 h^1_a{}^b h^1_{b;c} -$ 
 $2 h^{1ab} h^1_{a;b;c} + 2 h^{1ab} h^1_{ab}{}^c{}_{;c} - h^1_{ac;b} h^1_{ab;c} + \frac{3}{2} h^1_{ab;c} h^1_{ab;c} + 2 h^1_a{}^c h^1_{ab} \Gamma^d_{bc;d}$ 
```

---

Applying directly ExpandPerturbation

```
In[126]:=
  Perturbation[RicciScalarCD[], 2] // ExpandPerturbation // org
0.368023 Second
```

```
Out[126]=
-h2ab Rab + 2 h1ca h1ab Rbc + 2 h1ab h1cc;a;b + h2ab;a;b - h2aa;b - 2 h1ab h1c;c;b -  $\frac{1}{2}$  h1cc;b h1aa;b -
2 h1ab;a h1cb;c + 2 h1aa;b h1cb;c - 2 h1ab h1ca;b;c + 2 h1ab h1cab;c - h1ac;b h1ab;c +  $\frac{3}{2}$  h1ab;c h1ab;c
```

---

Comparison

```
In[127]:=
  ChangeCovD[% - %% // RiemannToChristoffel, PD, CD] // Simplification
0.208013 Second
```

```
Out[127]=
0
```

---

Third-order perturbation iteratively from metric perturbations

```
In[128]:=
  ChangeCovD[Perturbation[ricciscalar, 3] // ExpandPerturbation, PD, CD] // org
27.6217 Second
```

```
Out[128]=
-6 Γabc Γbad h1ce h1fd h1ef + 6 Γaab Γbcd h1ce h1fd h1ef + 6 Γabc Γbad h1ce h2de - 6 Γaab Γbcd h1ce h2de -
Γabc Γbad h3cd + Γaab Γbcd h3cd + h3ab Γcac;b - 3 h1ab h2ca Γdcd;b -  $\frac{9}{2}$  h1ab h1cd;a h1cd;b +
 $\frac{3}{2}$  h1ab h1c;a h1dd;b + 3 h1ab;a h2cc;b + 3 h2ab h1c;a;b + 3 h1ab h2cc;a;b + h3ab;a;b - h3aa;b -
3 h2ab h1c;c;b - 3 h1ab h2c;c;b -  $\frac{3}{2}$  h2cc;b h1aa;b - h3ab Γcab;c - 3 h1ab h2ca Γdbd;c -
6 h1ab h1dd;b h1ca;c - 6 h1ab h1ca;b h1dd;c - 6 h1ab;a h2cb;c + 3 h1aa;b h2cb;c - 3 h2ab h1ca;b;c -
6 h1ca h1ab h1dd;b;c - 3 h1ab h2ca;b;c + 3 h2ab h1cab;c + 3 h1ab h2cab;c + 6 h1ca h1ab h1db;d;c +
3 h1ab h1dd;c h1cab - 3 h2ac;b h1ab;c +  $\frac{9}{2}$  h2ab;c h1ab;c + 6 h1ab h2ca Γdbc;d +
6 h1ca h1ab h1db Γece;d + 6 h1ab h1ca;c h1db;d + 12 h1ab h1ca;b h1dc;d - 6 h1ab h1cab h1dc;d +
6 h1ab h1cd h1ac;b;d - 6 h1ab h1cd h1ab;c;d + 6 h1ca h1ab h1db;c;d - 6 h1ca h1ab h1dbc;d +
6 h1ab h1cd;b h1ca;d + 3 h1ab h1bd;c h1ca - 9 h1ab h1bc;d h1ca;d - 6 h1ca h1ab h1db Γecd;e
```

---

Applying directly ExpandPerturbation

```
In[129]:=
  Perturbation[RicciScalarCD[], 3] // ExpandPerturbation // org
1.5401 Second
```

```
Out[129]=
-h3aab Rab + 6 h1aab h2ac Rbc - 6 h1ac h1bab h1bd Rcd -  $\frac{9}{2}$  h1aab h1acd h1cdb +
 $\frac{3}{2}$  h1aab h1ca h1db + 3 h1aab h2cb + 3 h2ab h1ca h1b + 3 h1aab h2ca h1b + h3aab -
h3aab h1b - 3 h2ab h1ac h1cb - 3 h1aab h2ca h1cb -  $\frac{3}{2}$  h2cb h1aab - 6 h1aab h1db h1ac -
6 h1aab h1ac h1db - 6 h1aab h2bc h1c + 3 h1aab h2bc h1c - 3 h2ab h1ac h1bc - 6 h1ac h1aab h1db h1c -
3 h1aab h2ab h1c + 3 h2ab h1abc + 3 h1ab h2abc + 6 h1ac h1bd h1c + 3 h1ab h1dc h1abc -
3 h2acb h1abc +  $\frac{9}{2}$  h2abc h1abc + 6 h1aab h1ac h1bd + 12 h1aab h1ac h1cd -
6 h1aab h1abc h1cd + 6 h1aab h1acb h1bdc - 6 h1aab h1abc h1cdd + 6 h1ac h1aab h1bc h1d -
6 h1ac h1ab h1bcd + 6 h1aab h1cdb h1ac + 3 h1ab h1bdc h1ac - 9 h1ab h1bcd h1ac
```

---

Comparison

```
In[130]:=
  ChangeCovD[% - % // RiemannToChristoffel, PD, CD] // org
0.456029 Second
```

```
Out[130]=
3 h1aab h2ac  $\Gamma$ cdb - 3 h1aab h2ac  $\Gamma$ bdc
```

---

That expression is indeed only zero for a Levi–Civita connection. To show it we change back to partial derivatives and expand the Christoffels into derivatives of the metric:

```
In[131]:=
  ChangeCovD[% , CD, PD] // ToCanonical

Out[131]=
3 h1aab h2ac  $\Gamma$ cdb - 3 h1aab h2ac  $\Gamma$ bdc

In[132]:=
  % // ChristoffelToMetric // ToCanonical

0.228014 Second
```

```
Out[132]=
0
```

---

Fourth–order perturbation iteratively from metric perturbations

```
In[133]:=
  ChangeCovD[Perturbation[ricciscalar, 4] // ExpandPerturbation, PD, CD] // org;
83.7492 Second
```

---

Applying directly ExpandPerturbation

```
In[134]:=
  Perturbation[RicciScalarCD[], 4] // ExpandPerturbation // org;
5.08832 Second
```

---

Comparison

```
In[135]:=
  ChangeCovD[% - % // ToCanonical // RiemannToChristoffel, PD, CD] // Simplification
1.38009 Second
```

```
Out[135]=
4 h1ab (h3ac (Γdcd;b - Γdbd;c) + 3 h1ac h2bd (-Γede;c + Γece;d))
```

---

Again, this is zero for a Levi-Civita connection (we do not use the metric in the index-canonicalization process to avoid problems with the partial derivatives):

```
In[136]:=
  Simplification[ChangeCovD[%, CD, PD], UseMetricOnVBundle → None]
0.108007 Second
```

```
Out[136]=
4 h1ab (h3ac (Γdcd,b - Γdbd,c) + 3 h1ac h2bd (-Γede,c + Γece;d))
```

```
In[137]:=
  Simplification[% // ChristoffelToMetric, UseMetricOnVBundle → None]
0.440028 Second
```

```
Out[137]=
0
```

### 3.7. Perturbations of the Einstein tensor

To help in the recursive computation, we define the Einstein tensor in terms of the Christoffels instead of the metric.

```
In[138]:=
  einstein = EinsteinCD[-a, -b] // EinsteinToRicci // RiemannToChristoffel // Simplify
```

```
Out[138]=
Γccd Γdab - Γcad Γdcb - Γccb,a + Γcab,c -  $\frac{1}{2}$  gba gcd (Γeef Γfcd - Γecf Γfed - Γeed,c + Γecd,e)
```

The first-order perturbation from the Christoffel perturbations:

In[139]:=

**Perturbation[einstein] // ExpandPerturbation**

Out[139]=

$$\begin{aligned}
& -\frac{1}{2} \Gamma_{cb}^d (h_{d;a}^c - h_{ad}^{1c} + h_{a;d}^c) + \frac{1}{2} \Gamma_{ab}^d (h_{d;c}^c - h_{cd}^{1c} + h_{c;d}^c) + \\
& \frac{1}{2} \Gamma_{cd}^c (h_{b;a}^d + h_{a;b}^d - h_{ab}^{1d}) - \frac{1}{2} \Gamma_{ad}^c (h_{c;b}^d + h_{b;c}^d - h_{cb}^{1d}) + \\
& \frac{1}{2} (-h_{c;b,a}^c - h_{b;c,a}^c + h_{cb}^{1c},a) + \frac{1}{2} (h_{b;a,c}^c + h_{a;b,c}^c - h_{ab}^{1c},c) + \\
& \frac{1}{2} \left( -g^{cd} h_{ba}^{1c} (\Gamma_{ef}^e \Gamma_{cd}^f - \Gamma_{cf}^e \Gamma_{ed}^f - \Gamma_{ed,c}^e + \Gamma_{cd,e}^e) + g_{ba} h^{1cd} (\Gamma_{ef}^e \Gamma_{cd}^f - \Gamma_{cf}^e \Gamma_{ed}^f - \Gamma_{ed,c}^e + \Gamma_{cd,e}^e) \right. \\
& g_{ba} g^{cd} \left( -\frac{1}{2} \Gamma_{ed}^f (h_{f;c}^e - h_{cf}^{1e} + h_{c;f}^e) + \frac{1}{2} \Gamma_{cd}^f (h_{f;e}^e - h_{ef}^{1e} + h_{e;f}^e) + \right. \\
& \left. \frac{1}{2} \Gamma_{ef}^e (h_{d;c}^f + h_{c;d}^f - h_{cd}^{1f}) - \frac{1}{2} \Gamma_{cf}^e (h_{e;d}^f + h_{d;e}^f - h_{ed}^{1f}) + \right. \\
& \left. \frac{1}{2} (-h_{e;d,c}^e - h_{d;e,c}^e + h_{ed}^{1e},c) + \frac{1}{2} (h_{d;c,e}^e + h_{c;d,e}^e - h_{cd}^{1e},e) \right) \Big)
\end{aligned}$$

Applying directly ExpandPerturbation

In[140]:=

**Perturbation[EinsteinCD[-a, -b]] // ExpandPerturbation**

Out[140]=

$$\begin{aligned}
& \frac{1}{2} (-h_{c;b;a}^c - h_{b;c;a}^c + h_{bc}^{1c},a) + \frac{1}{2} (h_{b;a;c}^c + h_{a;b;c}^c - h_{ba}^{1c},c) + \frac{1}{2} \left( -h_{ab}^1 R - \right. \\
& \left. g_{ab} \left( -h^{1cd} R_{cd} + g^{cd} \left( \frac{1}{2} (-h_{e;d;c}^e - h_{d;e;c}^e + h_{de}^{1e},c) + \frac{1}{2} (h_{d;c;e}^e + h_{c;d;e}^e - h_{dc}^{1e},e) \right) \right) \right)
\end{aligned}$$

Comparison

In[141]:=

**ChangeCovD[%% - % // RiemannToChristoffel, PD, CD] // ToCanonical**

0.580037 Second

Out[141]=

0

The second-order perturbation from Christoffel perturbations:

In[142]:=

**Perturbation[einstein, 2] // ExpandPerturbation**

Out[142]=

$$\begin{aligned}
& \frac{1}{2} (h^1_{d;c} - h^1_{cd}{}^{;c} + h^1_{c;d}) (h^1_{b;a} + h^1_{a;b} - h^1_{ab}{}^{;d}) - \\
& \frac{1}{2} (h^1_{d;a} - h^1_{ad}{}^{;c} + h^1_{a;d}) (h^1_{c;b} + h^1_{b;c} - h^1_{cb}{}^{;d}) + \\
& \Gamma^c_{cd} \left( \frac{1}{2} (h^2_{b;a} + h^2_{a;b} - h^2_{ab}{}^{;d}) - h^1_{de} (h^1_{eb;a} + h^1_{ea;b} - h^1_{ab;e}) \right) - \\
& \Gamma^d_{cb} \left( \frac{1}{2} (h^2_{d;a} - h^2_{ad}{}^{;c} + h^2_{a;d}) - h^1_{ce} (h^1_{ed;a} + h^1_{ea;d} - h^1_{ad;e}) \right) - \\
& \Gamma^c_{ad} \left( \frac{1}{2} (h^2_{c;b} + h^2_{b;c} - h^2_{cb}{}^{;d}) - h^1_{de} (h^1_{ec;b} + h^1_{eb;c} - h^1_{cb;e}) \right) + \\
& \Gamma^d_{ab} \left( \frac{1}{2} (h^2_{d;c} - h^2_{cd}{}^{;c} + h^2_{c;d}) - h^1_{ce} (h^1_{ed;c} + h^1_{ec;d} - h^1_{cd;e}) \right) + \\
& (h^1_{dc;b} + h^1_{db;c} - h^1_{cb;d}) h^1_{cd,a} + \frac{1}{2} (-h^2_{c;b,a} - h^2_{b;c,a} + h^2_{cb}{}^{;c,a}) + \\
& h^1_{cd} (h^1_{dc;b,a} + h^1_{db;c,a} - h^1_{cb;d,a}) - (h^1_{db;a} + h^1_{da;b} - h^1_{ab;d}) h^1_{cd,c} + \\
& \frac{1}{2} (h^2_{b;a,c} + h^2_{a;b,c} - h^2_{ab}{}^{;c,c}) - h^1_{cd} (h^1_{db;a,c} + h^1_{da;b,c} - h^1_{ab;d,c}) + \frac{1}{2} \\
& \left( 2 h^1_{ba} h^1_{cd} (\Gamma^e_{ef} \Gamma^f_{cd} - \Gamma^e_{cf} \Gamma^f_{ed} - \Gamma^e_{ed,c} + \Gamma^e_{cd,e}) - g^{cd} h^2_{ba} (\Gamma^e_{ef} \Gamma^f_{cd} - \Gamma^e_{cf} \Gamma^f_{ed} - \Gamma^e_{ed,c} + \Gamma^e_{cd,e}) \right) \\
& g_{ba} (2 h^1_{cf1} h^1_{f1}{}^d - h^2_{cd}) (\Gamma^e_{ef} \Gamma^f_{cd} - \Gamma^e_{cf} \Gamma^f_{ed} - \Gamma^e_{ed,c} + \Gamma^e_{cd,e}) - \\
& 2 g^{cd} h^1_{ba} \left( -\frac{1}{2} \Gamma^f_{ed} (h^1_{f;c} - h^1_{cf}{}^{;e} + h^1_{c;f}) + \frac{1}{2} \Gamma^f_{cd} (h^1_{f;e} - h^1_{ef}{}^{;e} + h^1_{e;f}) + \right. \\
& \quad \frac{1}{2} \Gamma^e_{ef} (h^1_{d;c} + h^1_{c;d} - h^1_{cd}{}^{;f}) - \frac{1}{2} \Gamma^e_{cf} (h^1_{e;d} + h^1_{d;e} - h^1_{ed}{}^{;f}) + \\
& \quad \left. \frac{1}{2} (-h^1_{e;d,c} - h^1_{d;e,c} + h^1_{ed}{}^{;e,c}) + \frac{1}{2} (h^1_{d;c,e} + h^1_{c;d,e} - h^1_{cd}{}^{;e,e}) \right) + \\
& 2 g_{ba} h^1_{cd} \left( -\frac{1}{2} \Gamma^f_{ed} (h^1_{f;c} - h^1_{cf}{}^{;e} + h^1_{c;f}) + \frac{1}{2} \Gamma^f_{cd} (h^1_{f;e} - h^1_{ef}{}^{;e} + h^1_{e;f}) + \right. \\
& \quad \frac{1}{2} \Gamma^e_{ef} (h^1_{d;c} + h^1_{c;d} - h^1_{cd}{}^{;f}) - \frac{1}{2} \Gamma^e_{cf} (h^1_{e;d} + h^1_{d;e} - h^1_{ed}{}^{;f}) + \\
& \quad \left. \frac{1}{2} (-h^1_{e;d,c} - h^1_{d;e,c} + h^1_{ed}{}^{;e,c}) + \frac{1}{2} (h^1_{d;c,e} + h^1_{c;d,e} - h^1_{cd}{}^{;e,e}) \right) - g_{ba} g^{cd} \\
& \left( \frac{1}{2} (h^1_{f;e} - h^1_{ef}{}^{;e} + h^1_{e;f}) (h^1_{d;c} + h^1_{c;d} - h^1_{cd}{}^{;f}) - \frac{1}{2} (h^1_{f;c} - h^1_{cf}{}^{;e} + h^1_{c;f}) (h^1_{e;d} + \right. \\
& \quad \left. h^1_{d;e} - h^1_{ed}{}^{;f}) - \Gamma^f_{ed} \left( \frac{1}{2} (h^2_{f;c} - h^2_{cf}{}^{;e} + h^2_{c;f}) - h^{1ef10} (h^1_{f10f;c} + h^1_{f10c;f} - h^1_{cf;f10}) \right) \right. \\
& \quad \Gamma^f_{cd} \left( \frac{1}{2} (h^2_{f;e} - h^2_{ef}{}^{;e} + h^2_{e;f}) - h^{1ef11} (h^1_{f11f;e} + h^1_{f11e;f} - h^1_{ef;f11}) \right) + \\
& \quad \Gamma^e_{ef} \left( \frac{1}{2} (h^2_{d;c} + h^2_{c;d} - h^2_{cd}{}^{;f}) - h^{1ff12} (h^1_{f12d;c} + h^1_{f12c;d} - h^1_{cd;f12}) \right) - \\
& \quad \Gamma^e_{cf} \left( \frac{1}{2} (h^2_{e;d} + h^2_{d;e} - h^2_{ed}{}^{;f}) - h^{1ff13} (h^1_{f13e;d} + h^1_{f13d;e} - h^1_{ed;f13}) \right) + \\
& \quad (h^1_{f14e;d} + h^1_{f14d;e} - h^1_{ed;f14}) h^{1ef14,c} + \frac{1}{2} (-h^2_{e;d,c} - h^2_{d;e,c} + h^2_{ed}{}^{;e,c}) + \\
& \quad h^{1ef14} (h^1_{f14e;d,c} + h^1_{f14d;e,c} - h^1_{ed;f14,c}) - (h^1_{f15d;c} + h^1_{f15c;d} - h^1_{cd;f15}) h^{1ef15,e} + \\
& \quad \left. \frac{1}{2} (h^2_{d;c,e} + h^2_{c;d,e} - h^2_{cd}{}^{;e,e}) - h^{1ef15} (h^1_{f15d;c,e} + h^1_{f15c;d,e} - h^1_{cd;f15,e}) \right) \Big)
\end{aligned}$$



Applying directly `ExpandPerturbation`

`In[143]:=`

`Perturbation[EinsteinCD[-a, -b], 2] // ExpandPerturbation`

`Out[143]=`

$$\begin{aligned} & \frac{1}{2} (-h^2{}^c{}_{c;b;a} - h^2{}^c{}_{b;c;a} + h^2{}^c{}_{bc}{}^c{}_{;a}) + \frac{1}{2} (h^2{}^c{}_{b;a;c} + h^2{}^c{}_{a;b;c} - h^2{}^c{}_{ba}{}^c{}_{;c}) + \\ & 2 \left( \frac{1}{2} h^{1cd} (h^1{}_{dc;b;a} + h^1{}_{db;c;a} - h^1{}_{bc;d;a}) + \frac{1}{4} (h^1{}_{ec;b} + h^1{}_{eb;c} - h^1{}_{cb;e}) (h^1{}_{;a}{}^{ec} + h^1{}^e{}_a{}^{;c} - h^1{}^c{}_a{}^{;e}) \right) - \\ & 2 \left( \frac{1}{2} h^{1cd} (h^1{}_{db;a;c} + h^1{}_{da;b;c} - h^1{}_{ba;d;c}) + \frac{1}{4} (h^1{}_{eb;a} + h^1{}_{ea;b} - h^1{}_{ab;e}) (h^1{}_{;c}{}^{ec} + h^1{}^e{}_c{}^{;c} - h^1{}^c{}_c{}^{;e}) \right) + \\ & \frac{1}{2} \left( -h^2{}_{ab} R - g_{ab} \left( (2 h^1{}^{ce} h^1{}^d{}_e - h^2{}^{cd}) R_{cd} - \right. \right. \\ & \quad 2 h^{1cd} \left( \frac{1}{2} (-h^1{}^f{}_{f;d;c} - h^1{}^f{}_{d;f;c} + h^1{}_{df}{}^f{}_{;c}) + \frac{1}{2} (h^1{}^f{}_{d;c;f} + h^1{}^f{}_{c;d;f} - h^1{}_{dc}{}^f{}_{;f}) \right) + \\ & \quad g^{cd} \left( \frac{1}{2} (-h^2{}^{f1}{}_{f1;d;c} - h^2{}^{f1}{}_{d;f1;c} + h^2{}_{df1}{}^f{}_{;c}) + \frac{1}{2} (h^2{}^{f1}{}_{d;c;f1} + h^2{}^{f1}{}_{c;d;f1} - h^2{}_{dc}{}^{f1}{}_{;f1}) + \right. \\ & \quad 2 \left( \frac{1}{2} h^{1f1f10} (h^1{}_{f10f1;d;c} + h^1{}_{f10d;f1;c} - h^1{}_{df1;f10;c}) + \right. \\ & \quad \quad \left. \frac{1}{4} (h^1{}_{f11f1;d} + h^1{}_{f11d;f1} - h^1{}_{f1d;f11}) (h^1{}_{;c}{}^{f11f1} + h^1{}^{f11}{}_c{}^{;f1} - h^1{}^{f1}{}_c{}^{;f11}) \right) - \\ & \quad 2 \left( \frac{1}{2} h^{1f1f10} (h^1{}_{f10d;c;f1} + h^1{}_{f10c;d;f1} - h^1{}_{dc;f10;f1}) + \right. \\ & \quad \quad \left. \frac{1}{4} (h^1{}_{f11d;c} + h^1{}_{f11c;d} - h^1{}_{cd;f11}) (h^1{}_{;f1}{}^{f11f1} + h^1{}^{f11}{}_{f1}{}^{;f1} - h^1{}^{f1}{}_{f1}{}^{;f11}) \right) \left. \right) - \\ & 2 h^1{}_{ab} \left( -h^{1f12f13} R_{f12f13} + g^{f12f13} \left( \frac{1}{2} (-h^{1f14}{}_{f14;f13;f12} - h^{1f14}{}_{f13;f14;f12} + h^{1f14}{}_{f13f14}{}^{;f14}{}_{;f12}) + \right. \right. \\ & \quad \left. \left. \frac{1}{2} (h^{1f14}{}_{f13;f12;f14} + h^{1f14}{}_{f12;f13;f14} - h^{1f14}{}_{f13f12}{}^{;f14}{}_{;f14}) \right) \right) \end{aligned}$$

Comparison

`In[144]:=`

`ChangeCovD[% - % // RiemannToChristoffel, PD, CD] // ToCanonical`

3.45622 Second

`Out[144]=`

0

Note that `ExpandPerturbation` is really fast. The seventh-order perturbation of the Einstein tensor is computed in just a few seconds.

`In[145]:=`

`Perturbation[EinsteinCD[-a, -b], 7] // ExpandPerturbation;`

2.08413 Second

### 3.8. Perturbations of the Weyl tensor

The Weyl tensor is given here in terms of Riemann, Ricci and the Ricci scalar. Giving it in terms of the metric would be too slow.

Weyl in terms of Riemann:

In[146]:=

**WeylCD[-a, -b, -c, -d] // WeylToRiemann**

Out[146]=

$$-\frac{1}{2} g_{db} R_{ac} + \frac{1}{2} g_{cb} R_{ad} + \frac{1}{2} g_{da} R_{bc} - \frac{1}{2} g_{ca} R_{bd} - \frac{1}{6} g_{cb} g_{da} R + \frac{1}{6} g_{ca} g_{db} R + R_{abcd}$$

In[147]:=

**Perturbation[%] // ExpandPerturbation**

Out[147]=

$$\begin{aligned} & h^1_{ed} R_{abc}{}^e + g_{ed} \left( \frac{1}{2} (-h^1_{c;b;a} - h^1_{b;c;a} + h^1_{cb}{}^{;e}{}_{;a}) + \frac{1}{2} (h^1_{c;a;b} + h^1_{a;c;b} - h^1_{ca}{}^{;e}{}_{;b}) \right) + \\ & \frac{1}{2} \left( -h^1_{db} R_{ac} - g_{db} \left( \frac{1}{2} (-h^1_{e;c;a} - h^1_{c;e;a} + h^1_{ce}{}^{;e}{}_{;a}) + \frac{1}{2} (h^1_{c;a;e} + h^1_{a;c;e} - h^1_{ca}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{2} \left( h^1_{da} R_{bc} + g_{da} \left( \frac{1}{2} (-h^1_{e;c;b} - h^1_{c;e;b} + h^1_{ce}{}^{;e}{}_{;b}) + \frac{1}{2} (h^1_{c;b;e} + h^1_{b;c;e} - h^1_{cb}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{2} \left( h^1_{cb} R_{ad} + g_{cb} \left( \frac{1}{2} (-h^1_{e;d;a} - h^1_{d;e;a} + h^1_{de}{}^{;e}{}_{;a}) + \frac{1}{2} (h^1_{d;a;e} + h^1_{a;d;e} - h^1_{da}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{2} \left( -h^1_{ca} R_{bd} - g_{ca} \left( \frac{1}{2} (-h^1_{e;d;b} - h^1_{d;e;b} + h^1_{de}{}^{;e}{}_{;b}) + \frac{1}{2} (h^1_{d;b;e} + h^1_{b;d;e} - h^1_{db}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{6} \left( -g_{da} h^1_{cb} R - g_{cb} h^1_{da} R - g_{cb} g_{da} \left( -h^{1ef} R_{ef} + \right. \right. \\ & \quad \left. \left. g^{ef} \left( \frac{1}{2} (-h^{1f1}_{f1;f;e} - h^{1f1}_{f;f1;e} + h^{1f1}_{ff1}{}^{;f1}{}_{;e}) + \frac{1}{2} (h^{1f1}_{f;e;f1} + h^{1f1}_{e;f;f1} - h^{1fe}{}^{;f1}{}_{;f1}) \right) \right) \right) + \\ & \frac{1}{6} \left( g_{db} h^1_{ca} R + g_{ca} h^1_{db} R + g_{ca} g_{db} \left( -h^{1ef} R_{ef} + \right. \right. \\ & \quad \left. \left. g^{ef} \left( \frac{1}{2} (-h^{1f1}_{f1;f;e} - h^{1f1}_{f;f1;e} + h^{1f1}_{ff1}{}^{;f1}{}_{;e}) + \frac{1}{2} (h^{1f1}_{f;e;f1} + h^{1f1}_{e;f;f1} - h^{1fe}{}^{;f1}{}_{;f1}) \right) \right) \right) \end{aligned}$$

The first-order perturbation applying directly `ExpandPerturbation` is also internally computed along the same lines:

In[148]:=

**Perturbation[WeylCD[-a, -b, -c, -d], 1] // ExpandPerturbation**

Out[148]=

$$\begin{aligned} & h^1_{ed} R_{abc}{}^e + g_{ed} \left( \frac{1}{2} (-h^1_{c;b;a} - h^1_{b;c;a} + h^1_{cb}{}^{;e}{}_{;a}) + \frac{1}{2} (h^1_{c;a;b} + h^1_{a;c;b} - h^1_{ca}{}^{;e}{}_{;b}) \right) + \\ & \frac{1}{2} \left( -h^1_{db} R_{ac} - g_{db} \left( \frac{1}{2} (-h^1_{e;c;a} - h^1_{c;e;a} + h^1_{ce}{}^{;e}{}_{;a}) + \frac{1}{2} (h^1_{c;a;e} + h^1_{a;c;e} - h^1_{ca}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{2} \left( h^1_{da} R_{bc} + g_{da} \left( \frac{1}{2} (-h^1_{e;c;b} - h^1_{c;e;b} + h^1_{ce}{}^{;e}{}_{;b}) + \frac{1}{2} (h^1_{c;b;e} + h^1_{b;c;e} - h^1_{cb}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{2} \left( h^1_{cb} R_{ad} + g_{cb} \left( \frac{1}{2} (-h^1_{e;d;a} - h^1_{d;e;a} + h^1_{de}{}^{;e}{}_{;a}) + \frac{1}{2} (h^1_{d;a;e} + h^1_{a;d;e} - h^1_{da}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{2} \left( -h^1_{ca} R_{bd} - g_{ca} \left( \frac{1}{2} (-h^1_{e;d;b} - h^1_{d;e;b} + h^1_{de}{}^{;e}{}_{;b}) + \frac{1}{2} (h^1_{d;b;e} + h^1_{b;d;e} - h^1_{db}{}^{;e}{}_{;e}) \right) \right) + \\ & \frac{1}{6} \left( -g_{da} h^1_{cb} R - g_{cb} h^1_{da} R - g_{cb} g_{da} \left( -h^{1ef} R_{ef} + \right. \right. \\ & \quad \left. \left. g^{ef} \left( \frac{1}{2} (-h^{1f1}_{f1;f;e} - h^{1f1}_{f;f1;e} + h^{1f1}_{ff1}{}^{;f1}{}_{;e}) + \frac{1}{2} (h^{1f1}_{f;e;f1} + h^{1f1}_{e;f;f1} - h^{1fe}{}^{;f1}{}_{;f1}) \right) \right) \right) + \\ & \frac{1}{6} \left( g_{db} h^1_{ca} R + g_{ca} h^1_{db} R + g_{ca} g_{db} \left( -h^{1ef} R_{ef} + \right. \right. \\ & \quad \left. \left. g^{ef} \left( \frac{1}{2} (-h^{1f1}_{f1;f;e} - h^{1f1}_{f;f1;e} + h^{1f1}_{ff1}{}^{;f1}{}_{;e}) + \frac{1}{2} (h^{1f1}_{f;e;f1} + h^{1f1}_{e;f;f1} - h^{1fe}{}^{;f1}{}_{;f1}) \right) \right) \right) \end{aligned}$$

---

Comparison:

```
In[149]:=
  % - %% // ToCanonical
      0.184012 Second

Out[149]=
  0
```

## ■ 4. Perturbations of other objects. Densities

Usually we have other, user-defined, geometric objects defined on our manifold. For instance, matter fields. In this section we will see how to construct a tensor that is the perturbation of another one. We will also show that this can be done for densities as well.

---

Let us define a rank-3 tensor  $T$  and a rank-2 symmetric +1-density  $\Pi$ . Note that the characters of the indices at definition will be later privileged in the computation.

```
In[150]:=
  DefTensor[T[-a, b, c], M]
  DefTensor[Pi[b, c], M, Symmetric[{b, c}], WeightOfTensor -> AIndex]

  ** DefTensor: Defining tensor T[-a, b, c].
  ** DefTensor: Defining tensor Pi[b, c].
```

---

We can also define some constant symbol, whose perturbation would be automatically zero

```
In[152]:=
  DefConstantSymbol[kappa]

  ** DefConstantSymbol: Defining constant symbol kappa.
```

---

We define another two tensors which will represent the perturbations of the previous tensors, and hence must have the same density character. The syntax is parallel to that of `DefTensor`:

```
In[153]:=
  DefTensorPerturbation[t[LI[order], -a, b, c], T[-a, b, c], M]
  DefTensorPerturbation[P[LI[order], b, c],
  Pi[b, c], M, Symmetric[{b, c}], WeightOfTensor -> AIndex]

  ** DefTensor: Defining tensor t[LI[order], -a, b, c].
  ** DefTensor: Defining tensor P[LI[order], b, c].
```

---

Now we have

```
In[155]:=
  Perturbation[T[-a, b, c], 3]

Out[155]=
  tau_a^3 bc
```

```
In[156]:=
  Perturbation[Pi[a, b], 2]
```

```
Out[156]=
  P̂2 ab
```

Now we can do the perturbation of any expression containing these and the geometric tensors with any index positioning. When the indices are in the natural position, the perturbation will be directly computed

```
In[157]:=
  Perturbation[T[-a, b, c] Pi[d, e], 2]
```

```
Out[157]=
  P̂2 de Tabc + 2 P̂1 de τa1 bc + P̂de τa2 bc
```

But when they are in a different position we have to make use of ExpandPerturbation

```
In[158]:=
  Perturbation[T[-a, b, -c] Pi[-d, e], 2]
```

```
Out[158]=
  2 Δ[Tab c] Δ[ṽde] + Δ2 [ṽde] Tab c + Δ2 [Tab c] ṽde
```

```
In[159]:=
  % // ExpandPerturbation
```

```
Out[159]=
  Tab c (2 hdf1 fe P̂1 fe + gdf P̂2 fe + hdf2 ṽfe) +
  2 (gdf1 P̂1 f1e + hdf11 ṽf1e) (hcf1 Tabf + gcf τa1 bf) + ṽde (hcf2 Tabf + 2 hcf1 τa1 bf + gcf τa2 bf)
```

We can also do more difficult computations at higher order

```
In[160]:=
  RicciCD[-a, b] Pi[-b, d] +
  Pi[-a, c] RiemannCD[-c, b, -b, d] + Pi[-a, b] T[d, -b, c] T[-c, -e, e]
```

```
Out[160]=
  Tcee Tbd c ṽab - Rcd ṽac + Rab ṽbd
```

```
In[161]:=
  Perturbation[%, 3]
```

```
Out[161]=
  6 Δ[Tcee] Δ[Tbd c] Δ[ṽab] + 3 Δ[ṽbd] Δ2 [Rab] - 3 Δ[ṽac] Δ2 [Rcd] - 3 Δ[Rcd] Δ2 [ṽac] +
  3 Δ[Rab] Δ2 [ṽbd] + Δ3 [ṽbd] Rab - Δ3 [ṽac] Rcd + 3 Δ[ṽab] Δ2 [Tbd c] Tcee + 3 Δ[Tbd c] Δ2 [ṽab] Tcee +
  3 Δ[ṽab] Δ2 [Tcee] Tbd c + 3 Δ[Tcee] Δ2 [ṽab] Tbd c + Δ3 [ṽab] Tcee Tbd c + 3 Δ[Tbd c] Δ2 [Tcee] ṽab +
  3 Δ[Tcee] Δ2 [Tbd c] ṽab + Δ3 [Tbd c] Tcee ṽab + Δ3 [Tcee] Tbd c ṽab - Δ3 [Rcd] ṽac + Δ3 [Rab] ṽbd
```

```
In[162]:=
  % // ExpandPerturbation // ContractMetric // ToCanonical
```

```
7.81249 Second
```

Out[162]=

$$\begin{aligned}
 & \tilde{P}^{3db} R_{ab} + 3h^{2db} \tilde{P}^{1c} R_{bc} + 3h^{1db} \tilde{P}^{2c} R_{bc} - 6h^{1c} h^{1db} \tilde{P}^{1e} R_{ce} + 6h^{1b} h^{1dc} \tilde{P}^{1e} R_{ce} - \\
 & \tilde{P}^{3b} R_{ab}^d - 3h^{2ab} \tilde{P}^{1c} R_{bc}^d - 3h^{1ab} \tilde{P}^{2c} R_{bc}^d + 3h^{2f} \tilde{P}^{1af} T_{ce}^e T^{dbc} + 3h^{2f} \tilde{P}^{1bf} T_{ce}^e T^{dbc} + \\
 & 6h^{1f} h^{1fl} \tilde{P}^{1ffl} T_{ce}^e T^{dbc} + 3h^{1bf} \tilde{P}^{2af} T_{ce}^e T^{dbc} + 3h^{1af} \tilde{P}^{2bf} T_{ce}^e T^{dbc} + \tilde{P}^3_{ab} T_{ce}^e T^{dbc} + \\
 & 3h^{2ef} \tilde{P}^{1ab} T_{ce}^{ef} T^{dbc} + 6h^{1fl} h^{1ef} \tilde{P}^{1af1} T_{ce}^{ef} T^{dbc} + 6h^{1fl} h^{1ef} \tilde{P}^{1bf1} T_{ce}^{ef} T^{dbc} + 3h^{1ef} \tilde{P}^2_{ab} T_{ce}^{ef} T^{dbc} - \\
 & 6h^{1d} h^{1ff1} \tilde{P}^{1ac} T^{bce} T_e^{ff1} + 6h^{1df1} h^{1ef1} \tilde{P}^{1af} T^{bc} T_e^{ef} - 3h^{2de} \tilde{P}^{1af} T^{bc} T_e^{ef} - \\
 & 6h^{1de} h^{1fl} \tilde{P}^{1af1} T^{bc} T_e^{ef} - 6h^{1fl} h^{1de} \tilde{P}^{1ff1} T^{bc} T_e^{ef} - 3h^{1de} \tilde{P}^2_{af} T^{bc} T_e^{ef} + \\
 & h^3_{ef} T_{ce}^{ef} T^{dbc} \tilde{\Pi}_{ab} + h^3_{db} R_b^c \tilde{\Pi}_{ac} + 6h^{1bf10} h^{1df10} h^{1ff1} T^{bce} T_e^{ff1} \tilde{\Pi}_{ac} - 3h^{1ff1} h^{2d} T^{bce} T_e^{ff1} \tilde{\Pi}_{ac} - \\
 & 3h^{1d} h^{2ff1} T^{bce} T_e^{ff1} \tilde{\Pi}_{ac} - 3h^{1db} h^2_{bc} R_c^e \tilde{\Pi}_{ae} - 3h^{1c} h^{2db} R_c^e \tilde{\Pi}_{ae} + 6h^{1c} h^{1e} h^{1db} R_e^f \tilde{\Pi}_{af} + \\
 & h^3_{bf} T_{ce}^e T^{dbc} \tilde{\Pi}_{af} - 6h^{1df1} h^{1ef10} h^{1ff10} T^{bc} T_e^{ef} \tilde{\Pi}_{af} + 3h^{1ef1} h^{2df1} T^{bc} T_e^{ef} \tilde{\Pi}_{af} + \\
 & 3h^{1df1} h^{2ef1} T^{bc} T_e^{ef} \tilde{\Pi}_{af} - h^3_{de} T^{bc} T_e^{ef} \tilde{\Pi}_{af} + 3h^{1ef} h^{2f1} T_{ce}^{ef} T^{dbc} \tilde{\Pi}_{af1} + \\
 & 3h^{1fl} h^{2ef} T_{ce}^{ef} T^{dbc} \tilde{\Pi}_{af1} - 3h^{1fl} h^{2de} T^{bc} T_e^{ef} \tilde{\Pi}_{af1} - 3h^{1de} h^{2f1} T^{bc} T_e^{ef} \tilde{\Pi}_{af1} - \\
 & 6h^{1f10} h^{1d} h^{1ff1} T^{bce} T_e^{ff1} \tilde{\Pi}_{af10} + 6h^{1df1} h^{1ef1} h^{1f10} T^{bc} T_e^{ef} \tilde{\Pi}_{af10} - h^3_{ab} R^{dc} \tilde{\Pi}_{bc} + \\
 & 3h^{1dc} h^{2ab} R_c^e \tilde{\Pi}_{be} + 3h^{1ab} h^{2dc} R_c^e \tilde{\Pi}_{be} - 6h^{1ab} h^{1e} h^{1dc} R_e^f \tilde{\Pi}_{bf} + h^3_{af} T_{ce}^e T^{dbc} \tilde{\Pi}_{bf} + \\
 & 3h^{1ef} h^{2f1} T_{ce}^{ef} T^{dbc} \tilde{\Pi}_{bf1} + 3h^{1af} h^{2ef} T_{ce}^{ef} T^{dbc} \tilde{\Pi}_{bf1} - 6h^{1f10} h^{1d} h^{1ff1} T^{bce} T_e^{ff1} \tilde{\Pi}_{cf10} + \\
 & 3h^{1fl} h^{2f} h^{1e} T_{ce}^e T^{dbc} \tilde{\Pi}_{ff1} + 3h^{1af} h^{2f1} T_{ce}^e T^{dbc} \tilde{\Pi}_{ff1} + 6h^{1af} h^{1df10} h^{1ef10} T^{bc} T_e^{ef} \tilde{\Pi}_{ff1} - \\
 & 3h^{1de} h^{2f1} T^{bc} T_e^{ef} \tilde{\Pi}_{ff1} - 3h^{1fl} h^{2de} T^{bc} T_e^{ef} \tilde{\Pi}_{ff1} + 6h^{1af} h^{1fl} h^{1ef} T_{ce}^{ef} T^{dbc} \tilde{\Pi}_{ff1f10} - \\
 & 6h^{1fl} h^{1de} h^{1f10} T^{bc} T_e^{ef} \tilde{\Pi}_{ff1f10} + 6h^{1bf} \tilde{P}^{1af} T^{dbc} \tau_{ce}^1 + 6h^{1af} \tilde{P}^{1bf} T^{dbc} \tau_{ce}^1 + 3\tilde{P}^2_{ab} T^{dbc} \tau_{ce}^1 + \\
 & 3h^{2f} T^{dbc} \tilde{\Pi}_{af} \tau_{ce}^1 + 3h^{2f} T^{dbc} \tilde{\Pi}_{bf} \tau_{ce}^1 + 6h^{1f} h^{1fl} T^{dbc} \tilde{\Pi}_{ff1} \tau_{ce}^1 + 6h^{1ef} \tilde{P}^1_{ab} T^{dbc} \tau_{ce}^{ef} + \\
 & 3h^{2ef} T^{dbc} \tilde{\Pi}_{ab} \tau_{ce}^{ef} + 6h^{1fl} h^{1ef} T^{dbc} \tilde{\Pi}_{af1} \tau_{ce}^{ef} + 6h^{1fl} h^{1ef} T^{dbc} \tilde{\Pi}_{bf1} \tau_{ce}^{ef} + \\
 & 6\tilde{P}^1_{ab} \tau_{ce}^1 + 6h^{1bf} \tilde{\Pi}_{af} \tau_{ce}^1 + 6h^{1af} \tilde{\Pi}_{bf} \tau_{ce}^1 + 6h^{1ef} \tilde{\Pi}_{ab} \tau_{ce}^{ef} + 6h^{1ef} \tilde{\Pi}_{ab} \tau_{ce}^{ef} T^{1dbc} + \\
 & 6h^{1ef} \tilde{P}^1_{af} T^{bc} \tau_{de}^1 + 6h^{1af} \tilde{P}^1_{ef} T^{bc} \tau_{de}^1 + 3\tilde{P}^2_{ae} T^{bc} \tau_{de}^1 + 3h^{2e} T^{bc} \tilde{\Pi}_{af} \tau_{de}^1 + \\
 & 3h^{2f} T^{bc} \tilde{\Pi}_{ef} \tau_{de}^1 + 6h^{1f} h^{1fl} T^{bc} \tilde{\Pi}_{ff1} \tau_{de}^1 + 6h^{1ce} \tilde{P}^1_{af} T^{bce} \tau_{de}^1 + 3h^{2ce} T^{bce} \tilde{\Pi}_{af} \tau_{de}^1 + \\
 & 6h^{1ce} h^{1fl} T^{bce} \tilde{\Pi}_{af1} \tau_{de}^1 + 6h^{1fl} h^{1ce} T^{bce} \tilde{\Pi}_{ff1} \tau_{de}^1 - 6h^{1d} \tilde{P}^1_{ac} T^{bce} \tau_{de}^1 + \\
 & 6h^{1bf1} h^{1df1} T^{bce} \tilde{\Pi}_{ac} \tau_{de}^1 - 3h^{2d} T^{bce} \tilde{\Pi}_{ac} \tau_{de}^1 - 6h^{1c} h^{1d} T^{bce} \tilde{\Pi}_{af1} \tau_{de}^1 - \\
 & 6h^{1fl} h^{1d} T^{bce} \tilde{\Pi}_{cf1} \tau_{de}^1 - 6h^{1d} h^{1fl} T^{bce} \tilde{\Pi}_{ac} \tau_{de}^1 - 6h^{1de} \tilde{P}^1_{af} T^{bc} \tau_{de}^1 + \\
 & 6h^{1df1} h^{1ef1} T^{bc} \tilde{\Pi}_{af} \tau_{de}^1 - 3h^{2de} T^{bc} \tilde{\Pi}_{af} \tau_{de}^1 - 6h^{1de} h^{1fl} T^{bc} \tilde{\Pi}_{af1} \tau_{de}^1 - \\
 & 6h^{1fl} h^{1de} T^{bc} \tilde{\Pi}_{ff1} \tau_{de}^1 - 6h^{1de} \tilde{\Pi}_{af} \tau_{de}^1 - 6h^{1ce} h^{1d} T^{bce} \tilde{\Pi}_{af1} \tau_{de}^1 + \\
 & 3\tilde{P}^1_{ab} T^{dbc} \tau_{ce}^2 + 3h^{1bf} T^{dbc} \tilde{\Pi}_{af} \tau_{ce}^2 + 3h^{1af} T^{dbc} \tilde{\Pi}_{bf} \tau_{ce}^2 + 3\tilde{\Pi}_{ab} \tau_{de}^1 \tau_{ce}^2 + \\
 & 3h^{1ef} T^{dbc} \tilde{\Pi}_{ab} \tau_{ce}^2 + 3\tilde{\Pi}_{ab} \tau_{ce}^2 T^{2dbc} + 3\tilde{P}^1_{ae} T^{bc} \tau_{de}^2 + 3h^{1f} T^{bc} \tilde{\Pi}_{af} \tau_{de}^2 + \\
 & 3h^{1af} T^{bc} \tilde{\Pi}_{ef} \tau_{de}^2 + 3h^{1ce} T^{bce} \tilde{\Pi}_{af} \tau_{de}^2 - 3h^{1d} T^{bce} \tilde{\Pi}_{ac} \tau_{de}^2 - 3h^{1de} T^{bc} \tilde{\Pi}_{af} \tau_{de}^2 + \\
 & T^{dbc} \tilde{\Pi}_{ab} \tau_{ce}^3 + T^{bc} \tilde{\Pi}_{ae} \tau_{de}^3 + 3h^{1bc} \tilde{P}^{1de} h^{1bc;e;a} + \frac{3}{2} h^{2bc} \tilde{\Pi}^{de} h^{1bc;e;a} + \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h^{2bc;e;a} - \\
 & 3h^{1e} h^{1bc} \tilde{\Pi}^{df} h^{1ce;f;a} + \frac{3}{2} \tilde{P}^{1db} h^{1ce;ia} h^{1ce;ib} + \frac{3}{4} \tilde{\Pi}^{db} h^{2ce;ia} h^{1ce;ib} + \frac{3}{4} \tilde{\Pi}^{db} h^{1ce;ia} h^{2ce;ib} - \\
 & \frac{3}{2} \tilde{P}^{2db} h^{1c;ia;b} - \frac{3}{2} \tilde{P}^{1db} h^{2c;ia;b} - \frac{1}{2} \tilde{\Pi}^{db} h^{3c;ia;b} + \frac{3}{2} \tilde{P}^{2b} h^{1c;id;b} + \frac{3}{2} \tilde{P}^{1b} h^{2c;id;b} + \\
 & \frac{1}{2} \tilde{\Pi}_a^b h^{3c;id;b} + 3h^{1bc} \tilde{\Pi}^{de} h^{1e;fa} h^{1bf;c} - 3h^{1bc} \tilde{\Pi}^{de} h^{1f;ib} h^{1ef;c} + 3h^{1bc} \tilde{\Pi}_a^e h^{1df;ib} h^{1ef;c} + \\
 & \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{1ef;ib} h^{1ef;c} + \frac{3}{2} \tilde{P}^{1db} h^{1c;ia} h^{1e;ic} + \frac{3}{4} \tilde{\Pi}^{db} h^{2c;ia} h^{1e;ic} + \frac{3}{2} \tilde{P}^{1db} h^{1c;ib} h^{1e;ic} - \\
 & \frac{3}{2} \tilde{P}^1_{ab} h^{1dc;ib} h^{1e;ic} + \frac{3}{4} \tilde{\Pi}^{db} h^{2c;ib} h^{1e;ic} - \frac{3}{4} \tilde{\Pi}_a^b h^{2dc;ib} h^{1e;ic} - \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h^{1be;ia} h^{1f;ic} + \\
 & \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h^{1ae;ib} h^{1f;ic} - \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h^{1d;eb} h^{1f;ic} + \frac{3}{4} \tilde{\Pi}^{db} h^{1c;ia} h^{2e;ic} + \frac{3}{4} \tilde{\Pi}^{db} h^{1c;ib} h^{2e;ic} - \\
 & \frac{3}{4} \tilde{\Pi}_a^b h^{1dc;ib} h^{2e;ic} + \frac{3}{2} \tilde{P}^{2db} h^{1c;ia;c} - 3h^{1bc} \tilde{P}^{1de} h^{1be;ia;c} - \frac{3}{2} h^{2bc} \tilde{\Pi}^{de} h^{1be;ia;c} + \frac{3}{2} \tilde{P}^{1db} h^{2c;ia;c} - \\
 & \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h^{2be;ia;c} + \frac{1}{2} \tilde{\Pi}^{db} h^{3c;ia;c} + \frac{3}{2} \tilde{P}^{2db} h^{1c;ib;c} + 3h^{1bc} \tilde{P}^{1de} h^{1ae;ib;c} + \frac{3}{2} h^{2bc} \tilde{\Pi}^{de} h^{1ae;ib;c} -
 \end{aligned}$$

$$\begin{aligned}
& \frac{3}{2} \tilde{P}_a^{2b} h^{1dc}_{;b;c} - 3 h^{1bc} \tilde{P}_a^{1e} h^{1d}_{e;b;c} - \frac{3}{2} h^{2bc} \tilde{\Pi}_a^e h^{1d}_{e;b;c} - 3 h^{1db} \tilde{P}_a^{1c} h^{1e}_{e;b;c} - \\
& \frac{3}{2} h^{2db} \tilde{\Pi}_a^c h^{1e}_{e;b;c} + \frac{3}{2} \tilde{P}^{1db} h^{2c}_{;b;c} + \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h^{2ae}_{;b;c} - \frac{3}{2} \tilde{P}_a^{1b} h^{2dc}_{;b;c} - \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h^{2d}_{e;b;c} - \\
& \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{2e}_{e;b;c} + \frac{1}{2} \tilde{\Pi}^{db} h^{3c}_{;b;c} - \frac{1}{2} \tilde{\Pi}_a^b h^{3dc}_{;b;c} - \frac{3}{2} \tilde{P}^{2db} h^{1ab}_{;c} + \frac{3}{2} \tilde{P}_a^{2b} h^{1d}_{b;c} - \\
& \frac{3}{2} \tilde{P}^{1db} h^{2ab}_{;c} + \frac{3}{2} \tilde{P}_a^{1b} h^{2d}_{b;c} - \frac{1}{2} \tilde{\Pi}^{db} h^{3ab}_{;c} + \frac{1}{2} \tilde{\Pi}_a^b h^{3d}_{b;c} - \frac{3}{2} \tilde{P}_a^{2b} h^{1c;d}_{;c} + \\
& 3 h^{1bc} \tilde{P}_a^{1e} h^{1be}_{;id} + \frac{3}{2} h^{2bc} \tilde{\Pi}_a^e h^{1be}_{;id} + 3 h^{1ab} \tilde{P}_b^{1c} h^{1e}_{e;id} + \frac{3}{2} h^{2ab} \tilde{\Pi}_b^c h^{1e}_{e;id} - \\
& \frac{3}{2} \tilde{P}_a^{1b} h^{2c;d}_{;c} + \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h^{2be}_{;id} + \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{2e}_{e;id} - \frac{1}{2} \tilde{\Pi}_a^b h^{3c;d}_{;c} - 3 h^{1bc} \tilde{P}^{1de} h^{1ab}_{e;ic} - \\
& \frac{3}{2} h^{2bc} \tilde{\Pi}^{de} h^{1ab}_{e;ic} + 3 h^{1bc} \tilde{P}_a^{1e} h^{1d}_{b;e;ic} + \frac{3}{2} h^{2bc} \tilde{\Pi}_a^e h^{1d}_{b;e;ic} - \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h^{2ab}_{e;ic} + \\
& \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h^{2d}_{b;e;ic} - \frac{3}{2} \tilde{P}^{1db} h^{1e}_{e;ic} h^{1ab}_{;ic} - \frac{3}{4} \tilde{\Pi}^{db} h^{2e}_{e;ic} h^{1ab}_{;ic} + \frac{3}{2} \tilde{P}_a^{1b} h^{1e}_{e;ic} h^{1d}_{b;c} + \\
& \frac{3}{4} \tilde{\Pi}_a^b h^{2e}_{e;ic} h^{1d}_{b;c} - \frac{3}{4} \tilde{\Pi}^{db} h^{1e}_{e;ic} h^{2ab}_{;ic} + \frac{3}{4} \tilde{\Pi}_a^b h^{1e}_{e;ic} h^{2d}_{b;c} - \frac{3}{2} \tilde{P}_a^{1b} h^{1e}_{e;ic} h^{1b}_{c;d} - \\
& \frac{3}{4} \tilde{\Pi}_a^b h^{2e}_{e;ic} h^{1b}_{c;d} + \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h^{1f}_{f;ic} h^{1be}_{;id} - \frac{3}{2} \tilde{P}_a^{1b} h^{1ce}_{ib} h^{1ce;d} - \frac{3}{4} \tilde{\Pi}_a^b h^{2ce}_{ib} h^{1ce;d} - \\
& 3 h^{1bc} \tilde{\Pi}_a^e h^{1bf}_{f;ic} h^{1e}_{e;id} - \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{1ef}_{f;ic} h^{1ef;id} - \frac{3}{4} \tilde{\Pi}_a^b h^{1e}_{e;ic} h^{2c;d}_{;c} - \frac{3}{4} \tilde{\Pi}_a^b h^{1ce}_{ib} h^{2ce;d}_{;c} - \\
& \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h^{1f}_{f;ic} h^{1ab}_{e;ic} + 3 h^{1bc} \tilde{\Pi}^{de} h^{1bf}_{f;ic} h^{1e}_{e;ic} - 3 \tilde{P}^{1db} h^{1b}_{c;ia} h^{1e}_{e;ic} - \frac{3}{2} \tilde{\Pi}^{db} h^{2b}_{c;ia} h^{1e}_{e;ic} - \\
& 3 \tilde{P}^{1db} h^{1c}_{;ib} h^{1e}_{e;ic} + 3 \tilde{P}_a^{1b} h^{1dc}_{;ib} h^{1e}_{e;ic} - \frac{3}{2} \tilde{\Pi}^{db} h^{2c}_{;ib} h^{1e}_{e;ic} + \frac{3}{2} \tilde{\Pi}_a^b h^{2dc}_{;ib} h^{1e}_{e;ic} + \\
& 3 \tilde{P}^{1db} h^{1ab}_{;ic} h^{1e}_{e;ic} - 3 \tilde{P}_a^{1b} h^{1d}_{b;c} h^{1e}_{e;ic} + \frac{3}{2} \tilde{\Pi}^{db} h^{2ab}_{;ic} h^{1e}_{e;ic} - \frac{3}{2} \tilde{\Pi}_a^b h^{2d}_{b;c} h^{1e}_{e;ic} + \\
& 3 \tilde{P}_a^{1b} h^{1c;d}_{;c} h^{1e}_{e;ic} + \frac{3}{2} \tilde{\Pi}_a^b h^{2c;d}_{;c} h^{1e}_{e;ic} - 3 h^{1bc} \tilde{\Pi}^{de} h^{1f}_{f;ia} h^{1cf}_{f;ie} + 3 h^{1bc} \tilde{\Pi}_a^e h^{1bf}_{f;id} h^{1cf}_{f;ie} + \\
& \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h^{1f}_{f;ic} h^{1d}_{b;ie} - 3 h^{1bc} \tilde{\Pi}_a^e h^{1bf}_{f;ic} h^{1df}_{f;ie} + \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{1e}_{e;ib} h^{1f}_{f;ie} + \\
& \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{1e}_{e;ic} h^{1f}_{f;ie} - \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{1de}_{;ic} h^{1f}_{f;ie} - \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{1e}_{e;id} h^{1f}_{f;ie} - \\
& \frac{3}{2} \tilde{\Pi}^{db} h^{1b}_{c;ia} h^{2e}_{e;ic} - \frac{3}{2} \tilde{\Pi}^{db} h^{1c}_{;ib} h^{2e}_{e;ic} + \frac{3}{2} \tilde{\Pi}_a^b h^{1dc}_{;ib} h^{2e}_{e;ic} + \frac{3}{2} \tilde{\Pi}^{db} h^{1ab}_{;ic} h^{2e}_{e;ic} - \\
& \frac{3}{2} \tilde{\Pi}_a^b h^{1d}_{b;c} h^{2e}_{e;ic} + \frac{3}{2} \tilde{\Pi}_a^b h^{1c;d}_{;c} h^{2e}_{e;ic} + 3 h^{1b} h^{1bc} \tilde{\Pi}^{df} h^{1cf}_{f;ia;ie} + 3 h^{1db} \tilde{P}_a^{1c} h^{1e}_{e;ib;ie} + \\
& \frac{3}{2} h^{2db} \tilde{\Pi}_a^c h^{1e}_{e;ib;ie} - 3 h^{1ce} h^{1db} \tilde{\Pi}_a^f h^{1cf}_{f;ib;ie} + \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{2e}_{e;ib;ie} - 3 h^{1b} h^{1bc} \tilde{\Pi}^{df} h^{1af}_{f;ic;ie} + \\
& 3 h^{1db} \tilde{P}_a^{1c} h^{1e}_{e;ic;ie} + \frac{3}{2} h^{2db} \tilde{\Pi}_a^c h^{1e}_{e;ic;ie} + 3 h^{1ce} h^{1db} \tilde{\Pi}_a^f h^{1bf}_{f;ic;ie} - 3 h^{1ab} \tilde{P}_b^{1c} h^{1de}_{;ic;ie} - \\
& \frac{3}{2} h^{2ab} \tilde{\Pi}_b^c h^{1de}_{;ic;ie} + 3 h^{1b} h^{1bc} \tilde{\Pi}_a^f h^{1d}_{f;ic;ie} - 3 h^{1ab} h^{1ce} \tilde{\Pi}_b^f h^{1d}_{f;ic;ie} + 3 h^{1c} h^{1db} \tilde{\Pi}_a^e h^{1f}_{f;ic;ie} - \\
& 3 h^{1ab} h^{1dc} \tilde{\Pi}_b^e h^{1f}_{f;ic;ie} + \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{2e}_{e;ic;ie} - \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{2de}_{;ic;ie} - 3 h^{1bc} \tilde{P}_a^{1e} h^{1bc}_{;id;ie} - \\
& \frac{3}{2} h^{2bc} \tilde{\Pi}_a^e h^{1bc}_{;id;ie} - 3 h^{1ab} \tilde{P}_b^{1c} h^{1e}_{e;id;ie} - \frac{3}{2} h^{2ab} \tilde{\Pi}_b^c h^{1e}_{e;id;ie} - 3 h^{1b} h^{1bc} \tilde{\Pi}_a^f h^{1cf}_{f;id;ie} + \\
& 3 h^{1ab} h^{1ce} \tilde{\Pi}_b^f h^{1cf}_{f;id;ie} - \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h^{2bc}_{;id;ie} - \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{2c}_{;id;ie} - 3 h^{1db} \tilde{P}_a^{1c} h^{1bc}_{;ie;ie} - \\
& \frac{3}{2} h^{2db} \tilde{\Pi}_a^c h^{1bc}_{;ie;ie} + 3 h^{1ab} \tilde{P}_b^{1c} h^{1d}_{c;ie;ie} + \frac{3}{2} h^{2ab} \tilde{\Pi}_b^c h^{1d}_{c;ie;ie} - \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{2bc}_{;ie;ie} + \\
& \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{2d}_{c;ie;ie} + 3 h^{1b} h^{1bc} \tilde{\Pi}^{df} h^{1ac}_{f;ie} - 3 h^{1ce} h^{1db} \tilde{\Pi}_a^f h^{1bc}_{f;ie} - 3 h^{1b} h^{1bc} \tilde{\Pi}_a^f h^{1d}_{c;f;ie} + \\
& 3 h^{1ab} h^{1ce} \tilde{\Pi}_b^f h^{1d}_{c;f;ie} - 3 \tilde{P}^{1db} h^{1be}_{;ic} h^{1a}_{c;ie} - \frac{3}{2} \tilde{\Pi}^{db} h^{2be}_{;ic} h^{1a}_{c;ie} + 3 \tilde{P}^{1db} h^{1bc}_{;ie} h^{1a}_{c;ie} + \\
& \frac{3}{2} \tilde{\Pi}^{db} h^{2bc}_{;ie} h^{1a}_{c;ie} - \frac{3}{2} h^{1db} \tilde{\Pi}_a^c h^{1f}_{f;ie} h^{1bc}_{;ie} + \frac{3}{2} h^{1ab} \tilde{\Pi}_b^c h^{1f}_{f;ie} h^{1d}_{c;ie} + 3 \tilde{P}_a^{1b} h^{1be}_{;ic} h^{1dc}_{;ie} + \\
& \frac{3}{2} \tilde{\Pi}_a^b h^{2be}_{;ic} h^{1dc}_{;ie} - 3 \tilde{P}_a^{1b} h^{1bc}_{;ie} h^{1dc}_{;ie} - \frac{3}{2} \tilde{\Pi}_a^b h^{2bc}_{;ie} h^{1dc}_{;ie} - \frac{3}{2} \tilde{\Pi}^{db} h^{1be}_{;ic} h^{2c}_{a;ie} +
\end{aligned}$$

$$\begin{aligned}
& \frac{3}{2} \tilde{\Pi}^{db} h_{bc;e}^1 h_a^{2c;e} + \frac{3}{2} \tilde{\Pi}_a^b h_{be;c}^1 h^{2dc;e} - \frac{3}{2} \tilde{\Pi}_a^b h_{bc;e}^1 h^{2dc;e} - \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h_{e;ia}^1 h_{bc;f}^1 + \\
& \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h_{e;id}^1 h_{bc;f}^1 - \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h_{a;ie}^1 h_{bc;f}^1 + \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h_{ie}^{1df} h_{bc;f}^1 + \\
& 3 h^{1bc} \tilde{\Pi}^{de} h_{a;ib}^1 h_{ce;f}^1 - 3 h^{1bc} \tilde{\Pi}_a^e h_{ib}^{1df} h_{ce;f}^1 + 3 h^{1bc} \tilde{\Pi}^{de} h_{be;ia}^1 h_{c;f}^1 - 3 h^{1bc} \tilde{\Pi}^{de} h_{ae;ib}^1 h_{c;f}^1 + \\
& 3 h^{1bc} \tilde{\Pi}_a^e h_{e;ib}^1 h_{c;f}^1 - 3 h^{1bc} \tilde{\Pi}_a^e h_{be;id}^1 h_{c;f}^1 + 3 h^{1bc} \tilde{\Pi}^{de} h_{ab;ie}^1 h_{c;f}^1 - 3 h^{1bc} \tilde{\Pi}_a^e h_{b;ie}^1 h_{c;f}^1 - \\
& 3 h^{1db} \tilde{\Pi}_a^c h_{c;ib}^1 h_{e;f}^1 - 3 h^{1db} \tilde{\Pi}_a^c h_{b;ic}^1 h_{e;f}^1 + 3 h_a^1 b \tilde{\Pi}_b^c h_{ic}^{1de} h_{e;f}^1 + 3 h_a^1 b \tilde{\Pi}_b^c h_{c;e;d}^1 h_{e;f}^1 + \\
& 3 h^{1db} \tilde{\Pi}_a^c h_{bc;ie}^1 h_{e;f}^1 - 3 h_a^1 b \tilde{\Pi}_b^c h_{ic}^{1de} h_{e;f}^1 + 3 h^{1ce} h^{1db} \tilde{\Pi}_a^f h_{ce;b;f}^1 - 3 h_a^1 b \tilde{\Pi}_b^c h^{1db} \tilde{\Pi}_a^e h_{e;ic;f}^1 + \\
& 3 h_a^1 b h^{1dc} \tilde{\Pi}_b^e h_{e;ic;f}^1 + 3 h_b^1 h^{1bc} \tilde{\Pi}_a^f h_{ce;id;f}^1 - 3 h_a^1 b h^{1ce} \tilde{\Pi}_b^f h_{ce;id;f}^1 - 3 h_b^1 c h^{1db} \tilde{\Pi}_a^e h_{c;ie;f}^1 + \\
& 3 h_a^1 b h^{1dc} \tilde{\Pi}_b^e h_{c;ie;f}^1 + 3 h_b^1 c h^{1db} \tilde{\Pi}_a^e h_{ce;id;f}^1 - 3 h_a^1 b h^{1dc} \tilde{\Pi}_b^e h_{ce;id;f}^1 + 3 h^{1bc} \tilde{\Pi}^{de} h_{ef;c}^1 h_{ab}^{1;f} - \\
& 3 h^{1bc} \tilde{\Pi}^{de} h_{ce;f}^1 h_{ab}^{1;f} - 3 h^{1bc} \tilde{\Pi}^{de} h_{bf;ic}^1 h_{ae}^{1;f} + \frac{3}{2} h^{1bc} \tilde{\Pi}^{de} h_{bc;f}^1 h_{ae}^{1;f} - \\
& 3 h^{1db} \tilde{\Pi}_a^c h_{cf;ie}^1 h_{b;e;f}^1 + 3 h^{1db} \tilde{\Pi}_a^c h_{ce;f}^1 h_{b;e;f}^1 - 3 h^{1bc} \tilde{\Pi}_a^e h_{ef;ic}^1 h_b^{1d;f} + 3 h^{1bc} \tilde{\Pi}_a^e h_{ce;f}^1 h_b^{1d;f} + \\
& 3 h^{1bc} \tilde{\Pi}_a^e h_{bf;ic}^1 h_e^{1d;f} - \frac{3}{2} h^{1bc} \tilde{\Pi}_a^e h_{bc;f}^1 h_e^{1d;f} + 3 h_a^1 b \tilde{\Pi}_b^c h_{cf;ie}^1 h^{1de;f} - 3 h_a^1 b \tilde{\Pi}_b^c h_{ce;f}^1 h^{1de;f}
\end{aligned}$$

In the context of `xTensor``, we can define scalar functions. This is a function of two arguments:

```

In[163]:=
DefScalarFunction[F, 2]

** DefScalarFunction: Defining scalar function F.

```

It must have scalar arguments. For example,

```

In[164]:=
F[RicciScalarCD[], T[a, b, c] RicciCD[-a, -b] v[-c]]

Out[164]=
F[R, Rab Tabc vc]

```

The command `Perturbation` knows how to apply the chain rule when acting on this kind of function at any order. It is essentially implemented through the *Mathematica* total derivative command `Dt`.

```
In[165]:=
```

```
Perturbation[% , 3]
```

```
Out[165]=
```

$$\begin{aligned}
& (6 \Delta[R_{ab}] \Delta[T^{abc}] \Delta[v_c] + 3 \Delta[v_c] \Delta^2[T^{abc}] R_{ab} + \\
& \quad 3 \Delta[T^{abc}] \Delta^2[v_c] R_{ab} + 3 \Delta[v_c] \Delta^2[R_{ab}] T^{abc} + 3 \Delta[R_{ab}] \Delta^2[v_c] T^{abc} + \Delta^3[v_c] R_{ab} T^{abc} + \\
& \quad 3 \Delta[T^{abc}] \Delta^2[R_{ab}] v_c + 3 \Delta[R_{ab}] \Delta^2[T^{abc}] v_c + \Delta^3[T^{abc}] R_{ab} v_c + \Delta^3[R_{ab}] T^{abc} v_c) \\
& F^{(0,1)}[R, R_{ab} T^{abc} v_c] + 3 (\Delta[v_e] R_{ac} T^{ace} + \Delta[T^{ace}] R_{ac} v_e + \Delta[R_{ac}] T^{ace} v_e) \\
& (2 \Delta[T^{bdf}] \Delta[v_f] R_{bd} + 2 \Delta[R_{bd}] \Delta[v_f] T^{bdf} + \Delta^2[v_f] R_{bd} T^{bdf} + \\
& \quad 2 \Delta[R_{bd}] \Delta[T^{bdf}] v_f + \Delta^2[T^{bdf}] R_{bd} v_f + \Delta^2[R_{bd}] T^{bdf} v_f) F^{(0,2)}[R, R_{ab} T^{abc} v_c] + \\
& (\text{Scalar}[\Delta[v_c] R_{ab} T^{abc}] + \text{Scalar}[\Delta[T^{abc}] R_{ab} v_c] + \text{Scalar}[\Delta[R_{ab}] T^{abc} v_c])^3 \\
& F^{(0,3)}[R, R_{ab} T^{abc} v_c] + \Delta^3[R] F^{(1,0)}[R, R_{ab} T^{abc} v_c] + \\
& 3 \Delta^2[R] (\Delta[v_c] R_{ab} T^{abc} + \Delta[T^{abc}] R_{ab} v_c + \Delta[R_{ab}] T^{abc} v_c) F^{(1,1)}[R, R_{ab} T^{abc} v_c] + \\
& 3 \Delta[R] (2 \Delta[T^{abc}] \Delta[v_c] R_{ab} + 2 \Delta[R_{ab}] \Delta[v_c] T^{abc} + \Delta^2[v_c] R_{ab} T^{abc} + \\
& \quad 2 \Delta[R_{ab}] \Delta[T^{abc}] v_c + \Delta^2[T^{abc}] R_{ab} v_c + \Delta^2[R_{ab}] T^{abc} v_c) F^{(1,1)}[R, R_{ab} T^{abc} v_c] + \\
& 3 \Delta[R] (\text{Scalar}[\Delta[v_c] R_{ab} T^{abc}] + \text{Scalar}[\Delta[T^{abc}] R_{ab} v_c] + \text{Scalar}[\Delta[R_{ab}] T^{abc} v_c])^2 \\
& F^{(1,2)}[R, R_{ab} T^{abc} v_c] + 3 \Delta[R] \Delta^2[R] F^{(2,0)}[R, R_{ab} T^{abc} v_c] + \\
& 3 \text{Scalar}[\Delta[R]]^2 (\Delta[v_c] R_{ab} T^{abc} + \Delta[T^{abc}] R_{ab} v_c + \Delta[R_{ab}] T^{abc} v_c) F^{(2,1)}[R, R_{ab} T^{abc} v_c] + \\
& \text{Scalar}[\Delta[R]]^3 F^{(3,0)}[R, R_{ab} T^{abc} v_c]
\end{aligned}$$

Again we find the head `Scalar`, to avoid problems with powers of scalars formed by contracted products of tensors.

## ■ 5. Gauge transformations

Through all this notebook, we have implicitly supposed a family of manifolds  $M(\epsilon)$  equipped with metric  $g(\epsilon)$ , that depend on a dimensionless parameter  $\epsilon$ . The member with parameter  $\epsilon=0$  will be called the background spacetime. The objective of perturbation theory is to compare perturbed objects, for instance a tensor  $T(\epsilon)$ , with their background counterparts, in this case  $T(0)$ . The problem is that they are defined in different manifolds. Hence, we have to construct a mapping (the so-called gauge)  $\phi(\epsilon): M(\epsilon) \rightarrow M(0)$ , whose pull-back  $\phi^*(\epsilon)$  will take tensors defined in the tangent space of  $M(\epsilon)$  to the tangent space of  $M(0)$ . Then, the difference between perturbed and background tensors will define its perturbation. For instance, the full perturbation of the tensor  $T(0)$  is given by  $\phi^*(\epsilon)(T(\epsilon)) - T(0)$ . This expression will be expanded in a power series of the parameter  $\epsilon$ . Through this expansion, the different perturbative orders of the tensor  $T(0)$  will be defined. But, if we change the gauge  $\phi$  to a different one  $\psi$ , the new perturbation of the tensor will change to  $\psi^*(\epsilon)(T(\epsilon)) - T(0)$ , which will affect to all the perturbative orders. The theory and formulas for the power expansion of changes of gauge was developed by Bruni et al. in reference *Class. Quantum Grav.* 14, 2585 (1997).

In the context of `xPert`, we have constructed a command named `GaugeChange[expr,  $\xi$ ]` that performs the gauge transformation of the expression `expr` parameterized by the family of vector fields  $\xi[LI[n], a]$ , where `n` stands for the perturbative order.

We define the gauge generators. They must be a family of vector fields, one for each perturbative order. This is again taken into account by the label `LI`.

```
In[166]:=
```

```
DefTensor[ $\xi[LI[n], a], M]$ 
```

```
** DefTensor: Defining tensor  $\xi[LI[n], a]$ .
```



The background objects do not depend on the perturbative gauge chosen, so they do not change under a gauge transformation.

```
In[167]:=
  GaugeChange[
    g[-a, -b] RicciCD[c, d] + EinsteinCD[-a, c] EinsteinCD[-b, d] RicciScalarCD[], ξ]
Out[167]=
  gab Rcd + Gac Gbd R
```

The gauge transformation of the first-order metric perturbation is the following well-known formula.

```
In[168]:=
  GaugeChange[h[LI[1], -a, -b], ξ]
Out[168]=
  h1ab + ℒξ1 gab
```

The first-order gauge transformation of any other object is similar:

```
In[169]:=
  GaugeChange[τ[LI[1], -a, b, c], ξ]
Out[169]=
  τ1abc + ℒξ1 Tabc
In[170]:=
  % // ToCanonical
Out[170]=
  τ1abc + ℒξ1 Tabc
```

At second order this formula is more complicated

```
In[171]:=
  GaugeChange[h[LI[2], -a, -b], ξ]
Out[171]=
  h2ab + 2 (ℒξ1 h1ab) + ℒξ1 ℒξ1 gab + ℒξ2 gab
In[172]:=
  % // ToCanonical
Out[172]=
  h2ab + 2 (ℒξ1 h1ab) + ℒξ1 ℒξ1 gab + ℒξ2 gab
```

Note that this previous formula is applied when the indices are in the natural position and that the metric, since it is a background object, commutes with the gauge change. This means that the following objects are the same:

```
In[173]:=
  g[a, d] g[b, e] GaugeChange[τ[LI[3], -d, -e, c], ξ]
Out[173]=
  gad (τ3dbc + 3 (ℒξ1 τ2dbc) + 3 (ℒξ1 ℒξ1 τ1dbc) +
  ℒξ1 ℒξ1 ℒξ1 Tdbc + 3 (ℒξ1 ℒξ2 Tdbc) + 3 (ℒξ2 τ1dbc) + ℒξ3 Tdbc)
```

```

In[174]:=
  GaugeChange[τ[LI[3], a, b, c], ξ]

Out[174]=
  gad (τdbc + 3 (Lξ1 τdbc) + 3 (Lξ1 Lξ1 τdbc) +
    Lξ1 Lξ1 Lξ1 Tdbc + 3 (Lξ1 Lξ2 Tdbc) + 3 (Lξ2 τdbc) + Lξ3 Tdbc)

In[175]:=
  ToCanonical[% - %, UseMetricOnVBundle → None]

0.132008 Second

Out[175]=
  0

```

We can also do more difficult higher-order calculations. For instance,

```

In[176]:=
  GaugeChange[τ[LI[2], a, b, c] h[LI[5], -a, -b], ξ]

Out[176]=
  gad (τdbc + 2 (Lξ1 τdbc) + Lξ1 Lξ1 Tdbc + Lξ2 Tdbc)
  (h5ab + 5 (Lξ1 h4ab) + 10 (Lξ1 Lξ1 h3ab) + 10 (Lξ1 Lξ1 Lξ1 h2ab) + 5 (Lξ1 Lξ1 Lξ1 Lξ1 h1ab) +
  30 (Lξ1 Lξ1 Lξ2 h1ab) + Lξ1 Lξ1 Lξ1 Lξ1 Lξ1 gab + 10 (Lξ1 Lξ1 Lξ1 Lξ2 gab) + 10 (Lξ1 Lξ1 Lξ3 gab) +
  30 (Lξ1 Lξ2 h2ab) + 15 (Lξ1 Lξ2 Lξ2 gab) + 20 (Lξ1 Lξ3 h1ab) + 5 (Lξ1 Lξ4 gab) +
  10 (Lξ2 h3ab) + 15 (Lξ2 Lξ2 h1ab) + 10 (Lξ2 Lξ3 gab) + 10 (Lξ3 h2ab) + 5 (Lξ4 h1ab) + Lξ5 gab)

```

## ■ 6. Background field perturbations

Making use of `xPert`, we can also do background field perturbations. This approach is often used in quantum theory. There the perturbed metric is considered to be the background one plus a unique perturbation  $h$ . But when computing the inverse of the metric this unique perturbation will coupled to itself.

We can restrict our perturbation scheme to reproduce background field perturbations:

This can be implemented by forcing the vanishing of all perturbations of the metric of order greater than one, with the following rule:

```

In[177]:=
  bgfield = h[LI[order_], ___] => 0 /; order > 1;

```

Then, the perturbed metric would be

```

In[178]:=
  Perturbed[g[-a, -b], 4] /. bgfield

Out[178]=
  gab + ε h1ab

```

And its inverse

```
In[179]:=
(Perturbed[g[b, c], 4] // ExpandPerturbation) /. bgfield // org
```

```
Out[179]=
gbc - ε h1bc + ε2 h1ba h1ca - ε3 h1ad h1ba h1cd + ε4 h1ea h1ba h1cd h1de
```

This is consistent because it leads to

```
In[180]:=
%% // org
```

```
Out[180]=
δac + ε5 h1ab h1eb h1cd h1df h1ef
```

We can compute the perturbation of any other geometric object in this context. For example the fourth-order perturbed Riemann will take the following form

```
In[181]:=
(Perturbed[RiemannCD[-a, -b, -c, d], 4] // ExpandPerturbation) /. bgfield // org
```

1.6401 Second

```
Out[181]=
```

$$\begin{aligned}
& R_{abc}{}^d + \epsilon \left( -\frac{h^1{}^d{}_{c;bia}}{2} - \frac{h^1{}^d{}_{b;cia}}{2} + \frac{h^1{}^id{}_{bc;ia}}{2} + \frac{h^1{}^d{}_{c;ia;b}}{2} + \frac{h^1{}^d{}_{a;ic;b}}{2} - \frac{h^1{}^id{}_{ac;ib}}{2} \right) + \\
& \epsilon^2 \left( \frac{1}{2} h^1{}^{de} h^1{}_{ce;bia} + \frac{1}{2} h^1{}^{de} h^1{}_{be;cia} - \frac{1}{2} h^1{}^{de} h^1{}_{bc;eia} + \frac{1}{4} h^1{}^d{}_{e;ia} h^1{}^e{}_{c;ib} - \right. \\
& \quad \frac{1}{4} h^1{}^e{}_{c;ia} h^1{}^d{}_{e;ib} - \frac{1}{2} h^1{}^{de} h^1{}_{ce;a;b} - \frac{1}{2} h^1{}^{de} h^1{}_{ae;c;b} + \frac{1}{2} h^1{}^{de} h^1{}_{ac;e;b} - \frac{1}{4} h^1{}^d{}_{e;ib} h^1{}^e{}_{a;ic} + \\
& \quad \frac{1}{4} h^1{}^d{}_{e;ia} h^1{}^e{}_{b;ic} + \frac{1}{4} h^1{}_{ce;ib} h^1{}^e{}_{a;id} + \frac{1}{4} h^1{}_{be;ic} h^1{}^e{}_{a;id} - \frac{1}{4} h^1{}^e{}_{a;ic} h^1{}_{be}{}^id - \frac{1}{4} h^1{}_{ce;a} h^1{}^e{}_{b;id} - \\
& \quad \frac{1}{4} h^1{}^e{}_{a;id} h^1{}_{bc;ie} + \frac{1}{4} h^1{}^e{}_{a;ic} h^1{}^d{}_{b;ie} + \frac{1}{4} h^1{}^d{}_{e;ib} h^1{}_{ac}{}^ie + \frac{1}{4} h^1{}_{be}{}^id h^1{}_{ac}{}^ie - \frac{1}{4} h^1{}^d{}_{b;ie} h^1{}_{ac}{}^ie - \\
& \quad \left. \frac{1}{4} h^1{}_{ce;ib} h^1{}^d{}_{a;ie} - \frac{1}{4} h^1{}_{be;ic} h^1{}^d{}_{a;ie} + \frac{1}{4} h^1{}_{bc;ie} h^1{}^d{}_{a;ie} - \frac{1}{4} h^1{}^d{}_{e;ia} h^1{}_{bc}{}^ie + \frac{1}{4} h^1{}_{ce;a} h^1{}^d{}_{b;ie} \right) + \\
& \epsilon^3 \left( -\frac{1}{2} h^1{}^{de} h^1{}^f{}_{e} h^1{}_{cf;bia} - \frac{1}{2} h^1{}^{de} h^1{}^f{}_{e} h^1{}_{bf;cia} + \frac{1}{2} h^1{}^{de} h^1{}^f{}_{e} h^1{}_{bc;fia} - \frac{1}{4} h^1{}^{ef} h^1{}^d{}_{f;ia} h^1{}_{ce;ib} - \right. \\
& \quad \frac{1}{4} h^1{}^{de} h^1{}_{ef;ia} h^1{}^f{}_{c;ib} + \frac{1}{4} h^1{}^{ef} h^1{}_{ce;a} h^1{}^d{}_{f;ib} + \frac{1}{4} h^1{}^{de} h^1{}^f{}_{c;ia} h^1{}_{ef;ib} + \\
& \quad \frac{1}{2} h^1{}^{de} h^1{}^f{}_{e} h^1{}_{cf;a;ib} + \frac{1}{2} h^1{}^{de} h^1{}^f{}_{e} h^1{}_{af;ic;ib} - \frac{1}{2} h^1{}^{de} h^1{}^f{}_{e} h^1{}_{ac;f;ib} + \frac{1}{4} h^1{}^{ef} h^1{}^d{}_{f;ib} h^1{}_{ae;ic} + \\
& \quad \frac{1}{4} h^1{}^{de} h^1{}_{ef;ib} h^1{}^f{}_{a;ic} - \frac{1}{4} h^1{}^{ef} h^1{}^d{}_{f;ia} h^1{}_{be;ic} - \frac{1}{4} h^1{}^{de} h^1{}_{ef;ia} h^1{}^f{}_{b;ic} - \frac{1}{4} h^1{}^{ef} h^1{}_{cf;ib} h^1{}^id{}_{ae} - \\
& \quad \frac{1}{4} h^1{}^{ef} h^1{}_{bf;ic} h^1{}^id{}_{ae} + \frac{1}{4} h^1{}^{ef} h^1{}_{cf;a} h^1{}_{be}{}^id + \frac{1}{4} h^1{}^{ef} h^1{}_{ae;ic} h^1{}_{bf}{}^id - \frac{1}{4} h^1{}^{ef} h^1{}^d{}_{f;ib} h^1{}_{ac;ie} - \\
& \quad \frac{1}{4} h^1{}^{ef} h^1{}_{bf}{}^id h^1{}_{ac;ie} + \frac{1}{4} h^1{}^{ef} h^1{}_{cf;ib} h^1{}^d{}_{a;ie} + \frac{1}{4} h^1{}^{ef} h^1{}_{bf;ic} h^1{}^d{}_{a;ie} - \frac{1}{4} h^1{}^{de} h^1{}^f{}_{c;ib} h^1{}_{af;ie} - \\
& \quad \frac{1}{4} h^1{}^{de} h^1{}^f{}_{b;ic} h^1{}_{af;ie} + \frac{1}{4} h^1{}^{ef} h^1{}^d{}_{f;ia} h^1{}_{bc;ie} - \frac{1}{4} h^1{}^{ef} h^1{}_{cf;ia} h^1{}^d{}_{b;ie} + \frac{1}{4} h^1{}^{de} h^1{}^f{}_{c;ia} h^1{}_{bf;ie} + \\
& \quad \frac{1}{4} h^1{}^{de} h^1{}^f{}_{a;ic} h^1{}_{bf;ie} + \frac{1}{4} h^1{}^{de} h^1{}^f{}_{c;ib} h^1{}_{ae;if} + \frac{1}{4} h^1{}^{de} h^1{}^f{}_{b;ic} h^1{}_{ae;if} + \frac{1}{4} h^1{}^{ef} h^1{}^id{}_{ae} h^1{}_{bc;if} - \\
& \quad \frac{1}{4} h^1{}^{ef} h^1{}^d{}_{a;ie} h^1{}_{bc;if} - \frac{1}{4} h^1{}^{ef} h^1{}_{ae;ic} h^1{}^d{}_{b;if} + \frac{1}{4} h^1{}^{ef} h^1{}_{ac;ie} h^1{}^d{}_{b;if} - \frac{1}{4} h^1{}^{de} h^1{}^f{}_{c;ia} h^1{}_{be;if} - \\
& \quad \left. \frac{1}{4} h^1{}^{de} h^1{}^f{}_{a;ic} h^1{}_{be;if} - \frac{1}{4} h^1{}^{de} h^1{}_{ef;ib} h^1{}^if{}_{ac} - \frac{1}{4} h^1{}^{de} h^1{}_{bf;ie} h^1{}^if{}_{ac} + \frac{1}{4} h^1{}^{de} h^1{}_{be;if} h^1{}^if{}_{ac} + \right)
\end{aligned}$$

$$\begin{aligned}
& \left( \frac{1}{4} h^{1de} h^{1ef;a} h^{1bc}{}^{;f} + \frac{1}{4} h^{1de} h^{1af;e} h^{1bc}{}^{;f} - \frac{1}{4} h^{1de} h^{1ae;f} h^{1bc}{}^{;f} \right) + \\
\epsilon^4 & \left( \frac{1}{2} h^{1de} h^{1ef} h^{1f1} h^{1cf1;b;a} + \frac{1}{2} h^{1de} h^{1ef} h^{1f1} h^{1bf1;c;a} - \frac{1}{2} h^{1de} h^{1ef} h^{1f1} h^{1bc;f1;a} + \right. \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1d}_{f1;a} h^{1cf;b} + \frac{1}{4} h^{1de} h^{1ff1} h^{1ef1;a} h^{1cf;b} + \frac{1}{4} h^{1de} h^{1ef} h^{1ff1;a} h^{1c}_{f1;b} - \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1cf;a} h^{1d}_{f1;b} - \frac{1}{4} h^{1de} h^{1ff1} h^{1cf;a} h^{1ef1;b} - \frac{1}{4} h^{1de} h^{1ef} h^{1c}_{f1;a} h^{1ff1;b} - \\
& \frac{1}{2} h^{1de} h^{1ef} h^{1f1} h^{1cf1;a;b} - \frac{1}{2} h^{1de} h^{1ef} h^{1f1} h^{1af1;c;b} + \frac{1}{2} h^{1de} h^{1ef} h^{1f1} h^{1ac;f1;b} - \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1d}_{f1;b} h^{1af;c} - \frac{1}{4} h^{1de} h^{1ff1} h^{1ef1;b} h^{1af;c} - \frac{1}{4} h^{1de} h^{1ef} h^{1ff1;b} h^{1a}_{f1;c} + \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1d}_{f1;a} h^{1bf;c} + \frac{1}{4} h^{1de} h^{1ff1} h^{1ef1;a} h^{1bf;c} + \frac{1}{4} h^{1de} h^{1ef} h^{1ff1;a} h^{1b}_{f1;c} + \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1cf1;b} h^{1af}{}^{;d} + \frac{1}{4} h^{1ef1} h^{1ef} h^{1bf1;c} h^{1af}{}^{;d} - \frac{1}{4} h^{1ef1} h^{1ef} h^{1cf1;a} h^{1bf}{}^{;d} - \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1af;c} h^{1bf1}{}^{;d} + \frac{1}{4} h^{1de} h^{1ff1} h^{1cf1;b} h^{1af1;e} + \frac{1}{4} h^{1de} h^{1ff1} h^{1bf1;c} h^{1af1;e} - \\
& \frac{1}{4} h^{1de} h^{1ff1} h^{1cf1;a} h^{1bf1;e} - \frac{1}{4} h^{1de} h^{1ff1} h^{1af;c} h^{1bf1;e} + \frac{1}{4} h^{1ef1} h^{1ef} h^{1d}_{f1;b} h^{1ac;f} + \\
& \frac{1}{4} h^{1de} h^{1ff1} h^{1ef1;b} h^{1ac;f} + \frac{1}{4} h^{1ef1} h^{1ef} h^{1bf1}{}^{;d} h^{1ac;f} + \frac{1}{4} h^{1de} h^{1ff1} h^{1bf1;e} h^{1ac;f} - \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1cf1;b} h^{1a}{}^{;f} - \frac{1}{4} h^{1ef1} h^{1ef} h^{1bf1;c} h^{1a}{}^{;f} + \frac{1}{4} h^{1de} h^{1ef} h^{1c}_{f1;b} h^{1af1;f} + \\
& \frac{1}{4} h^{1de} h^{1ef} h^{1b}_{f1;c} h^{1af1;f} - \frac{1}{4} h^{1ef1} h^{1ef} h^{1d}_{f1;a} h^{1bc;f} - \frac{1}{4} h^{1de} h^{1ff1} h^{1ef1;a} h^{1bc;f} - \\
& \frac{1}{4} h^{1de} h^{1ff1} h^{1af1;e} h^{1bc;f} + \frac{1}{4} h^{1ef1} h^{1ef} h^{1cf1;a} h^{1b}{}^{;f} - \frac{1}{4} h^{1de} h^{1ef} h^{1c}_{f1;a} h^{1bf1;f} - \\
& \frac{1}{4} h^{1de} h^{1ef} h^{1a}_{f1;c} h^{1bf1;f} - \frac{1}{4} h^{1de} h^{1ff1} h^{1cf1;b} h^{1ae;f1} - \frac{1}{4} h^{1de} h^{1ff1} h^{1bf1;c} h^{1ae;f1} + \\
& \frac{1}{4} h^{1de} h^{1ff1} h^{1bc;f} h^{1ae;f1} - \frac{1}{4} h^{1de} h^{1ef} h^{1c}_{f1;b} h^{1af;f1} - \frac{1}{4} h^{1de} h^{1ef} h^{1b}_{f1;c} h^{1af;f1} - \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1af}{}^{;d} h^{1bc;f1} + \frac{1}{4} h^{1ef1} h^{1ef} h^{1a}{}^{;f} h^{1bc;f1} + \frac{1}{4} h^{1ef1} h^{1ef} h^{1af;c} h^{1b}{}^{;f1} - \\
& \frac{1}{4} h^{1ef1} h^{1ef} h^{1ac;f} h^{1b}{}^{;f1} + \frac{1}{4} h^{1de} h^{1ff1} h^{1cf1;a} h^{1be;f1} + \frac{1}{4} h^{1de} h^{1ff1} h^{1af;c} h^{1be;f1} - \\
& \frac{1}{4} h^{1de} h^{1ff1} h^{1ac;f} h^{1be;f1} + \frac{1}{4} h^{1de} h^{1ef} h^{1c}_{f1;a} h^{1bf;f1} + \frac{1}{4} h^{1de} h^{1ef} h^{1a}_{f1;c} h^{1bf;f1} + \\
& \frac{1}{4} h^{1de} h^{1ef} h^{1ff1;b} h^{1ac}{}^{;f1} + \frac{1}{4} h^{1de} h^{1ef} h^{1bf1;f} h^{1ac}{}^{;f1} - \frac{1}{4} h^{1de} h^{1ef} h^{1bf;f1} h^{1ac}{}^{;f1} - \\
& \left. \frac{1}{4} h^{1de} h^{1ef} h^{1ff1;a} h^{1bc}{}^{;f1} - \frac{1}{4} h^{1de} h^{1ef} h^{1af1;f} h^{1bc}{}^{;f1} + \frac{1}{4} h^{1de} h^{1ef} h^{1af;f1} h^{1bc}{}^{;f1} \right)
\end{aligned}$$

## ■ 7. Timings

This section performs a number of experiments to analyze the efficiency of the system. All timings have been performed with a Pentium IV 3GHz Linux box with 2Gbytes of RAM running *Mathematica* 5.2.

---

We load some Graphics ` packages

```
In[182]:=
<< Graphics`Graphics`
```

```
In[183]:=
<< Graphics`MultipleListPlot`
```

## 6.1. Timings with perturbations of the Riemann tensor

We compute the perturbations of the Riemann tensor in terms of metric perturbations up to tenth order. We measure the time `ExpandPerturbation` takes to calculate each order and make a plot of these timings versus the perturbative order. We also consider the the graph of the number of terms versus perturbative order.

---

Perturbations of the Riemann tensor up to tenth order. Save the results of timing and the perturbative expression itself in pairs (ti, ri).

```
In[184]:=
  {t1, r1} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 1]] // ToCanonical];

In[185]:=
  {t2, r2} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 2]] // ToCanonical];
    0.244015 Second

In[186]:=
  {t3, r3} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 3]] // ToCanonical];
    0.888055 Second

In[187]:=
  {t4, r4} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 4]] // ToCanonical];
    2.77217 Second

In[188]:=
  {t5, r5} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 5]] // ToCanonical];
    7.9685 Second

In[189]:=
  {t6, r6} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 6]] // ToCanonical];
    21.5853 Second

In[190]:=
  {t7, r7} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 7]] // ToCanonical];
    60.7638 Second

In[191]:=
  {t8, r8} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 8]] // ToCanonical];
    159.066 Second

In[192]:=
  {t9, r9} = AbsoluteTiming[
    ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 9]] // ToCanonical];
    408.962 Second
```

```
In[193]:=
{t10, r10} = AbsoluteTiming[
  ExpandPerturbation[Perturbation[RiemannCD[-a, -b, -c, d], 10]] // ToCanonical];
1134.67 Second
```

---

The following is the list of the times measured above

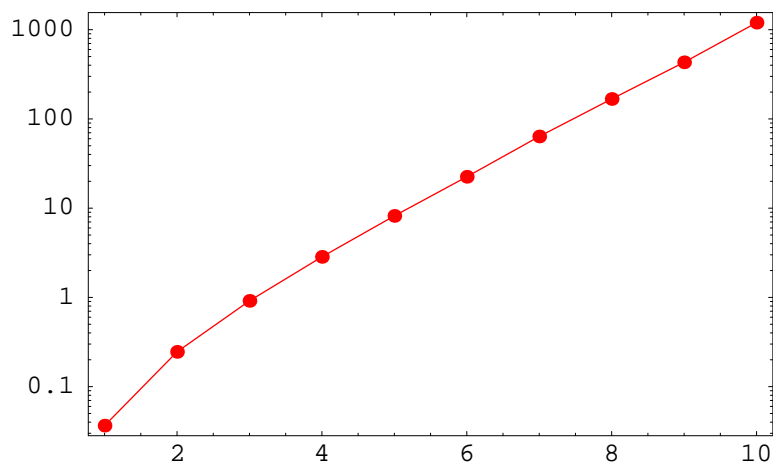
```
In[194]:=
times = {t1, t2, t3, t4, t5, t6, t7, t8, t9, t10} /. Second → 1
```

```
Out[194]=
{0.036555, 0.245376, 0.914745, 2.846096, 8.191186,
 22.481673, 63.588681, 167.305122, 431.159441, 1199.161759}
```

---

We plot then in a logarithmic plot

```
In[195]:=
timesplot = Show[LogListPlot[times, PlotStyle → Hue[0],
  Frame → True, Axes → False, PlotJoined → True, DisplayFunction → Identity],
  LogListPlot[times, PlotStyle → {Hue[0], PointSize[0.02]}, Frame → True,
  Axes → False, SymbolShape → PlotSymbol[Star, Filled → False],
  DisplayFunction → Identity], DisplayFunction → $DisplayFunction]
```



```
Out[195]=
- Graphics -
```

---

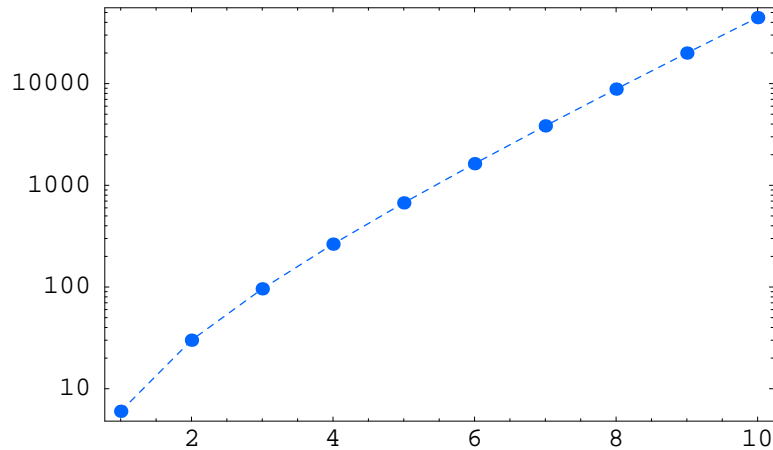
This is the list of the lengths of each perturbative expression

```
In[196]:=
lengths = Length/@Expand[{r1, r2, r3, r4, r5, r6, r7, r8, r9, r10}]
1.38809 Second
```

```
Out[196]=
{6, 30, 96, 264, 672, 1632, 3840, 8832, 19968, 44544}
```

We also plot then in a logarithmic plot

```
In[197]:=
lengthsplot = Show[LogListPlot[lengths, PlotStyle → {Hue[0.6], Dashing[{0.01, 0.01]}],
  Frame → True, Axes → False, PlotJoined → True, DisplayFunction → Identity],
  LogListPlot[lengths, PlotStyle → {Hue[0.6], PointSize[0.02]}, Frame → True,
  Axes → False, DisplayFunction → Identity], DisplayFunction → $DisplayFunction]
```



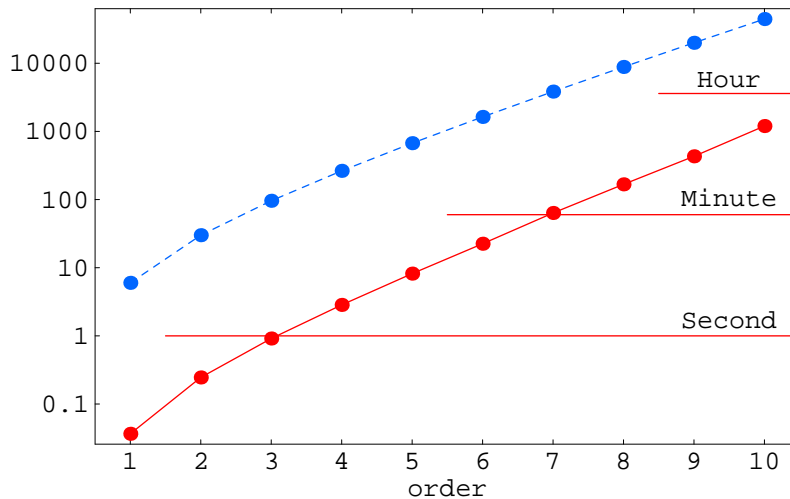
```
Out[197]=
- Graphics -
```

Prepare ticks distribution:

```
In[198]:=
tickat[n_] := {n, n, {0.00625, 0}, {GrayLevel[0], AbsoluteThickness[0.25]}};
ticks1 = tickat /@ Range[10];
logtickat[n_] := {Log[10, n], If[n < 1, N[n], n],
  {0.00625, 0}, {GrayLevel[0], AbsoluteThickness[0.25]}};
ticks2 = logtickat /@ (10^Range[-4, 7]);
tickat[n_] := {n, "", {0.00625, 0}, {GrayLevel[0], AbsoluteThickness[0.25]}};
ticks3 = tickat /@ Range[10];
logtickat[n_] :=
  {Log[10, n], "", {0.00625, 0}, {GrayLevel[0], AbsoluteThickness[0.25]}};
ticks4 = logtickat /@ (10^Range[-4, 7]);
ticks = {ticks1, ticks2, ticks3, ticks4};
```

and show both graphics together

```
In[207]:=
plot = Show[timesplot,
Graphics[{Hue[0], Line[{{1.5, 0}, {10.5, 0}}]}],
Graphics[{Hue[0], Line[{{5.5, Log[10, 60]}, {10.5, Log[10, 60]}]}]},
Graphics[{Hue[0], Line[{{8.5, Log[10, 3600]}, {10.5, Log[10, 3600]}]}]}],
lengthsplot, Graphics[Text["Second", {9.5, 0.2}]],
Graphics[Text["Minute", {9.5, 2}]], Graphics[Text["Hour", {9.5, 3.75}]],
PlotRange -> {{0.5, 10.5}, All}, ImageSize -> 300,
FrameLabel -> {order, None}, FrameTicks -> ticks]
```



```
Out[207]=
- Graphics -
```

## 6.2. Timings with the Leibnitz rule

We now apply Perturbation on products of different numbers of objects. We then graphically represent the time it takes for the expression to be expanded versus the perturbative order and the number of factors in the product. The  $n$ -th order perturbation of the Leibnitz rule is implemented in the context of xPert through a closed formula. We will see that this is faster than the chain rule, that is implemented essentially via the *Mathematica* command Dt.

Driver to average a number of experiments. For timings below a second the computation is repeated a number of times:

```
In[208]:=
average[lists_List] := Plus @@ lists / Length[lists];
threshold = 1;
experiment[number_Integer, order_Integer] :=
Module[{timing = experiment1[number, order], repeat},
If[timing[[1, 1]] < threshold,
repeat = Max[1, Ceiling[threshold / timing[[1, 1]]]];
average@Table[experiment1[number, order], {repeat}],
timing] /. Second -> 1
];
```

Generate products of  $n > 1$  arguments:

```
In[211]:=
Table[prod[n] = Times @@ Table[Unique[x], {n}], {n, 1, 10}];
```



---

and measure the time Perturbation takes to act on them, and the number of terms produced:

```

In[212]:=
  experiment1[number_, 1] :=
    AbsoluteTiming[Length@Expand@Perturbation[prod[number]]];
  experiment1[number_, order_] :=
    AbsoluteTiming[Length@Expand@Perturbation[prod[number], order]];

In[214]:=
  Llist[1] = experiment[#, 1] & /@ Range[2, 10]

      8.67654 Second

Out[214]=
  {{0.000211, 2}, {0.000332, 3}, {0.000558, 4}, {0.000827, 5},
   {0.001268, 6}, {0.001962, 7}, {0.003622, 8}, {0.005935, 9}, {0.010834, 10}}

In[215]:=
  Llist[2] = experiment[#, 2] & /@ Range[2, 10]

      8.74455 Second

Out[215]=
  {{0.000267, 3}, {0.000582, 6}, {0.001080, 10}, {0.001836, 15},
   {0.002867, 21}, {0.004534, 28}, {0.007061, 36}, {0.011499, 45}, {0.017347, 55}}

In[216]:=
  Llist[3] = experiment[#, 3] & /@ Range[2, 10]

      8.20451 Second

Out[216]=
  {{0.000333, 4}, {0.000897, 10}, {0.001975, 20}, {0.003841, 35}, {0.006783, 56},
   {0.011084, 84}, {0.018257, 120}, {0.028401, 165}, {0.045301, 220}}

In[217]:=
  Llist[4] = experiment[#, 4] & /@ Range[2, 10]

      8.75255 Second

Out[217]=
  {{0.000399, 5}, {0.001293, 15}, {0.003245, 35}, {0.007178, 70}, {0.015798, 126},
   {0.026700, 210}, {0.046359, 330}, {0.074470, 495}, {0.117205, 715}}

In[218]:=
  Llist[5] = experiment[#, 5] & /@ Range[2, 10]

      9.81261 Second

Out[218]=
  {{0.000461, 6}, {0.001817, 21}, {0.005189, 56}, {0.013916, 126}, {0.030750, 252},
   {0.060912, 462}, {0.127107, 792}, {0.219005, 1287}, {0.352478, 2002}}

In[219]:=
  Llist[6] = experiment[#, 6] & /@ Range[2, 10]

      11.5687 Second

Out[219]=
  {{0.000597, 7}, {0.002702, 28}, {0.008821, 84}, {0.025633, 210}, {0.057140, 462},
   {0.136063, 924}, {0.260038, 1716}, {0.510616, 3003}, {0.951655, 5005}}

```

In[220]:=

```
Llist[7] = experiment[#, 7] & /@Range[2, 10]
```

11.6247 Second

Out[220]=

```
{ {0.000694, 8}, {0.003526, 36}, {0.012212, 120}, {0.038457, 330}, {0.102152, 792},
  {0.246908, 1716}, {0.551533, 3432}, {1.080255, 6435}, {2.041409, 11440} }
```

In[221]:=

```
Llist[8] = experiment[#, 8] & /@Range[2, 10]
```

15.277 Second

Out[221]=

```
{ {0.000898, 9}, {0.004340, 45}, {0.017768, 165}, {0.052049, 495}, {0.156999, 1287},
  {0.401607, 3003}, {0.906162, 6435}, {1.957846, 12870}, {3.988170, 24310} }
```

In[222]:=

```
Llist[9] = experiment[#, 9] & /@Range[2, 10]
```

22.4534 Second

Out[222]=

```
{ {0.000821, 10}, {0.005109, 55}, {0.023245, 220}, {0.085614, 715}, {0.274254, 2002},
  {0.735193, 5005}, {1.903072, 11440}, {4.452425, 24310}, {8.864111, 48620} }
```

In[223]:=

```
Llist[10] = experiment[#, 10] & /@Range[2, 10]
```

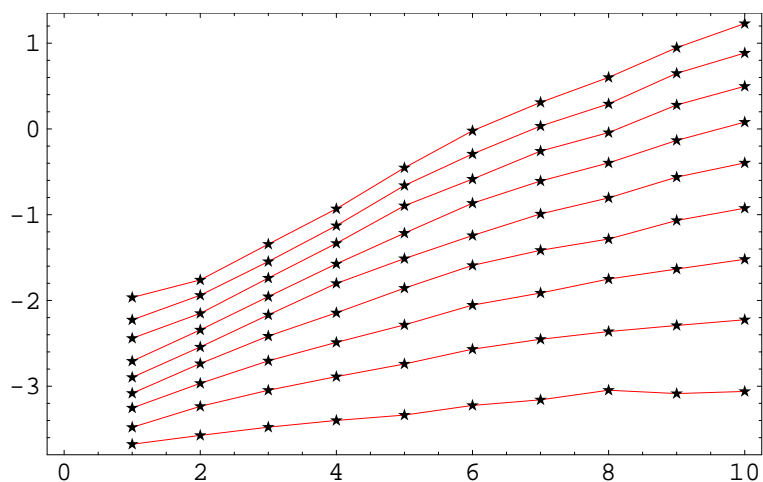
34.3101 Second

Out[223]=

```
{ {0.000869, 11}, {0.005958, 66}, {0.030211, 286}, {0.118979, 1001}, {0.401679, 3003},
  {1.200730, 8008}, {3.144754, 19448}, {7.678046, 43758}, {16.930981, 92378} }
```

In[224]:=

```
Ltimingsplot = MultipleListPlot[
  Transpose@Log[10, Map[First, Array[Llist, {10}], {2}]], PlotJoined → True,
  Axes → False, Frame → True, SymbolShape → PlotSymbol[Star], PlotStyle → Hue[0]]
```

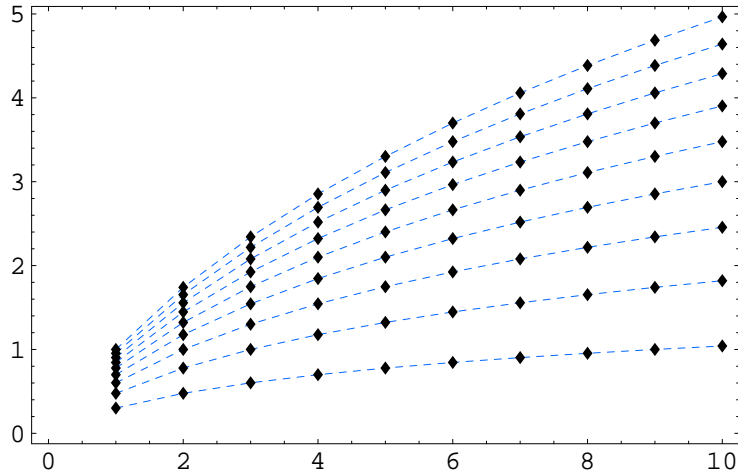


Out[224]=

- Graphics -

```
In[225]:=
```

```
Ltermsplot = MultipleListPlot[Transpose@Log[10, Map[Last, Array[Llist, {10}], {2}]],
  PlotJoined → True, Axes → False, Frame → True, SymbolShape → PlotSymbol[Diamond],
  PlotStyle → {{Hue[0.6], Dashing[{0.01, 0.01]}}}]
```

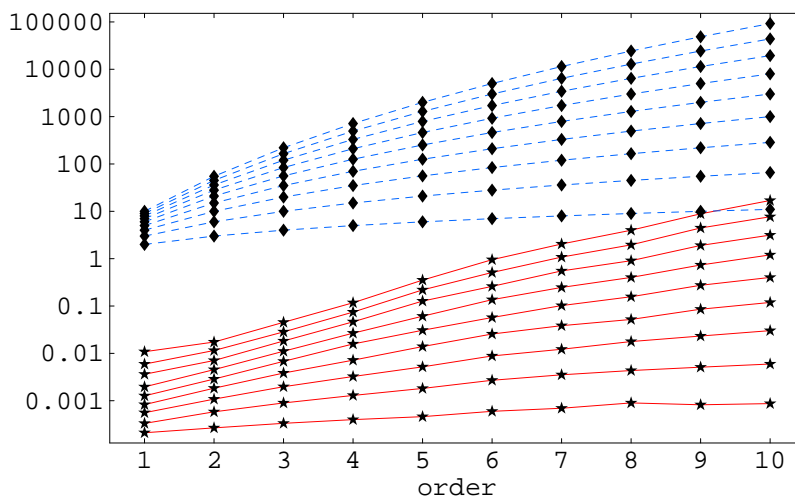


```
Out[225]=
```

```
- Graphics -
```

```
In[226]:=
```

```
Lplot = Show[Ltimingsplot, Ltermsplot, FrameLabel → {order, None},
  FrameTicks → ticks, PlotRange → {{0.5, 10.5}, All}, ImageSize → 300]
```



```
Out[226]=
```

```
- Graphics -
```

## 6.2. Timings with the chain rule

In this section we measure the time it takes to the command `Perturbation` to act on a scalar function of up to 10 arguments through the chain rule. We analyze it up to the fifth perturbative order. In the context of `xPert`, we have implemented the  $n$ th order perturbation of the chain rule through the `Mathematica` command `Dt`.

---

Generate functions of n arguments and perturb them:

```
In[227]:=
  Table[func[n] = F@@Table[Unique[x], {n}], {n, 1, 10}];
```

```
In[228]:=
  experiment1[number_, 1] :=
    AbsoluteTiming[Length@Expand@Perturbation[func[number]]];
  experiment1[number_, order_] :=
    AbsoluteTiming[Length@Expand@Perturbation[func[number], order]];
```

```
In[230]:=
  Flist[1] = experiment[#, 1] & /@Range[10]

  9.70461 Second
```

```
Out[230]=
  {{0.000401, 2}, {0.000665, 2}, {0.000944, 3}, {0.001162, 4}, {0.001436, 5},
  {0.001694, 6}, {0.001974, 7}, {0.002245, 8}, {0.002539, 9}, {0.002852, 10}}
```

---

The first case is a product of two objects and hence its length must be corrected:

```
In[231]:=
  Flist[1] = ReplacePart[Flist[1], 1, {1, 2}]
```

```
Out[231]=
  {{0.000401, 1}, {0.000665, 2}, {0.000944, 3}, {0.001162, 4}, {0.001436, 5},
  {0.001694, 6}, {0.001974, 7}, {0.002245, 8}, {0.002539, 9}, {0.002852, 10}}
```

```
In[232]:=
  Flist[2] = experiment[#, 2] & /@Range[10]

  9.07657 Second
```

```
Out[232]=
  {{0.000587, 2}, {0.001233, 5}, {0.001967, 9}, {0.002857, 14}, {0.003935, 20},
  {0.005273, 27}, {0.006778, 35}, {0.008543, 44}, {0.010484, 54}, {0.012736, 65}}
```

```
In[233]:=
  Flist[3] = experiment[#, 3] & /@Range[10]

  9.97662 Second
```

```
Out[233]=
  {{0.000870, 3}, {0.002557, 10}, {0.005340, 22}, {0.009621, 40}, {0.015884, 65},
  {0.024725, 98}, {0.036190, 140}, {0.051157, 192}, {0.071447, 255}, {0.096454, 330}}
```

```
In[234]:=
  Flist[4] = experiment[#, 4] & /@Range[10]

  13.8609 Second
```

```
Out[234]=
  {{0.001459, 5}, {0.006338, 20}, {0.016207, 51}, {0.037493, 105}, {0.076312, 190},
  {0.128472, 315}, {0.216543, 490}, {0.386452, 726}, {0.697105, 1035}, {0.947616, 1430}}
```

```
In[235]:=
  Flist[5] = experiment[#, 5] & /@Range[10]

26.4657 Second

Out[235]=
  {{0.002552, 7}, {0.016583, 36}, {0.059041, 108},
  {0.168279, 252}, {0.395561, 506}, {0.857346, 918}, {1.602924, 1547},
  {2.933435, 2464}, {5.042421, 3753}, {8.547611, 5512}}

In[236]:=
  Flist[6] = experiment[#, 6] & /@Range[10]

174.911 Second

Out[236]=
  {{0.004004, 11}, {0.040087, 65}, {0.195557, 221},
  {0.701540, 574}, {2.025705, 1265}, {4.913787, 2492}, {10.945729, 4522},
  {23.051759, 7704}, {44.932045, 12483}, {83.759705, 19415}}

In[237]:=
  Flist[7] = experiment[#, 7] & /@Range[8]

330.277 Second

Out[237]=
  {{0.005622, 15}, {0.089162, 110}, {0.681502, 429}, {3.055511, 1240},
  {10.884433, 2990}, {33.225293, 6372}, {82.240041, 12405}, {196.821582, 22528}}

In[238]:=
  Flist[8] = experiment[#, 8] & /@Range[6]

302.719 Second

Out[238]=
  {{0.010850, 22}, {0.267966, 185}, {2.461819, 810},
  {14.321746, 2580}, {61.641596, 6765}, {222.491062, 15525}}

In[239]:=
  Flist[9] = experiment[#, 9] & /@Range[5]

415.602 Second

Out[239]=
  {{0.013778, 30}, {0.672458, 300},
  {8.219750, 1479}, {61.320062, 5180}, {348.810322, 14725}}

In[240]:=
  Flist[10] = experiment[#, 10] & /@Range[4]

381.524 Second

Out[240]=
  {{0.215057, 42}, {2.882951, 481}, {34.553778, 2640}, {349.965110, 10108}}

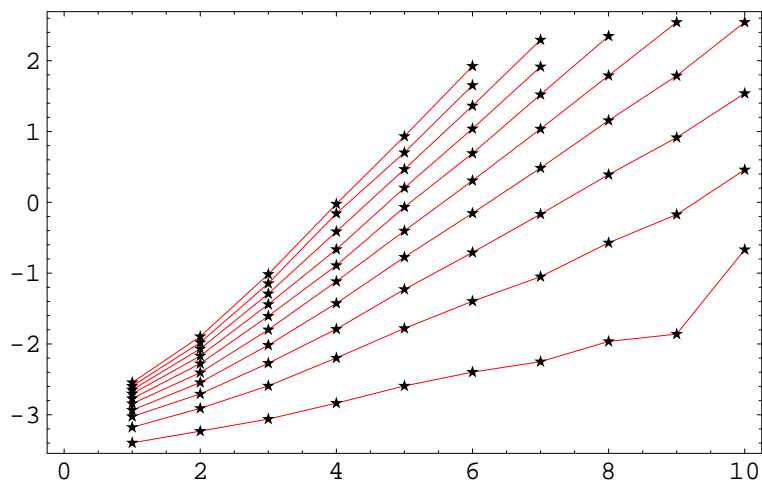
In[241]:=
  MaxMemoryUsed[]

Out[241]=
  1048894992
```

Graphs:

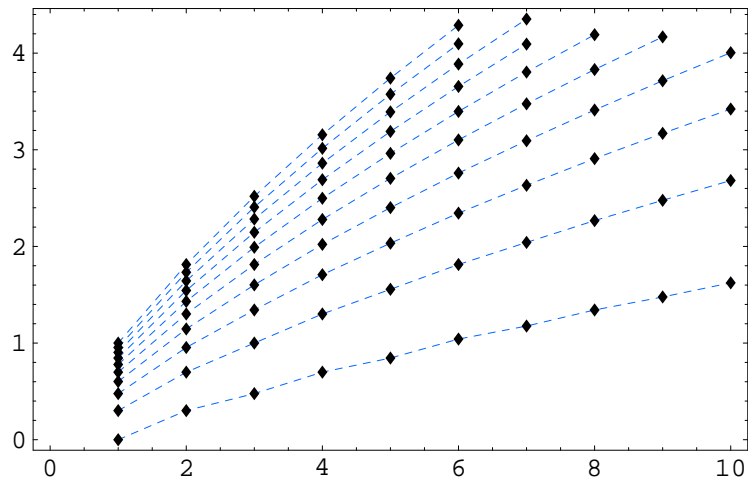
```
In[242]:=
tlist = Transpose[PadRight[#, 10, Null] & /@ Array[Flist, {10}]] /. Null -> Sequence[];
```

```
In[243]:=
Ftimingsplot = MultipleListPlot[Log[10, Map[First, tlist, {2}]], PlotJoined -> True,
  Axes -> False, Frame -> True, SymbolShape -> PlotSymbol[Star], PlotStyle -> Hue[0]]
```



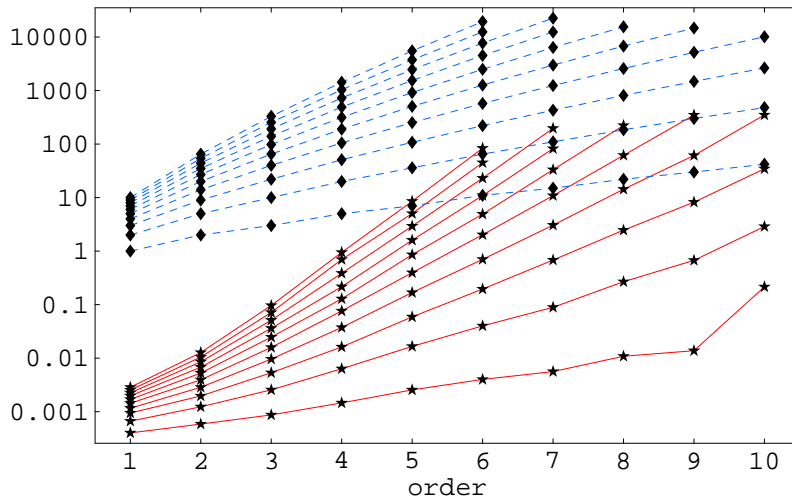
```
Out[243]=
- Graphics -
```

```
In[244]:=
Ftermsplot = MultipleListPlot[Log[10, Map[Last, tlist, {2}]],
  PlotJoined -> True, Axes -> False, Frame -> True, SymbolShape -> PlotSymbol[Diamond],
  PlotStyle -> {{Hue[0.6], Dashing[{0.01, 0.01}]}}]
```



```
Out[244]=
- Graphics -
```

```
In[245]:=
Fplot = Show[Ftimingsplot, Ftermsplot, FrameLabel -> {order, None},
  FrameTicks -> ticks, PlotRange -> {{0.5, 10.5}, All}, ImageSize -> 300]
```



```
Out[245]=
- Graphics -
```

## ■ 8. Final comments

**Note:** For further information about xPert `'`, and to be kept informed about new releases, you may contact the authors electronically at [brizuela@iem.cfmac.csic.es](mailto:brizuela@iem.cfmac.csic.es), [jmm@iem.cfmac.csic.es](mailto:jmm@iem.cfmac.csic.es) and [mena@iem.cfmac.csic.es](mailto:mena@iem.cfmac.csic.es). Suggestions and comments are welcome and very much appreciated!

This is xPertDoc.nb, the docfile of xPert `'`, currently in version 1.0.0.

```
In[246]:=
?xAct`xPert`*
```

**xAct`xPert`**

DefMetricPerturbation	GaugeChange	Perturbed
DefTensorPerturbation	GeneralPerturbation	\$PerturbationParameter
Disclaimer	Perturbation	\$Version
ExpandPerturbation	PerturbationOrder	\$xTensorVersionExpected

0.180011 Second